

## Electro-optic detection of subwavelength terahertz spot sizes in the near field of a metal tip

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We report on a method to obtain a subwavelength resolution in terahertz time-domain imaging. In our method, a sharp copper tip is used to locally distort and concentrate the THz electric field. The distorted electric field, present mainly in the near field of the tip, is electro-optically measured in an (100) oriented GaP crystal. By raster scanning the tip along the surface of the crystal, we find the smallest THz spot size of 18  $\mu\text{m}$  for frequencies from 0.1 to 2.5 THz. For our peak frequency of 0.15 THz, this corresponds to a resolution of  $\lambda/110$ . Our setup has the potential to reach a resolution down to a few  $\mu\text{m}$ . © 2002 American Institute of Physics. [DOI: 10.1063/1.1503404]

In recent years, terahertz time-domain spectroscopy (TDS) has become a popular method to obtain information on the far-infrared properties of molecules, semiconductors, and other materials,<sup>1-4</sup> and has shown to be a promising technique to image objects.<sup>5-7</sup> Unfortunately, diffraction, which limits the resolution to about a wavelength, is an obstacle for a number of interesting new THz applications. Breaking the diffraction limit would, for instance, make it possible to perform THz spectroscopy on extremely small volumes, and, ultimately on living cells. Work on the label free probing of deoxyribonucleic acid (DNA),<sup>8</sup> shows that THz radiation can be used to distinguish between different biomolecules.

One way to overcome the diffraction limit is to limit the detection area with an aperture. For example, Mitrofanov *et al.*<sup>9</sup> have shown that by using a detector with an integrated aperture, near field images can be made with a spatial resolution determined by the aperture size, and not by the THz wavelength. Unfortunately, this method relies on extensive micro-fabrication and requires two parallel plates, namely one carrying the object and the other the detector, to be scanned along each other at close distance, which is difficult to do in practice. Also, since the aperture has a finite thickness, waveguide effects strongly attenuate long-wavelength components,<sup>10</sup> which makes their detection difficult. Another approach to obtain subwavelength resolution<sup>11</sup> uses a laser to modulate the THz transmission of a small area on a semiconductor surface near the object. The modulation on the THz beam is measured in the far field, and an improved spatial resolution, down to 50  $\mu\text{m}$  is observed. In this method, however, signals are small and superimposed on a large background, complicating the detection.

Here, we present a method to create and detect a local THz source with a subwavelength resolution, which has the potential to image a microscopic object by raster scanning it between the source and the detector. In our method, THz pulses are incident on a sharp metal tip, having subwavelength dimensions, which is held in close proximity to an (100) oriented GaP crystal. A synchronized probe pulse is

tightly focused into the crystal close to the tip and, using the electro-optic effect in the same crystal, *only* detects the THz polarization component that is *perpendicular* to the (100) surface. As the incident THz polarization is parallel to the crystal surface, the perpendicular component is present *only* near the tip, which functions as a lightning rod where field lines are bent and concentrated. We find that, directly under the tip, the measured THz spot size has a diameter of only 18  $\mu\text{m}$ , corresponding to  $\lambda/110$  for a wavelength of 2 mm. Our method to obtain a subwavelength resolution has four advantages: (i) the resolution is determined by the spot diameter of our probe laser, not by the incident THz beam size, (ii) our measurements are background free, as we only detect the component near the tip that is perpendicularly polarized to the crystal surface, (iii) our detection setup measures in the near field, where the electric field can be strong. This provides sufficiently large signals and thus reduces the measurement time, and (iv) our method has the more practical requirement of scanning a tip along a surface, instead of two surfaces parallel to each other.

Our setup to generate THz pulses, is described elsewhere.<sup>12</sup> As we show in Fig. 1, the THz pulses, which have an average power of 40  $\mu\text{W}$ , are focused onto a copper tip, which has been sharpened in a FeCl etch bath. A 300  $\mu\text{m}$  thick, (100) oriented GaP detection crystal is placed within a distance of several microns of the tip. The crystal surface closest to the tip is coated with a 300 nm thick Ge layer on top off a 133 nm thick SiO<sub>2</sub> layer. A synchronized, 800 nm probe pulse, derived from the same laser that generates the THz pulse, is focused to about a 20  $\mu\text{m}$  spot size on the coating, initially counter propagating with respect to the THz pulse. The coating, which has a negligible effect on the THz beam, absorbs any probe light before it reaches the tip. This eliminates any measurement errors that would have been caused by probe light scattered off the tip. After the reflection at the coating, the probe pulse is sent through a  $\lambda/4$  plate, and a Wollaston prism to a differential detector. When a THz electric field is present in the crystal, the electro-optic effect changes the probe polarization from linear to elliptical. The amount of ellipticity is proportional to the instantaneous THz electric field, and is measured by the differential detec-

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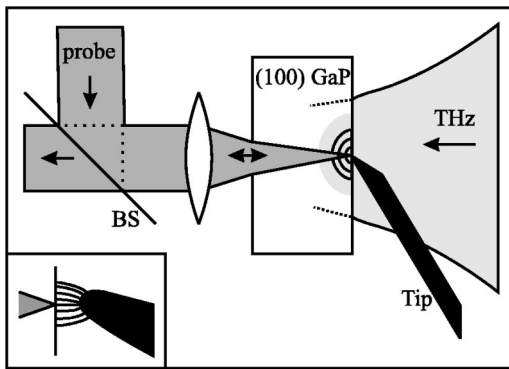


FIG. 1. Schematic picture of the setup used. The THz radiation comes in from the right-hand side, and scatters at the tip. The probe beam comes in from the left-hand side, and is focused on the crystal–air interface of an (100) GaP crystal. At this interface, the probe beam polarization is changed by the THz electric field near the tip. The probe beam reflects back at the crystal interface, and is sent to the differential detection setup (not shown). The inset on the bottom left side zooms in on the tip, and schematically shows the near-field electric-field lines near the tip, which have a component perpendicular to the crystal surface.

tor setup. We emphasize that in an (110) oriented crystal, which is the conventionally used orientation in electro-optic detection, the probe polarization is only affected by electric fields parallel to the crystal surface.<sup>13</sup> Here, however, using an (100) oriented crystal, the probe polarization is *only* changed by electric fields perpendicular to the crystal surface<sup>14</sup> making the probe “blind” to the parallel polarized electric field of the incident THz pulse. Near the tip, however, as is illustrated in the inset of Fig. 1, the THz near field must have a component perpendicular to the crystal surface, since the electric field near a metal, in this case the tip, has to be perpendicular to the metal surface.

The inset of Fig. 2 shows the THz electric field as a function of the delay between the probe and THz pulse, measured directly under the copper tip, which is pressed against the crystal. The power spectrum of this pulse is plotted in Fig. 2. It contains frequencies from 0.1 to 2.5 THz and peaks around 0.15 THz. Although this measurement was obtained in less than a minute, we note that, in our setup, a measurement of a 25 ps long THz electric-field pulse is obtained in 20 ms with a dynamic range of 50.

Figure 3 shows the measured peak THz electric field as a

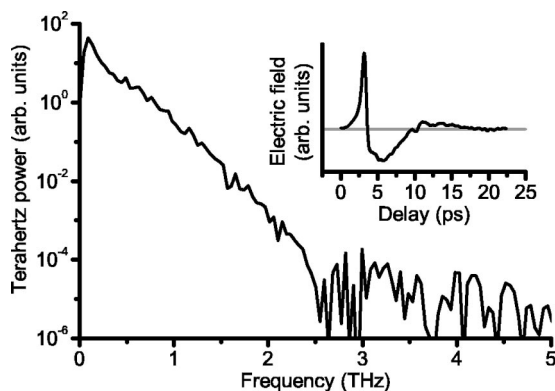


FIG. 2. THz power spectrum, calculated from the THz electric field as a function of delay shown in the inset. The measurement was performed with the tip pressed against the crystal, directly above the probe beam. This measurement was obtained in less than a minute.

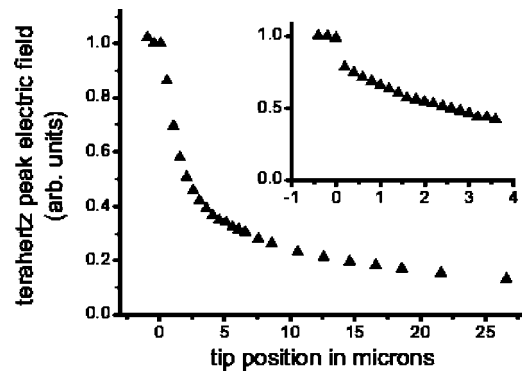


FIG. 3. Measured peak THz electric field as a function of the tip–crystal distance. The inset shows a separate measurement of the region close to the crystal.

function of the tip–crystal separation. The measurement covers a total distance of about 30  $\mu\text{m}$ . The inset of Fig. 3 shows an additional measurement, which zooms in on the region close to the crystal. We have chosen the 0  $\mu\text{m}$  position arbitrarily as the position where the signal reaches its maximum. Figure 3 shows that the THz electric field decreases to 30% of its maximum value, when the tip–crystal distance increases to 6.5  $\mu\text{m}$ , which clearly indicates near-field behavior. When the tip is close to the crystal ( $\sim 1 \mu\text{m}$ ), it is abruptly attracted to the crystal, presumably by electrostatic forces, which gives rise to a sudden increase in the signal at 0  $\mu\text{m}$ . The sharp, but smooth, decrease of the THz electric field when the distance to the crystal is increased only a few  $\mu\text{m}$ , is related to results obtained by Knoll *et al.*<sup>15</sup> at the much shorter wavelength of 10  $\mu\text{m}$ . They found that the scattering cross section of a tip close to a dielectric half space is strongly increased due to an image charge induced in the dielectric medium. It is likely, that these results also apply in our case, and that an increase of the scattering cross section is related to the strong enhancement of the local THz electric field when the tip approaches the crystal.

In Figs. 4(a)–4(d), we plot the integrated power of the THz pulses, measured while raster scanning the tip across the crystal surface, for four different tip–crystal separations. The four measurements clearly show a rapid increase in the THz spot size for increasing tip–crystal distances. This is quantitatively expressed by the results in Fig. 4(e), where we plot the measured THz spot size, defined as the full width at half maximum (FWHM) of the intensity, as a function of the tip–crystal distance. The 0  $\mu\text{m}$  point is chosen arbitrarily as the position where the spot size reaches its minimum. Figure 4(e) clearly shows that the spot diameter increases more than a factor of 6 when the tip–crystal separation increases to 30  $\mu\text{m}$ . This again strongly suggests that, close to the crystal, we are measuring in the near field of the tip. Figure 4 (e) also shows that subwavelength resolution is easily obtained even for relatively large tip–crystal distances. The smallest observed spot diameter is 18  $\mu\text{m}$ . A spectral analysis (not shown here) clearly proves that the spot diameters are independent of frequency, with the smallest spot corresponding to a maximum resolution of  $\lambda/110$  at 0.15 THz.

Earlier measurements with a probe beam focal spot size of 50  $\mu\text{m}$  gave the smallest THz spot size of 50  $\mu\text{m}$ . Our presently measured smallest spot size is 18  $\mu\text{m}$ , again similar

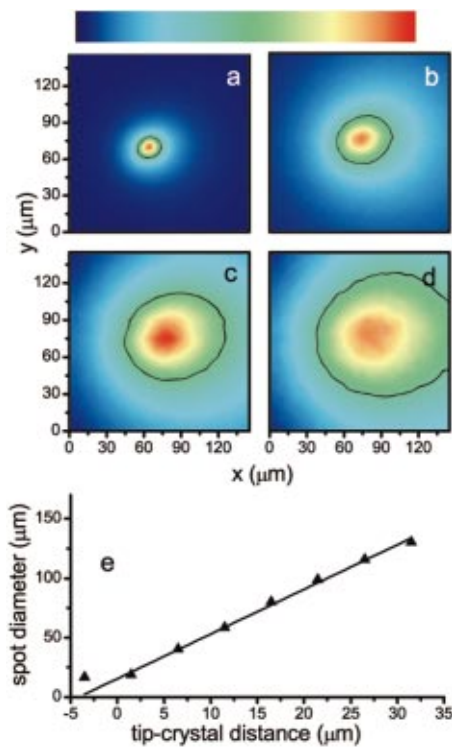


FIG. 4. (Color) THz intensity as a function of the  $xy$  position of the tip along the crystal. The scans (a) and (b) are  $50 \times 50$  pixel images with a  $3 \times 3 \mu\text{m}^2$  pixel size, and (c) and (d) are  $30 \times 30$  pixel images with a  $5 \times 5 \mu\text{m}^2$  pixel size. In the pictures (a), (b), (c), and (d), the tip is positioned at 0, 10, 20, and  $30 \mu\text{m}$  distance from the crystal, respectively. The color scale is shown at the top. The color levels run from zero (navy blue) to maximum intensity (red), and are separately normalized for each scan. The black line is the FWHM of the spot. Graph (e) shows the measured FWHM spot diameter as a function of tip–crystal distance. The solid line is a linear fit through all points but the first.

to the probe beam focal diameter of  $20 \mu\text{m}$ , suggesting that the observed THz spot sizes are determined by the focal diameter of the probe. This indicates that a resolution better than  $18 \mu\text{m}$  is achievable in our setup if we focus our probe tighter. Ultimately, we envisage an experiment in which a thin sample is placed on the crystal underneath the tip. When the crystal is raster scanned underneath the tip, a two-dimensional THz image of the sample is obtained. A high resolution, however, is only possible, if the physical size of the tip is smaller than the desired resolution. Ultimately, we expect that the resolution of our setup will reach about  $1 \mu\text{m}$ , limited by the wavelength of our probe beam.

The advantages of our method to use THz TDS to produce near-field images are that it is coherent, intrinsically broadband, and uses only  $40 \mu\text{W}$  of input power. Coherent

detection offers the advantage of measuring both amplitude and phase, while broadband detection allows us to obtain information on a wide range of frequencies in one measurement. The chance that a sample is heated significantly or even damaged is negligible in our setup, since we use low input powers. We note that the time-dependent nature of THz pulses makes it possible to use our setup for visible pump-THz probe measurements on microscopic systems.

In conclusion, we have presented a method to obtain a subwavelength resolution in THz imaging by electro-optically measuring the THz electric field close to a sharp copper tip in a (100) oriented GaP crystal. We demonstrate a smallest THz spot size of  $18 \mu\text{m}$ , determined by the probe beam diameter, and argue that our method should be able to reach a resolution of  $1 \mu\text{m}$ , limited by the focal diameter of the probe beam.

*Note added in proof.* After submission of this manuscript, we have further improved our setup. Our presently measured smallest THz spot size is now  $8 \mu\text{m}$ .

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