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# Effect of Inherent Damping of the Series Elastic Element on Rendering Performance and Passivity of Interaction Control

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*We study a realistic model of series elastic actuation (SEA) under velocity-sourced impedance control (VSIC), where the inherent damping of the series elastic element is considered during the analysis, even when only the elasticity of the series damped elastic element is used to estimate the interaction forces. We establish a fundamental rendering limitation when the viscous damping of the physical filter is considered in the plant model and prove that passive rendering of stiffness levels that are higher than the stiffness of the physical filter, as well as passive rendering of Voigt models whose damping levels exceed the physical damping of the plant, are possible. We introduce passive physical equivalents of the closed-loop SEA systems with inherent series damping while rendering Kelvin-Voigt, spring, and null impedance models to provide an intuitive understanding of the passivity bounds and to enable rigorous comparisons of rendering performance among various closed-loop systems with different plant models (including or omitting the series damping) and/or controllers (utilizing different interaction force estimates). We present a comprehensive set of experiments to verify our results and demonstrate the effect of including/omitting the damping of the physical filter in the model of SEA. [DOI: 10.1115/1.4068463]*

**Keywords:** physical human–robot interaction, interaction control, series elastic actuation, haptic rendering, fundamental limits of performance, Cauchy's residue theorem, physical realizations, coupled stability, effective impedance

## 1 Introduction

Series elastic actuation (SEA) is an interaction control paradigm that can provide safe and natural interactions with high stability robustness, and rendering fidelity. SEA involves introducing a compliant element between the actuator and the interaction port and utilizing the model of this compliance for closed-loop force control [1–3]. The purposeful introduction of the compliant element relaxes the stringent stability constraints on the controller gains that arise due to sensor-actuator noncollocation and actuator bandwidth restrictions [4–6] and provides high stability robustness for interaction control. Furthermore, high rendering performance can be achieved by actively compensating for the dynamics of the compliant element through the use of its model. On the negative side, the control effort required to compensate for the compliant element increases rapidly for high-frequency interactions, resulting in actuator saturation and limiting the control bandwidth of SEA.

Velocity-sourced impedance control (VSIC) is a widely used controller for SEA that utilizes a cascaded architecture. The inner

motion control loop of VSIC effectively compensates for the parasitic forces, resulting in favorable rendering performance [1,3,7,8]. Additionally, VSIC does not require a dynamic model of the plant, allowing for empirically tuned controller gains.

In the literature, extensive research has been conducted on establishing passivity conditions of SEA under VSIC [9–12], and the necessary and sufficient conditions for passivity have been determined for linear spring and null impedance rendering [13,14]. It has also been proven that the passively renderable stiffness of SEA under any causal controller is limited by the physical stiffness of its compliant element [15]. Moreover, it has been shown that while Kelvin–Voigt (abbreviated as Voigt) models (linear spring and damper connected in parallel) cannot be passively rendered using SEA under VSIC with positive controller gains [11], passive Voigt model rendering that compensates for plant damping is possible with SEA under VSIC through the use of negative controller gains [14]. The passivity analysis of SEA has also been extended to model-based control, such as model reference force control [16].

Series elastic actuation relies on the assumption that its compliant element is an ideal spring and estimates the interaction forces through the deflections of this spring element for use in the closed-loop force control. However, this assumption is unrealistic, as some form of energy dissipation is inherent to all physical spring

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implementations. Furthermore, this dissipation is parallel to the series elastic element and cannot be included in the environment model during the coupled stability analysis, as commonly done for the end-effector mass. The inclusion of the physical dissipation of the elastic element in the model of the SEA has major consequences on the high-frequency response of the system model, as seen from the interaction port. One such model extends the SEA paradigm to series *damped* elastic actuation (SDEA) by incorporating a viscous dissipation element parallel to the series elastic element [17–21].

The frequency responses of SDEA and SEA are significantly different. When the causal controllers roll off, the dynamics of the uncontrolled plant are recovered; hence, the high-frequency responses of SDEA/SEA as seen from the interaction port are dominated by the dynamics of their serially attached physical filters. Accordingly, at high frequencies, the dynamics of SDEA are dominated by its damping (even for low damping coefficients), while SEA behaves as a linear spring. Consequently, SDEA possesses *distinct* stability properties and haptic rendering performance compared to SEA since the high-frequency dynamics of the plant impose fundamental limitations on the achievable control performance of any closed-loop system.

Series *damped* elastic actuation can estimate interaction forces through the sum of the forces induced on the physical spring and damper pair. Incorporating the series physical damping into the plant and its model within the closed-loop controller has been demonstrated to enhance the force control bandwidth of SEA [17]. Moreover, this approach has been shown to provide additional benefits, such as enhancing energy efficiency [19], reducing unwanted oscillations [20], mitigating the requirement for derivative control terms [21], and enabling passive Voigt model rendering with positive controller gains [14,22,23].

Passivity analysis of SDEA has been studied in the literature, but the closed-form analytic passivity conditions derived from these studies are complex and difficult to interpret [18,22–24]. Previous research has shown that Voigt models can be passively rendered using SDEA. For example, the passive range of virtual stiffness and damping parameters for SDEA under a cascaded impedance controller with an inner torque loop acting on a velocity-compensated plant and load dynamics has been presented in Ref. [24]. Similarly, a passivity analysis of SDEA under an unconventional basic impedance controller has been presented in Ref. [18]. In this controller, a force sensor is employed after the end-effector inertia to measure human interface forces and these force measurements are used in the closed-loop force control, in addition to the series-damped elastic element. The passivity analysis with this controller indicates that a sufficient level of viscous damping is required in the compliant element to ensure the passivity of stiffness rendering.

Mengilli et al. [22] have provided sufficient conditions for the passivity of SDEA under VSIC for the null impedance, linear spring, and Voigt model rendering. Furthermore, they have extended these results to absolute stability analysis and derived necessary and sufficient conditions for two-port passivity of SDEA under VSIC with a virtual coupler [23].

Recently, authors have established necessary and sufficient conditions for the passivity of SEA and SDEA (together abbreviated as S(D)EA) under VSIC while rendering Voigt, spring, and null impedance models [14]. Moreover, the use of passive physical equivalents for S(D)EA under VSIC has been advocated in Ref. [14] to establish an intuitive understanding of the passivity bounds and to highlight the effect of different plant parameters and controller terms on the closed-loop performance. Furthermore, the passive physical equivalents are shown to be instrumental in enabling objective comparisons of closed-loop systems featuring different controller architectures. In Ref. [25], authors have extended their analysis to study the effects of low-pass filtering on SEA under VSIC.

This study follows an analysis technique similar to Refs. [14,25], as it relies on frequency domain analysis to derive closed-form analytical solutions for necessary and sufficient conditions that

ensure passivity, and utilizes passive physical equivalents to establish an understanding of the passivity bounds, to determine the parasitic elements, and to compare the effect of different plant and controller dynamics on the coupled stability and haptic rendering performance of the closed-loop system.

On the other hand, this study significantly extends the previously established results on the passivity analysis and passive physical realizations of S(D)EA [14,15,22,23], by establishing a fundamental rendering limitation for SDEA, proving that the inclusion of the damping of physical filter in the model enables passive rendering of Voigt models that exceed the stiffness and damping levels of the plant, and systematically studying the impact of including/omitting the damping force induced on the serial compliant element under VSIC on the rendering performance. Our novel contributions can be listed as follows:

- Utilizing a Cauchy integral, we establish a fundamental rendering limitation for all SDEA systems under causal controllers and prove that the inclusion of viscous damping of the physical filter in the model enables passive rendering of stiffness levels that are higher than the physical stiffness of the serial filter, as well as passive rendering of Voigt models whose damping levels exceed the physical damping of the plant. This relation also explains why Voigt model rendering with SEA is restricted to damping compensation.
- We present necessary and sufficient conditions for the passivity of SDEA under VSIC while rendering Voigt, spring, and null impedance models, when the damping force on the physical filter is *not* used for closed-loop control.
- We derive passive physical equivalents of various closed-loop SDEA systems and rigorously study the effect of omitting/including the damping force induced on the physical filter of SDEA in the system model on the coupled stability and closed-loop rendering performance.
- We provide comprehensive comparisons of rendering performance among various closed-loop systems with different plant models including or omitting the series damping and/or controllers utilizing different interaction force estimates through their passive physical equivalents and K-B plots.
- We demonstrate the validity and applicability of our theoretical results through a comprehensive set of experiments using a custom brake pedal with SDEA.

## 2 Preliminaries

A schematic representation of a single-axis plant with SDEA in the absence of its controller is presented in Fig. 1. The figure illustrates the actuator's reflected inertia  $J_m$ , viscous friction  $B_m$  which includes the motor damping, the physical spring  $K$  and the viscous damper  $B_f$  arranged in parallel between the end effector and the actuator. The actuator and end-effector velocities are represented by  $\omega_m$  and  $\omega_{end}$ , respectively. The symbol  $\tau_m$  denotes the actuator torque, while  $M_{end}$  represents the inertia of the end effector. Consideration is given to a lumped-parameter linear time-invariant (LTI) model, so nonlinear effects, including backlash and actuator saturation, are neglected.

The torque applied to the physical filter, represented as a damped compliant element, is equal to the sum of the torques exerted on the linear spring  $K$  and the viscous damper  $B_f$ . Figure 1 can be regarded

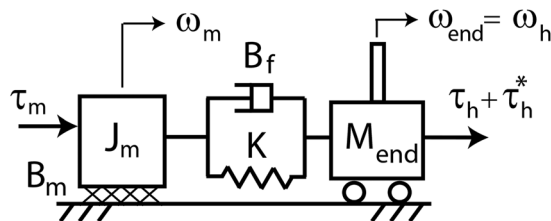


Fig. 1 Schematic representation of an SDEA plant





**Table 1 Comparison of physical realizations of various models of S(D)EA under VSIC**

	Voigt rendering	Spring rendering	Null rendering
SEA <sup>a</sup>	(a)	(b)	(c)
SDEA <sup>a</sup>	(d)	(e)	(f)
SDEA <sub>K<sub>fb</sub></sub>	(g)	(h)	(i)

<sup>a</sup>The realizations of S(D)EA under VSIC presented in (a)–(f) tabs are adapted from [14].

motion control loop is asymptotically stable. Then, the following inequalities serve as necessary and sufficient conditions for establishing the passivity of  $Z_{\text{Voigt}}^{\text{SDEA}_{K_{fb}}}(s)$ .

- $K \geq K_{\text{ref}} \frac{\alpha}{(\alpha+1)} \frac{B_m + G_m - B_f \alpha}{B_m + G_m + B_{\text{ref}} \alpha}$ , and
- $-(B_m + G_m) < \alpha B_{\text{ref}}$ , and
- $0 < \frac{\alpha}{\alpha+1} K_{\text{ref}}$ , and
- $0 < (\alpha + 1)$ , and
- $0 < (B_m + G_m)$ , and
- $-2\sqrt{\gamma} \leq B_f (B_m + G_m + B_{\text{ref}} \alpha) (B_f + B_m + G_m) - \alpha J_m (B_f (K + K_{\text{ref}}) + K B_{\text{ref}})$

where  $\gamma = B_f J_m^2 (K (\alpha + 1) (K (B_m + G_m + B_{\text{ref}} \alpha) + B_f \alpha K_{\text{ref}}) - \alpha K K_{\text{ref}} (B_f + B_m + G_m))$ .

The proof of Proposition 2 is presented in Appendix B.

**Remark 1** (1). If  $B_{\text{ref}}$  ( $B_{\text{ref}}$  and  $K_{\text{ref}}$ ) is set to zero in Eq. (1), the output impedance transfer function becomes the same as the transfer function for spring (null impedance) rendering. Hence, Proposition 2 covers spring and null impedance rendering as its special cases.

Together with Condition (i), a simpler but more conservative set of sufficient conditions can be derived by imposing the following inequality instead of Condition (ii) of Proposition 2:

$$J_m \leq \frac{B_f (B_f + B_m + G_m) (B_m + G_m + B_{\text{ref}} \alpha)}{\alpha (K (B_{\text{ref}} + B_f) + B_f K_{\text{ref}})} \quad (2)$$

**4.2 Passive Physical Equivalent of SDEA<sub>K<sub>fb</sub></sub> Under Velocity-Sourced Impedance Control.** A realization of Eq. (1) characterizing SDEA<sub>K<sub>fb</sub></sub> under VSIC during Voigt model rendering is presented in Table 1(g). The parameters of this realization include  $k_{1v} = K - \frac{K_{\text{ref}} \alpha}{\alpha+1} - \frac{B_f (B_f - B_{\text{ref}} \alpha)}{J_m}$  and  $c_{4v} = \frac{B_m - B_f + G_m + B_{\text{ref}} \alpha - B_f \alpha}{\alpha+1} - \frac{K_{\text{ref}} \alpha (B_m + G_m - B_f \alpha)}{K (\alpha+1)^2}$ . The rest of the terms are relatively complicated; hence, they are presented as a MATLAB script that allows for a numerical means of checking for the non-negativeness of each element<sup>2</sup>.

The third row of Table 1 presents realizations of SDEA<sub>K<sub>fb</sub></sub> under VSIC while rendering Voigt, spring, and null impedance models. Table 1 indicates that there exists continuity among the realizations; the realization of spring rendering and null impedance rendering with SDEA<sub>K<sub>fb</sub></sub> can be recovered from Voigt model rendering with SDEA<sub>K<sub>fb</sub></sub> by setting  $B_{\text{ref}} = 0$  and  $B_{\text{ref}} = K_{\text{ref}} = 0$ , respectively. Realizations of SEA can also be recovered from the realizations of SDEA<sub>K<sub>fb</sub></sub> by setting  $B_f = 0$ .

**4.3 Effective Impedance Analysis.** After removing the rendered stiffness and physical damper pair  $\frac{\alpha}{\alpha+1} K_{\text{ref}} - B_f$  and the serial coupling filter  $k_{1v}$ , the effective impedance of the realization in Table 1(g) indicates that the effective damping converges to  $c_{4v}$  at low frequencies, while it approaches to  $c_{4v} + c_{5v}$  at high frequencies. Accordingly, for these elements,  $c_{4v}$  is the dominant damping at low

<sup>2</sup>The MATLAB script of the parameters of the realization in Table 1g is available for download at [https://hmi.sabanciuniv.edu/SDEAKfb\\_realization.m](https://hmi.sabanciuniv.edu/SDEAKfb_realization.m).

frequencies, and  $c_{5v}$  is added to  $c_{4v}$  as the frequency increases. Similarly, the effective parasitic inductance converges to  $b_{4v}$  at low frequencies, while it approaches zero at high frequencies. While the effect of  $b_{4v}$  becomes higher as the frequency increases, this is balanced by the fact that effective inductance goes to zero at high frequencies.

**4.4 Haptic Rendering Performance.** The physical realization of  $SDEA_{K_{fb}}$  during Voigt model rendering in Table 1(g) indicates three main branches in parallel: a spring-damper pair  $\frac{\alpha}{(\alpha+1)} K_{ref} - B_f$  in parallel, and a branch capturing the dynamics governed by a complex topology of damper-inductance terms that are coupled to the system in series through a spring. The effective impedances at low and high frequencies of the complex topology of damper-inductance terms are provided in the Sec. 4.3.

The parallel spring-damper pair of  $\frac{\alpha}{(\alpha+1)} K_{ref} - B_f$  indicates that  $SDEA_{K_{fb}}$  under VSIC can render the desired spring levels for proper selections of  $K_{ref}$ . Rendering of desired damping levels is more involved as both  $c_{4v}$  and  $c_{5v}$  depend on  $B_{ref}$ . At low frequencies, the effective impedance of the whole system approaches  $B_f + c_{4v} (B_{vir})$ ; hence,  $B_{ref}$  can be selected to render desired damping levels. Even though the presented realization is only valid for positive values of  $c_{4v}$ , passivity conditions indicate that it is possible to render desired damping levels that are lower than the damping of the series elastic element by selecting  $c_{4v}$  negative, such that  $0 > c_{4v} \geq -B_f$ .

Note that the inerter-damping terms are coupled to  $\frac{\alpha}{(\alpha+1)} K_{ref} - B_f$  in parallel through the coupling filter  $k_{1v}$  and this coupling becomes stronger for lower choices of  $K_{ref}$ . As frequency increases,  $c_{5v}$  is added to rendered damping as a parasitic effect. Unlike the case in SDEA realization,  $SDEA_{K_{fb}}$  realization does not have a pure inerter term that dominates the parasitic dynamics at high frequencies; similar to SEA realization,  $SDEA_{K_{fb}}$  realization has an inerter term  $b_{4v}$  that adds frequency dependent inductance at low frequencies.

A comparison of the effective impedances of the realization of  $SDEA_{K_{fb}}$  with SDEA under VSIC indicates that while the effective parasitic inductance of  $SDEA_{K_{fb}}$  goes to zero, the effective parasitic inductance of SDEA goes to  $b_{2v}$  at high frequencies; hence, the parasitic inductance of SDEA is higher than that of  $SDEA_{K_{fb}}$ . A numerical comparison of the effective parasitic damping of  $SDEA_{K_{fb}}$  with SDEA under VSIC is presented in Sec. 5.

**4.5 Comparison of Passivity Bounds of  $SDEA_{K_{fb}}$  With Series Damped Elastic Actuation.** For the simplicity of the analysis, a comparison is made for the case when all controller gains are taken positive. Comparison of the necessary conditions presented for SDEA in Ref. [14] and Condition (i) of Proposition 2 that impose upper bounds on  $K_{ref}$  for  $SDEA_{K_{fb}}$  indicates that the bound for  $SDEA_{K_{fb}}$  is more relaxed as follows:

$$\begin{aligned} K &\geq K_{ref} \frac{\alpha}{(\alpha+1)} \frac{B_m + G_m}{B_m + G_m + B_{ref} \alpha} \\ &\geq K_{ref} \frac{\alpha}{(\alpha+1)} \frac{B_m + G_m - B_f \alpha}{B_m + G_m + B_{ref} \alpha} \end{aligned} \quad (3)$$

Accordingly, Eq. (3) shows that  $SDEA_{K_{fb}}$  can passively render virtual springs that are stiffer than the virtual springs that can be passively rendered by SDEA. In fact, with  $SDEA_{K_{fb}}$  under VSIC, it is possible to passively render virtual stiffness levels that exceed the stiffness of the physical filter.

Comparison of the sufficient conditions of SDEA and  $SDEA_{K_{fb}}$  presented in Ref. [14] and Eq. (2) indicate that

$$\begin{aligned} J_m &\leq \frac{B_f (B_m + G_m + B_{ref} \alpha) [B_m + G_m + B_f (1 + \alpha)]}{(B_f K_{ref} + B_{ref} K) \alpha} \\ &\leq \frac{B_f (B_m + G_m + B_{ref} \alpha) (B_m + G_m + B_f)}{[B_f K_{ref} + K (B_{ref} + B_f)] \alpha} \end{aligned} \quad (4)$$

**Table 2 Plant parameters utilized in simulations**

Parameter (unit)	Value	Parameter (unit)	Value
$J_m$ (kgm <sup>2</sup> )	0.002	$B_m$ (Nms/rad)	1.22
$K$ (Nm/rad)	360	$B_f$ (Nms/rad)	0.35

Consequently, while  $SDEA_{K_{fb}}$  can passively render a larger range of virtual springs compared to SDEA according to Eq. (3),  $SDEA_{K_{fb}}$  also possesses a more strict passivity bound on  $J_m$  as expressed in Eq. (4).

## 5 Evaluations of Rendering Performance

This section studies the haptic rendering performance by providing comparisons of Bode plots of the S(D)EA models under VSIC during Voigt model, linear spring, and null impedance rendering.

Table 2 presents the plant parameters utilized in the simulations. The controller gains of VSIC are taken as  $G_m = 10$  Nms/rad and  $G_t = 5$  rad/(sNm). The Voigt model parameters are chosen as  $K_{ref} = 150$  Nm/rad and  $B_{ref} = 0.1$  Nms/rad, respectively. For SEA,  $B_{ref}$  is set to  $-0.0307$  Nms/rad such that  $c_{1v} = 0.1$  Nms/rad.

Performance evaluations of rendering with S(D)EA under VSIC, with the insights provided by their passive physical equivalents, have been presented in Ref. [14].

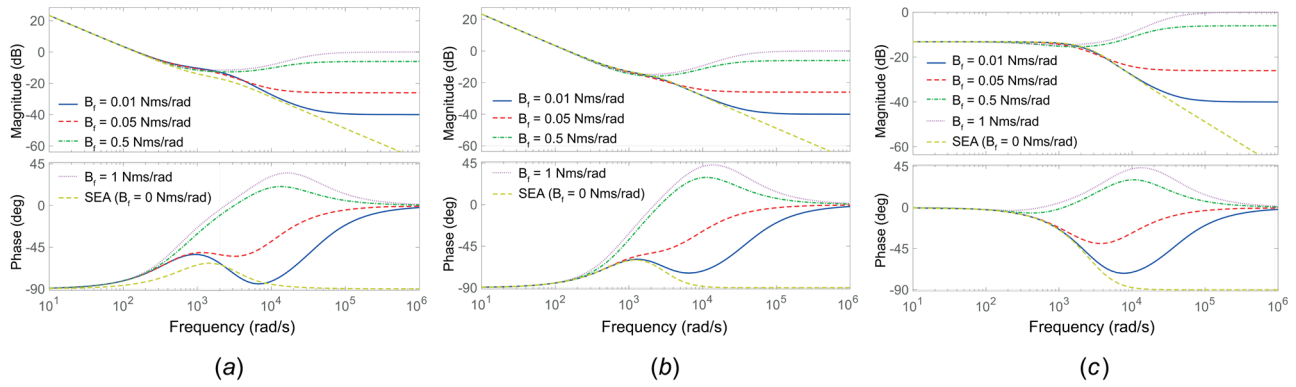
**5.1 Effect Physical Filter Damping on Series Damped Elastic Actuation and  $SDEA_{K_{fb}}$ .** Tables 1(d)–1(f) present the passive physical equivalents of SDEA under VSIC during Voigt model, spring, and null impedance rendering, respectively. Figure 3 exhibits the performance of Voigt model, spring, and null impedance rendering for SDEA, while employing proportional controllers. Results in Fig. 3(a) indicate that the Voigt model rendering performs poorly for the lowest value of  $B_f$ , but there is no noticeable distinction between rendering performance for sufficiently high  $B_f$ , such as when  $B_f = 0.5$  Nms/rad and 1 Nms/rad. Moreover, Fig. 3(b) reveals that when  $B_f$  is minimum, the performance bandwidth of spring rendering is the largest, but the fidelity of spring rendering diminishes for larger virtual springs. Similarly, Fig. 3(c) shows that the null impedance rendering performs poorly for the lowest  $B_f$ , but there is no significant difference between the rendering performance when  $B_f = 0.5$  Nms/rad and 1 Nms/rad.

Tables 1(g)–1(i) present the passive physical equivalents of  $SDEA_{K_{fb}}$  under VSIC during Voigt model, spring, and null impedance rendering, respectively. All passive physical equivalents have the damping of the physical filter  $B_f$  as a term parallel to all other terms. Passive physical equivalents make it explicit that  $B_f$  is directly included in the effective damping of  $SDEA_{K_{fb}}$ .

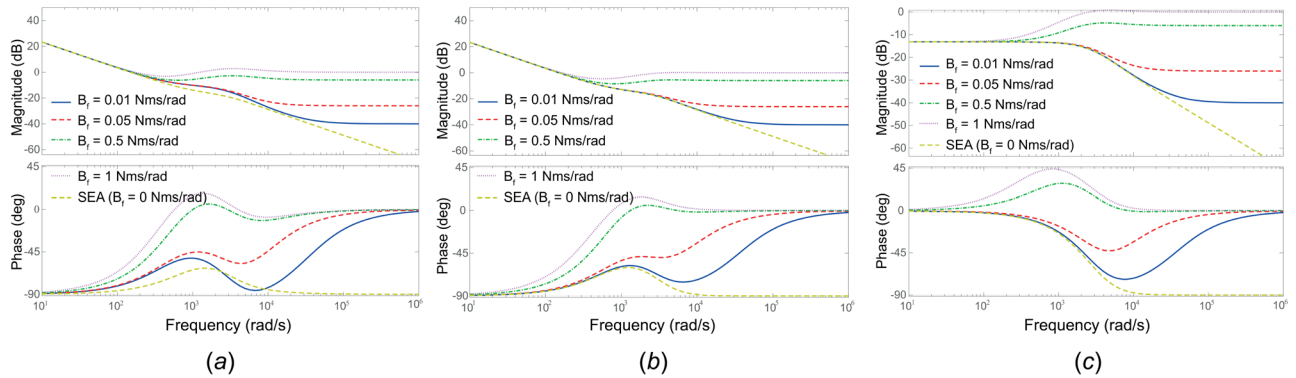
Figure 4 illustrates the effects of varying  $B_f$  on the performance of Voigt model, spring, and null impedance rendering for  $SDEA_{K_{fb}}$ . Figures 4(a) and 4(b) indicate that decreasing  $B_f$  results in a larger performance bandwidth for virtual spring rendering. Similarly, Fig. 4(c) demonstrates that decreasing  $B_f$  leads to improvements in the null impedance rendering.

Figures 3 and 4 illustrate that the rendering performance of SDEA increases for higher  $B_f$ , while the performance of  $SDEA_{K_{fb}}$  decreases for higher  $B_f$ . This observation can be explained by examining the realizations of both systems. Specifically,  $B_f$  acts as an uncontrolled term that runs parallel to controlled damping of  $c_{4v}$  for  $SDEA_{K_{fb}}$  and higher  $B_f$  values require larger adjustments to  $c_{4v}$  to overcome effects of  $B_f$ . Conversely, for SDEA, there is a direct control action on the  $B_f$  that can be used to improve the rendering performance.

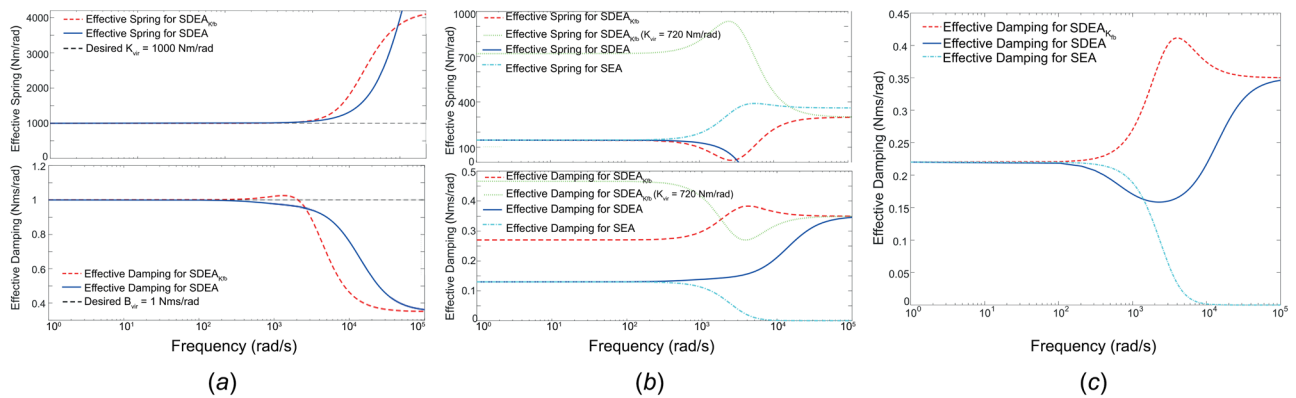
Figures 3 and 4 also demonstrate that SEA approaches its physical stiffness  $K$  at high frequencies, while the presence of even a low  $B_f$  changes the high-frequency behavior, with both SDEA and  $SDEA_{K_{fb}}$  converging to  $B_f$ .



**Fig. 3 Effect of physical filter damping  $B_f$  on performance SDEA during Voigt, spring, and null impedance rendering: (a) Voigt model rendering, (b) spring rendering, and (c) null impedance rendering**



**Fig. 4 Effect of unmeasured  $B_f$  on the performance of SDEA $_{K_b}$  during Voigt, spring, and null impedance rendering: (a) Voigt model rendering, (b) spring rendering, and (c) null impedance rendering**



**Fig. 5 Effective impedance comparisons of the S(D)EA models under VSIC: (a) effective impedances for SDEA and SDEA $_{K_b}$  during Voigt model rendering, (b) effective impedances for S(D)EA and SDEA $_{K_b}$  during spring rendering, and (c) effective damping of S(D)EA and SDEA $_{K_b}$  during null impedance**

## 5.2 Comparisons of Series Damped Elastic Actuation, SDEA $_{K_b}$ , and Series Elastic Actuation

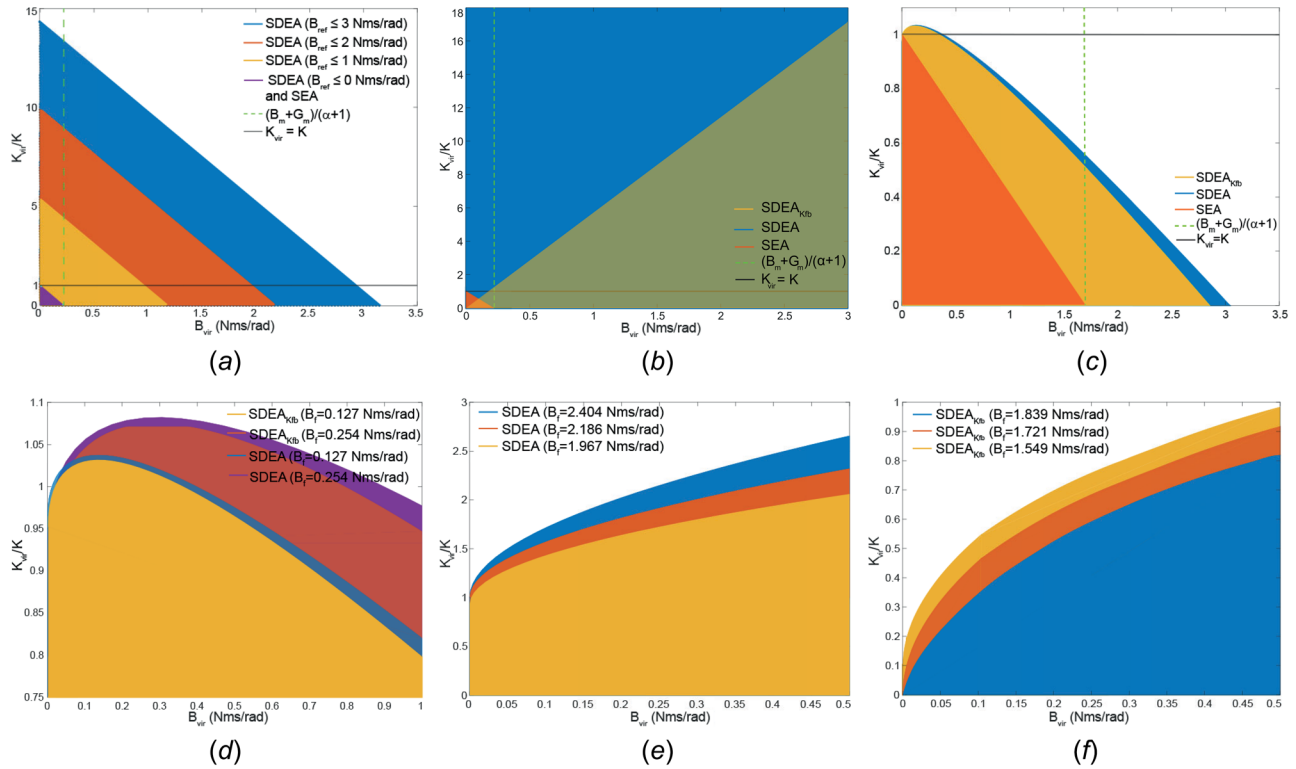
**5.2.1 Voigt Model Rendering.** The passive realizations of the S(D)EA models during Voigt model rendering are presented in Tables 1(a), 1(d), and 1(g), respectively.

Table 1(d) presents a realization featuring a controllable damping term alongside a controllable stiffness term, allowing for direct rendering of desired virtual springs and dampers. In contrast, Table 1(g) exhibits realizations with controllable stiffness terms for rendering desired virtual springs, but the damping terms include the parasitic damping  $B_f$  of the physical filter in parallel with the

controllable damping terms. Although this configuration still facilitates the rendering of desired damping levels for SDEA $_{K_b}$ , the control is more indirect and require suppression of  $B_f$ .

Figure 5(a) depicts comparisons of the Voigt model rendering performance of SDEA and SDEA $_{K_b}$  in terms of their effective impedances. SEA is not presented in Fig. 5(a) because damping augmentation is not possible for SEA. Desired damping and stiffness values are set to 1 Nms/rad and 1000 Nm/rad for Fig. 5(a), respectively. The appropriate  $B_{ref}$  are selected such that  $B_f + c_{4v} = B_{vir} = 1$  Nms/rad for SDEA $_{K_b}$  and  $\frac{\alpha}{\alpha+1} B_{ref} + \sigma(B_m + G_m) = B_{vir} = 1$  Nms/rad for SDEA. The desired damping value is almost 5 times the parasitic damping value in Table 1(f) and the desired stiffness value





**Fig. 6 K-B plots of the S(D)EA models under VSIC: (a) K-B plots for S(D)EA under VSIC for var.  $B_{ref}$ , (b) K-B plots for S(D)EA & SDEA $_{Kfb}$  for high  $B_f$ , (c) K-B plots for SDEA & SDEA $_{Kfb}$  for low  $B_f$ , (d) K-B plots for SDEA & SDEA $_{Kfb}$  for various  $B_f$ , (e) K-B plots of SDEA around the critical  $B_f$ , and (f) K-B plots of SDEA $_{Kfb}$  around the critical  $B_f$**

is almost three times the physical stiffness of the plant. The results indicate that SDEA $_{Kfb}$  has a similar performance bandwidth to SDEA in terms of achieving the desired damping level and virtual stiffness, but SDEA has slightly better accuracy, as shown in Fig. 5(a).

Figure 6(a) presents K-B plots for S(D)EA for various  $B_{ref}$  values. The lowest  $B_{ref}$  value for the K-B plots of SEA and SDEA is set to  $-\frac{B_m+G_m}{\alpha}$  according to the passivity condition. The area under the K-B plot of SDEA increases as  $B_{ref}$  increases, as shown in Fig. 6(a). During physical implementations, the unmodeled actuator saturation will introduce an upper bound to this K-B plot. The K-B plot of SEA is a subset of that of SDEA since only *damping compensation* is possible with SEA, while *damping augmentation* is also possible with SDEA. In particular, SEA can passively render virtual damping levels from zero to  $\frac{B_m+G_m}{\alpha+1}$ , while its virtual stiffness needs to simultaneously decrease from its highest value of  $K$  to zero, that is, the highest virtual damping rendering requires zero stiffness, while the highest stiffness rendering necessitates zero virtual damping.

Figure 6(b) depicts K-B plots of the S(D)EA models for relatively high values of  $B_f$ . In this figure, the K-B plot of SDEA covers the whole plot, as no actuator saturation is imposed, while the K-B plot of SEA only covers the damping compensation region as in Fig. 6(a). The K-B plot of SDEA $_{Kfb}$  is depicted by the yellow triangle, which indicates that SDEA $_{Kfb}$  can passively render high virtual stiffness levels with high virtual damping, but unlike the case for SDEA, SDEA $_{Kfb}$  cannot passively render high  $K_{vir}$  with low  $B_{vir}$ .

Figures 6(c)–6(f) depict K-B plots of the S(D)EA models using the physical filter damping  $B_f = 0.127$  Nms/rad of the experimental setup and gains presented in Sec. 6. Unlike the case in K-B plots presented in Fig. 6(b), higher  $K_{vir}$  is limited by  $B_{vir}$  for SDEA and SDEA $_{Kfb}$ , as presented in Fig. 6(c). SDEA $_{Kfb}$  and SDEA have similar performance for the lower  $B_{vir}$ , but SDEA can reach a higher  $B_{vir}$  than SDEA $_{Kfb}$  for the same  $K_{vir}$  levels.  $B_{vir}$  and  $K_{vir}$  ranges for SEA are constrained by  $\frac{B_m+G_m}{\alpha+1}$  and physical stiffness  $K$ , respectively, as in Fig. 6(b).

Figure 6(d) illustrates K-B plots for SDEA and SDEA $_{Kfb}$  for various levels of  $B_f$ . It demonstrates that as  $B_f$  decreases, the  $K_{vir}/K$

ratio approaches unity, while the peak stiffness level decreases. Moreover, the total area increases for higher  $B_f$ , for both SDEA and SDEA $_{Kfb}$ , with the peak stiffness of SDEA surpassing that of SDEA $_{Kfb}$ . Figures 6(a)–6(d), indicate that  $B_f$  directly influences the shape of K-B plots.

Figure 6(e) presents K-B plots of SDEA, as  $B_f$  approaches the critical point where the change in passivity bounds occurs. In particular, when  $B_f$  exceeds a critical value, Condition (i) of Proposition 1 in Ref. [14] becomes more conservative than Condition (v) of Proposition 1. Beyond this  $B_f$  value, the K-B plots in Figs. 6(c) and 6(e) begin to resemble the K-B plots in Fig. 6(b). Similarly, Fig. 6(f) depicts the scenario where  $B_f$  approaches its critical for SDEA $_{Kfb}$ . This limit is determined as  $\frac{B_m+G_m}{\alpha}$ , which alleviates the need to impose Condition (i) of Proposition 2. Beyond this critical value, the K-B plots in Figs. 6(c) and 6(f) begin to resemble the K-B plots Fig. 6(b).

As  $B_f$  approaches zero, the K-B plots of SDEA and SDEA $_{Kfb}$  converge to the K-B plot of SEA.

**5.2.2 Spring Rendering.** Tables 1(b), 1(e), and 1(h) present the passive physical equivalents for spring rendering with the S(D)EA models. All passive physical equivalents have a controllable stiffness term, in parallel to all other terms. The realization in Table 1(h) has the filter damping  $B_f$  as parallel parasitic dynamics to the controllable stiffness, negatively affecting the rendering performance.

Figure 5(b) depicts the spring rendering performances of the S(D)EA models in terms of effective impedances. While the effective damping of SDEA and SEA is similar, SDEA $_{Kfb}$  exhibits higher effective damping, indicating higher parasitic damping effects. Figure 5(b) demonstrates that at low frequencies, effective damping equals  $c_{2s} + B_f$ . Moreover, Fig. 5(b) illustrates that SDEA and SEA can render the desired virtual stiffness over a wider frequency range compared to SDEA $_{Kfb}$ . Specifically, SDEA $_{Kfb}$  renders a virtual stiffness level that is twice as high as the physical stiffness, depicted in Fig. 5(b). At high frequencies, the effective stiffness of SDEA

**Table 3 Parameters of the brake pedal with SDEA**

Parameter (unit)	Value	Parameter (unit)	Value
$J_m$ (kgm <sup>2</sup> )	0.0024	$B_m$ (Nms/rad)	0.0177
$K$ (Nm/rad)	121.8	$B_f$ (Nms/rad)	0.0127

converges to zero,  $SDEA_{K_{fb}}$  converges to a nonzero value, and SEA converges to its physical stiffness  $K$ , while effective damping of SDEA and  $SDEA_{K_{fb}}$  converges to  $B_f$  and SEA converges to zero.

**5.2.3 Null Impedance Rendering.** Tables 1(c), 1(f), and 1(i) present the realizations for null impedance rendering for the S(D)EA models. In the realization of SDEA and  $SDEA_{K_{fb}}$ ,  $B_f$  term is uncontrollable, but it is serially connected to the parasitic device dynamics, similar to the physical stiffness, for SDEA, while  $B_f$  is parallel to other terms for  $SDEA_{K_{fb}}$ . Figure 5(c) depicts the comparison of the S(D)EA models during null impedance rendering in terms of effective damping of the models. Effective damping of all three models approach  $\frac{B_m+G_m}{\alpha+1}$  (which is also equal to  $c_{1n} + B_f$ ) at low frequency range, but the effective damping levels of SDEA and SEA get lower at the intermediate frequency range, while the effective damping of  $SDEA_{K_{fb}}$  increases. At high frequencies, effective damping converges to  $B_f$  for SDEA and  $SDEA_{K_{fb}}$ , while SEA converges to zero.

## 6 Experimental Validations

This section presents experimental validations of the theoretical passivity bounds and haptic rendering performance. The experimental setup comprises a custom single-axis brake pedal with SDEA, based on prior designs [35,36]. The pedal is actuated by a brushless DC motor equipped with an optical encoder and Hall-effect sensors. The torque output of the motor is amplified by a transmission ratio of 1:39.5. The spring of the series elastic element is implemented as a compliant cross-flexure joint, whose deflections are measured by an encoder to estimate the interaction torques. An eddy-current damper is implemented using permanent magnets and an aluminum plate. The parameters of the SDEA plant are experimentally determined as in Table 3 [14].

Velocity-sourced impedance control controllers are implemented using a real-time operating system with a sampling rate of 1 kHz using an industrial PC connected to an EtherCAT bus. Unless specified otherwise, the controller gains of VSIC are taken as  $G_m = 0.0576$  Nms/rad and  $G_t = 30$  rad/(sNm) throughout the experiments.

The transmission ratio of the plant was set to unity while deriving the theoretical passivity bounds. The results can be extended to systems with a transmission ratio of  $n$  by applying the following transformations to form an equivalent system:  $J_{meq} = n^2 J_m$ ,  $B_{meq} = n^2 B_m$ ,  $G_{meq} = n^2 G_m$ , and  $G_{teq} = 1/n G_t$ .

**6.1 Verification of Theoretical Passivity Bounds of  $SDEA_{K_{fb}}$ .** The passivity of a system is investigated by studying the coupled stability of interactions when the system is exposed to the most destabilizing environments [37]. This approach allows for the conclusion of system's passivity if there exists no set of ideal springs or inertias that destabilizes the system. For S(D)EA, the most destabilizing environments are inertial [13,14,23,38].

To verify the passivity bounds in Sec. 4, the SDEA brake pedal underwent testing with various inertial environments. The coupled stability conclusions for each data point in these plots were made after a search in the parameter space. Four distinct masses were coupled to the end effector of the SDEA brake pedal, and impacts were imposed on the end effector to excite the system at all possible frequencies. A line search was conducted along the  $y$ -axis, starting from 25% below the most conservative theoretical passivity bound and  $K_{vir}$  was increased with a resolution of 0.5 Nm/rad. For each trial parameter set, if no violation of the coupled stability was observed

after five trials at each end-effector inertia, then it was concluded that the experimental evidence indicates passivity. Otherwise, if any violation of coupled stability was observed, then the parameter set is active. Supplementary materials including videos of experiments are available at our website<sup>3</sup>.

**6.1.1 Voigt Model Rendering.** In these experiments, we have investigated the coupled stability of  $SDEA_{K_{fb}}$  under VSIC during Voigt model rendering when the controllers are proportional. To validate the necessary and sufficient conditions provided in Proposition 2, we have tested various  $K_{ref}$  and  $B_{ref}$  values.

Figure 7(a) illustrates the experimental  $K_{vir}$ - $B_{ref}$  plot for the brake pedal with  $SDEA_{K_{fb}}$  under VSIC. The theoretical passivity bound based on Condition (i) of Proposition 2 is represented by the magenta line, while the bound based on Condition (ii) of Proposition 2 is depicted by the blue line. The experimental results confirm the theoretically predicted passivity boundary. Specifically, the experimental data conform to the dashed blue line, which is the more conservative necessary condition. The experimental results closely align with the theoretical values, with an error of about 8%. The experimental results are expected to be more conservative due to unmodeled Coulomb friction and hysteresis effects that result in additional dissipation in the physical system that cannot be accommodated in the LTI plant model.

Figure 7(b) presents the experimentally determined K-B plot of  $SDEA_{K_{fb}}$  under VSIC, where  $B_{vir}$  is computed from  $c_{4v} + B_f$  according to Table 1(g). Figure 7(b) shows that  $B_{vir}$  can be compensated with  $SDEA_{K_{fb}}$  by choosing the appropriate  $B_{ref}$ . It can also be verified that  $K_{vir}$  can exceed the physical stiffness for lower  $B_{vir}$  levels, and  $K_{vir}$  decreases as  $B_{vir}$  increases.

**6.1.2 Spring Rendering.** In these experiments, the coupled stability of  $SDEA_{K_{fb}}$  under VSIC during spring rendering when the controllers are proportional has been studied. We have selected various  $K_{ref}$  values for several  $G_t$  gains according to the conditions given in Proposition 2 when  $B_{ref} = 0$ .

Figure 7(c) presents the experimental  $K_{vir}$ - $G_t$  plot for the  $SDEA_{K_{fb}}$  brake pedal under VSIC. The theoretical passivity bounds derived from Conditions (i)-(ii) of Proposition 2 when  $B_{ref} = 0$  are depicted as the magenta and blue lines, respectively. It can be observed that the two bounds are very close to each other for the given the parameters of the brake pedal. The experimental results validate the theoretically predicted passivity boundary; the theoretical bounds are approximately 6.5% more conservative than the experimental results.

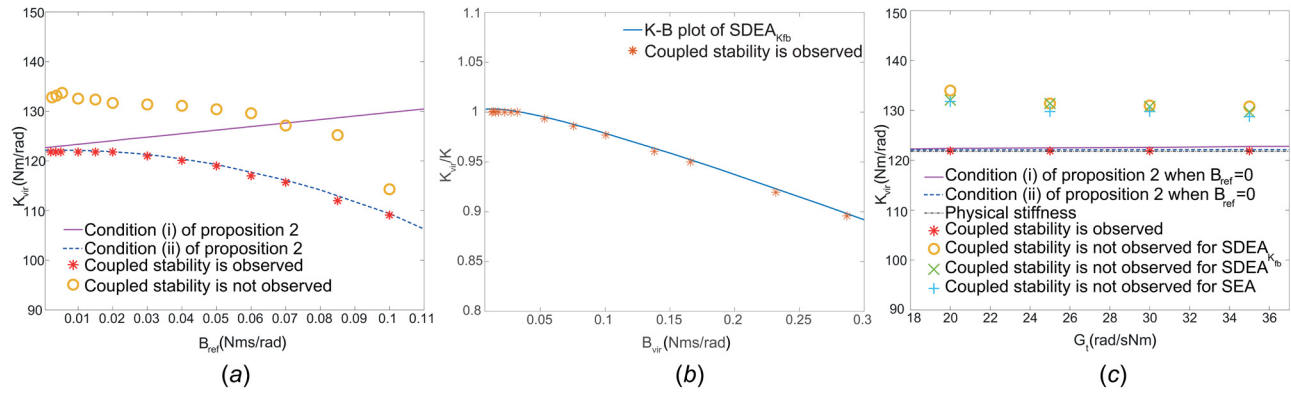
Figure 7(c) also illustrates the experimentally determined bounds for coupled stability of the S(D)EA models under VSIC. As expected from the theoretical analysis, the passivity bounds for virtual spring stiffness are the highest for  $SDEA_{K_{fb}}$ .

**6.1.3 Null Impedance Rendering.** In these experiments, the coupled stability of  $SDEA_{K_{fb}}$  under VSIC during null impedance rendering when the controllers are proportional has been studied. According to the theoretical bounds,  $SDEA_{K_{fb}}$  during null impedance rendering is passive for all  $G_t$  values as established in Proposition 2 with  $B_{ref} = K_{ref} = 0$ . Our experiments validate this prediction by displaying coupled stability for all the  $G_t$  gains tested.

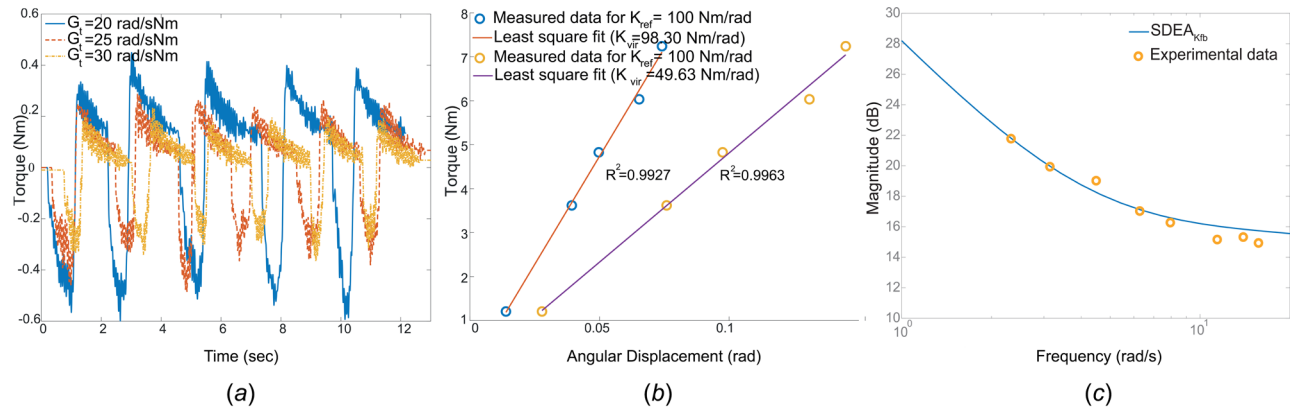
**6.2 Evaluation of Haptic Rendering Fidelity.** In this subsection, we experimentally evaluate the performance of  $SDEA_{K_{fb}}$  under VSIC during null impedance, spring, and Voigt model rendering. We also compare the haptic rendering performance of  $SDEA_{K_{fb}}$  with SDEA under VSIC.

**6.2.1 Null Impedance Rendering Performance.** The null impedance rendering is valuable as it allows the user to move the end-effector with minimal resistance. An experimental validation of the null impedance rendering performance of  $SDEA_{K_{fb}}$  under VSIC is demonstrated in Fig. 8(a), utilizing three different levels of the

<sup>3</sup>[https://hmi.sabanciuniv.edu/SDEAKfb\\_EffectofDamping\\_Supplementary.mp4](https://hmi.sabanciuniv.edu/SDEAKfb_EffectofDamping_Supplementary.mp4)



**Fig. 7** Passivity bounds versus experimental coupled stability results: (a)  $B_{ref}$ – $K_{vir}$  plot for SDEA $_{K_{fb}}$  during Voigt model rendering, (b) K-B plot of SDEA $_{K_{fb}}$ , and (c)  $G_t$ – $K_{vir}$  plot for SDEA $_{K_{fb}}$  during spring rendering



**Fig. 8** (a) Null impedance rendering performance of SDEA $_{K_{fb}}$  for  $G_t = 20, 25$ , and  $30$  rad/(sNm), (b) virtual stiffness rendering performance of SDEA $_{K_{fb}}$  for  $K_{ref} = 50$  and  $100$  Nm/rad, and (c) Voigt model rendering performance of SDEA $_{K_{fb}}$  when  $B_{vir}$  and  $K_{vir}$  were set as  $0.83$  Nms/rad and  $25$  Nm/rad, respectively

torque controller gain  $G_t$ . As the value of the torque controller gain  $G_t$  is increased from  $20$  rad/(sNm) to  $30$  rad/(sNm), the force required to move the brake pedal reduces from  $1.49\%$  to  $0.91\%$  of the  $40$  Nm torque output capability of the brake pedal. The experimental results in Fig. 8(a) align well with the analysis in Sec. 5, which demonstrates the positive influence of increasing torque controller gain  $G_t$  on null impedance rendering performance.

**6.2.2 Spring Rendering Performance.** The implementation of virtual constraints to restrict users from accessing undesired regions of the workspace is commonly achieved through spring rendering, making it a crucial control mode. The experimental verification of the spring rendering performance for two levels of desired virtual stiffness, with  $K_{ref}$  values of  $50$  Nm/rad and  $100$  Nm/rad, is presented in Fig. 8(b). To determine the rendered stiffness of the SDEA $_{K_{fb}}$  under VSIC, predetermined torques are applied to the end effector, and the resulting deflections are measured. A least-squares fit to the experimental data reveals that the values of  $K_{vir}$  are  $49.63$  Nm/rad with  $R^2$  of  $0.99$  and  $98.30$  Nm/rad with  $R^2$  of  $0.99$  with  $0.73\%$  and  $0.25\%$  the observed errors, respectively.

These experiments are also conducted with  $G_t = 25$  rad/(sNm). The performance of spring rendering is evaluated for two levels of virtual stiffness:  $K_{ref} = 50$  Nm/rad and  $100$  Nm/rad. A least-square fit to the experimental data reveals that for  $K_{ref} = 50$  Nm/rad, the value of  $K_{vir}$  is  $49.86$  Nm/rad with  $R^2 = 0.99$ , resulting in an error of  $1.48\%$ . For  $K_{ref} = 100$  Nm/rad, the value of  $K_{vir}$  is  $97.88$  Nm/rad with  $R^2 = 0.99$ , resulting in an error of  $0.39\%$ . These findings align with the analysis presented in Sec. 5, indicating the beneficial effect of increasing the torque controller gain  $G_t$  and desired virtual stiffness  $K_{ref}$  on the spring rendering performance.

Experiments were carried out to evaluate the system's ability to produce the proper virtual spring forces when the end effector is excited by dynamic human inputs during spring rendering using SDEA and SDEA $_{K_{fb}}$ . Both systems were tested using a  $K_{ref}$  value of  $50$  Nm/rad, and the normalized root-mean-square error (NRMSE) values were recorded as  $3.97\%$  for SDEA and  $5.09\%$  for SDEA $_{K_{fb}}$ , respectively. The experimental results align with the theoretical analysis, which predicts that the NRMSE of SDEA $_{K_{fb}}$  would be higher than that of SDEA due to the unmeasured  $B_f$ .

**6.2.3 Voigt Model Rendering Performance.** Figure 8(c) presents experimental verification of the Voigt model rendering performance of SDEA $_{K_{fb}}$ . In this experiment, the end effector of the brake pedal was excited by an ideal velocity source imposing sine waves at eight distinct frequencies, ranging from  $1$  rad/s to  $16$  rad/s, while the brake pedal was rendering a Voigt model. The resulting interaction forces were recorded and presented as the Bode plot in Fig. 8(c). The average error between the experimental data and theoretical predictions in this Bode plot was computed as  $4.5\%$ , indicating high fidelity Voigt models rendering as predicted by theoretical results.

Experimental validation of the interaction performance of the SDEA $_{K_{fb}}$  brake pedal, rendering a Voigt model with  $K_{ref} = 100$  Nm/rad and  $B_{ref} = 0.01$  Nms/rad, under dynamic user inputs, was also conducted. Desired interaction torques derived from the Voigt model and the corresponding interaction torques measured using the series damped elastic element were compared. The resulting NRMSE of  $1.1\%$  confirms the accuracy of the rendering. A similar experiment was carried out for the SDEA brake pedal, yielding an NRMSE of  $1\%$ . As anticipated from the theoretical analysis, the higher NRMSE of SDEA $_{K_{fb}}$ , compared to SDEA is attributed to the



parasitic damping effect of the filter and the absence of damping on the computed force in the series elastic element.

## 7 Conclusion

We have established a fundamental rendering limitation for SEA when the inherent damping of the series elastic element is considered in the system model. We have proven that the viscous damping of the physical filter enables passive rendering of stiffness levels that are higher than the physical stiffness of the filter, as well as rendering Voigt models whose damping levels exceed that of the physical damping of the plant. This result has practical consequences, as SEA with an ideal series spring is an idealization, and every physical spring implementation possesses some level of dissipation. Furthermore, this result can guide practical designs of systems with SEA, as intentional utilization of larger physical damping levels in the series elastic element can extend the range of passively renderable virtual environments.

We have studied a widely used controller for SEA, where the forces induced on the series damping element are neglected in the feedback. Note that this model, called  $SDEA_{K_{fb}}$  under VSIC, corresponds to the most common SEA controllers; however, our analysis considers the inherent damping of the series elastic element in the plant model. We have established necessary and sufficient conditions for the passivity of  $SDEA_{K_{fb}}$  under VSIC during Voigt model, linear spring, and null impedance rendering. We have compared these results with the passivity conditions for S(D)EA under VSIC, to establish the effect of unmeasured series physical damping on the results. Our results provide the proper passivity bounds for a realistic model of SEA under VSIC.

We have derived passive physical equivalents of  $SDEA_{K_{fb}}$  under VSIC while rendering Voigt, linear springs, and the null impedance models. The passive physical equivalents not only make the rendered impedance and parasitic dynamics of the closed-loop system explicit but also enable the objective comparisons among the S(D)EA models in terms of their rendering performance. Through the passive physical equivalents, we have rigorously compared the effects of measuring and omitting the damping force of the series elastic element on the closed-loop rendering performance of the system. We have shown that if the force acting on  $B_f$  is used in feedback control (SDEA under VSIC), this additional information has beneficial effects on haptic rendering performance, enabling direct control of damping during passive Voigt model renderings. On the other hand, if force on  $B_f$  is not used in feedback control ( $SDEA_{K_{fb}}$  under VSIC), then  $B_f$  acts as an additive effect in addition to more indirectly controllable effective damping terms that can still be adjusted to passively render Voigt models.

We have also compared SDEA and  $SDEA_{K_{fb}}$  under VSIC in terms of their effective impedances and shown that parasitic damping effects are higher for  $SDEA_{K_{fb}}$ . The maximum passively renderable virtual stiffness is limited by the physical stiffness of the filter for SDEA under VSIC [14,15,22], while [15] has proven that the unmeasured damping effect on the physical filter of  $SDEA_{K_{fb}}$  can help exceed this physical spring stiffness bound. We have extended this analysis to Voigt model rendering without the assumption of positive controller gains, and studied the performance tradeoff through realizations.

Table 1 illustrates the realizations of the S(D)EA models under VSIC while rendering null impedance, springs, and Voigt models. It is important to note the continuity among these realizations. For example, the null impedance realization for SEA under VSIC can be obtained from the Voigt model realization for SDEA or  $SDEA_{K_{fb}}$  by setting  $B_{ref}$ ,  $K_{ref}$ , and  $B_f$  to zero.

Throughout the paper, proportional controllers are preferred to simplify the analysis while capturing the main effects of interest. Introducing more complex controllers significantly complicates realizations, making their interpretation much harder.

While the passivity conditions and physical realizations are informative, they also impose relatively conservative bounds to ensure the coupled stability of interactions. More relaxed coupled stability conditions can be established using less conservative analysis

techniques, such as time domain passivity [39] or complementary stability [40–42]. However, these techniques depend on numerical calculations or optimizations; hence, they cannot offer closed-form analytical solutions and a general understanding of the tradeoffs.

Our future work involves extending the passive realizations to fractional order control systems [43,44], as interpretations of such controllers can significantly benefit from physical intuition.

## Acknowledgment

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## Funding Data

- TUBITAK (Grant Nos. 216M200 and 23AG003; Funder ID: 10.13039/501100004410).

## Data Availability Statement

The datasets generated and supporting the findings of this article are obtainable from the corresponding author upon reasonable request.

## Appendix A: Proof of Proposition 1

The proof of Proposition 1 can be presented as follows:

Since both the low- and high-frequency dynamics have a pole at the origin,  $Z_{SDEA_{cl}}$  possesses a pole at the origin and with an equal number of remaining poles and zeros. The passivity of  $Z_{SDEA_{cl}}$  implies that it can have no right-hand plane (RHP) poles or zeros. Furthermore, any complex imaginary poles are simple with positive real residues.

The following complex integral is considered  $\oint [Z_{SDEA_{cl}}(s) - B_f] ds$  along the directed closed contour depicted in Fig. 9. Given that the contour includes infinitesimal indentations around the pole at the origin and possible (representative) simple complex poles on the imaginary axis, Cauchy's residue theorem ensures that the complex integral evaluates to zero, since the integrand is analytic everywhere along and inside the contour.

The contour can be divided into six parts: The semi-infinite negative (*I*) and positive (*II*) imaginary axes (depicted in black), the outer semicircle (*III*) at infinity (depicted in blue), and the infinitesimal semicircular indentation (*IV*) around the origin and representative infinitesimal semicircular indentations (*V–VI*) around a pair of simple complex conjugate poles (depicted in red).

Since  $Z_{SDEA_{cl}}$  models the impedance of a physical system, the real part of this complex function is even, while the imaginary part is odd. Accordingly, for the complex the integrals taken along the negative

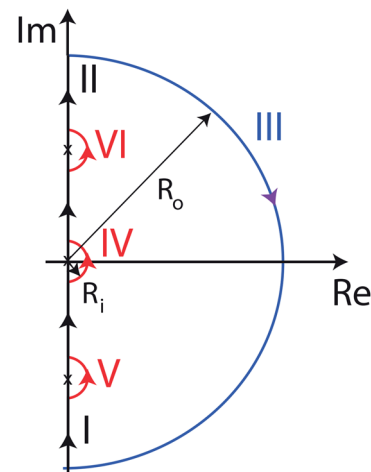


Fig. 9 The directed closed contour (with indentations around the simple poles) used for the complex integral



(I) and positive (II) axes, the imaginary parts cancel out, while the real parts are added to result in  $\int_{\text{I+II}} [Z_{\text{SDEA-cl}}(s) - B_f] ds = 2i \int_0^\infty [\Re\{Z_{\text{SDEA-cl}}(\omega)\} - B_f] d\omega$ .

Since  $Z_{\text{SDEA-cl}}(s \rightarrow \infty) = \frac{K}{s} + B_f$ , the integral taken along the outer semicircle at infinity evaluates to  $\lim_{R_o \rightarrow \infty} \int_{\text{III}} \frac{K}{s} ds = \lim_{R_o \rightarrow \infty} \int_{\pi/2}^{-\pi/2} \frac{KiR_o e^{i\theta}}{R_o e^{i\theta}} d\theta = -i\pi K$ .

Similarly, since  $Z_{\text{SDEA-cl}}(s \rightarrow 0) = \frac{K_{\text{vir}}}{s} + B_{\text{vir}}$ , the integral taken on the infinitesimal semicircle around zero evaluates to  $\lim_{R_i \rightarrow 0} \int_{\text{IV}} \left[ \frac{K_{\text{vir}}}{s} + (B_{\text{vir}} - B_f) \right] ds = \lim_{R_i \rightarrow 0} \int_{-\pi/2}^{\pi/2} \left[ \frac{K_{\text{vir}}}{R_i e^{i\theta}} + (B_{\text{vir}} - B_f) \right] iR_i e^{i\theta} d\theta = i\pi K_{\text{vir}}$ .

Finally, if simple complex poles exist, the infinitesimal semicircular indentations (V – VI) around these complex poles will contribute to the integral with a positive value  $i\pi 2\Delta \geq 0$ , due to their (equal) positive real residues  $\Delta$ . Collecting all terms together and rearranging, one can express

$$\pi(K_{\text{vir}} - K) + 2\pi\Delta + 2 \int_0^\infty [\Re\{Z_{\text{SDEA-cl}}(\omega)\} - B_f] d\omega = 0 \quad (\text{A1})$$

The available bandwidth  $\Omega_a$  of any causal system is defined as the finite frequency after which the closed-loop system converges to its plant dynamics [45]. In particular,  $\Re\{Z_{\text{SDEA-cl}}(\omega)\}$  can be made arbitrarily close to  $B_f$  for  $\omega > \Omega_a$  for all plants with SDEA for a sufficiently large (but finite)  $\Omega_a$ .

Similarly, the performance bandwidth  $\omega_p$  is defined as the frequency at which the virtual environment can be rendered as  $Z_{\text{SDEA-cl}}(s) = \frac{K_{\text{vir}}}{s} + B_{\text{vir}}$ . There exists a sufficiently small  $\omega_p < \Omega_a$  for which the approximation error converges to zero.

Hence, the bounds in the definite integral of Eq. (A1) can be replaced with  $\omega_p$  and  $\Omega_a$ . Re-arranging, one can express

$$\int_{\omega_p}^{\Omega_a} \Re\{Z_{\text{SDEA-cl}}(\omega)\} d\omega + \pi\Delta = \frac{\pi}{2}(K - K_{\text{vir}}) + \Omega_a B_f - \omega_p B_{\text{vir}} \quad (\text{A2})$$

Passivity implies  $\Re\{Z_{\text{SDEA-cl}}(\omega)\} \geq 0$  for all  $\omega$  and  $\Delta \geq 0$ . Consequently, Eq. (A2) indicates that

$$\frac{\pi}{2} K + \Omega_a B_f \geq \frac{\pi}{2} K_{\text{vir}} + \omega_p B_{\text{vir}}$$

concluding the proof for all causal LTI controllers.

The result can be generalized to other causal controllers, including nonlinear and time-varying ones [46], since for an LTI plant, LTI controllers can achieve optimal performance. Consequently, for an LTI plant, optimal performance obtained by a causal stabilizing LTI controller determines the performance limits that cannot be improved by any other causal stabilizing controller. Accordingly, as long as the SDEA plant is LTI, the rendering limitation expressed in Eq. (A2) is valid for all casual controllers.

## Appendix B: Proof of Proposition 2

The proof of Proposition 2 can be presented as follows: According to the positive realness theorem [26,47],

- $Z(s)$  has no poles in the right half plane: If Routh–Hurwitz stability criterion is applied, there is no pole in the open right half plane for  $Z_{\text{voigt}}^{\text{SDEA-}K_{\text{rb}}}(s)$ , since  $(\alpha + 1)$  is non-negative, and the inner motion control loop is asymptotically stable.
- Any poles of  $Z(s)$  on the imaginary axis are simple with positive and real residues: There exist no poles on the imaginary axis, except at  $s = 0$ , as long as  $(\alpha + 1)$  is non-negative. For the pole at  $s = 0$ , the residue equals to  $\frac{\alpha}{\alpha+1} K_{\text{ref}}$ , which should be positive. If  $(\alpha + 1) = 0$ , then the impedance transfer function has double poles at  $s = 0$ , and the positive realness theorem is violated. Therefore,  $(\alpha + 1) = 0$  violates passivity.

- $\Re[Z(jw)] \geq 0$  for all  $w$ : The sign of  $\Re[Z_{\text{voigt}}^{\text{SDEA-}K_{\text{rb}}}(jw)]$  can be checked by that of  $H(jw) = d_6 w^6 + d_4 w^4 + d_2 w^2$  with

$$d_2 = K(K(B_m + G_m + B_{\text{ref}}\alpha) + B_f K_{\text{ref}}\alpha)(\alpha + 1) - K K_{\text{ref}}\alpha(B_f + B_m + G_m) \quad (\text{B1})$$

$$d_4 = B_f(B_f + B_m + G_m)(B_m + G_m + B_{\text{ref}}\alpha) - J_m\alpha(K(B_{\text{ref}} + B_f) + B_f K_{\text{ref}}) \quad (\text{B2})$$

$$d_6 = B_f J_m^2 \quad (\text{B3})$$

Here,  $d_6$  is positive, since  $B_f$  is positive. Condition (i) of Proposition 2 is imposed by the non-negative  $d_2$ . The last necessary and sufficient condition can be derived from  $d_4 \geq -2\sqrt{d_2 d_6}$  [23] as Condition (vi) of Proposition 2, which is never satisfied when  $B_{\text{ref}} = -\frac{B_m + G_m}{\alpha}$ .

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