

APPLICATIONS OF PRECONDITIONED NEWTON-KRYLOV METHODS

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- 1 Motivation
- 2 Numerical Tools
- 3 Details of the Analysis
- 4 Results and Discussion
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Role of CFD

- Analysis
- Design (Direct and Inverse Approach)
- Optimization

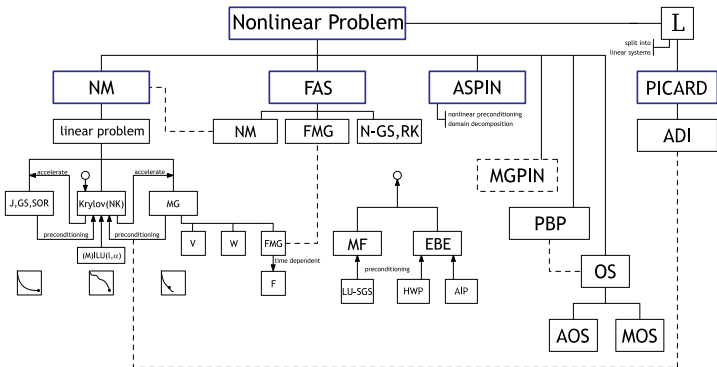


Nonlinear Equations

- Discretization yields a system of nonlinear equations
- Nonlinear-GS (or Jacobi)
- Newton's Method (Inexact NM)
- Full Approximation Scheme (Multigrid)
- Picard - (governing equations one-by-one)



Nonlinear Solver



Newton's Method

Used to linearize the non-linear system of equations.

1 $J(x^k)\Delta x^{k+1} = -f(x^k)$

2 $x^{k+1} = x^k + \Delta x^{k+1}$

3 $J_{ij} = \frac{\partial f_i}{\partial x_j}$

- Quadratically convergent from good starting guesses.
- Not globally convergent
- Requires the solution of the linear steps



Krylov Methods and Preconditioners

Following linear Krylov sub-space solvers are used.

- BiCGSTAB (Bi-conjugate Gradient Stabilized)
- GMRES (Generalized Minimal Residual)

Combined with the preconditioners

- Jacobi
- SGS (Symmetric Gauss-Seidel)
- ILU(l) (Incomplete LU decomposition)



Matrix-free algorithms

Using directional differencing the Jacobian matrix vector multiplications are carried out using only the f vector.

$$Jv = \frac{f(x + \epsilon v) - f(x)}{\epsilon} \quad (1)$$

where

$$\epsilon = \sigma^{1/2} / \|x\| [8] \quad (2)$$

if $\|x\| = 0$, the result of the matrix vector product is set identically to zero.



Matrix-free algorithms

Advantages:

- Low memory requirement
- Faster matrix vector multiplications
- Suitable for using storage schemes and some preconditioners such as Jacobi, SGS.

Disadvantages:

- Preconditioners requiring J explicitly can not be used (ILU).



Compressed Storage Schemes

CCS (Compressed Column Storage):

- Value vector: Stores the subsequent non-zeros of the matrix rows.
- Row indicator: Stores the row indexes of the elements in the value vector.
- Column pointer: Stores the locations in the value vector that start a column.



Grid Generation - 2D

The Winslow Equations

$$g_{22} \frac{\partial^2 x}{\partial \xi^2} - 2g_{12} \frac{\partial^2 x}{\partial \xi \partial \eta} + g_{11} \frac{\partial^2 x}{\partial \eta^2} = -g \left(P \frac{\partial x}{\partial \xi} + Q \frac{\partial x}{\partial \eta} \right) \quad (3)$$

$$g_{22} \frac{\partial^2 y}{\partial \xi^2} - 2g_{12} \frac{\partial^2 y}{\partial \xi \partial \eta} + g_{11} \frac{\partial^2 y}{\partial \eta^2} = -g \left(P \frac{\partial y}{\partial \xi} + Q \frac{\partial y}{\partial \eta} \right) \quad (4)$$

where $g = g_{11}g_{22} - g_{12}g_{21}$ and $P(\xi, \eta)$, $Q(\xi, \eta)$ are suitably selected control functions.

A set of possible control functions was proposed by Thompson, Thames, and Mastin (The TTM Method):

$$P(\xi, \eta) = - \sum_{n=1}^N a_n \frac{(\xi - \xi_n)}{|\xi - \xi_n|} e^{-c_n |\xi - \xi_n|} - \sum_{i=1}^I b_i \frac{(\xi - \xi_i)}{|\xi - \xi_i|} e^{-d_i [(\xi - \xi_i)^2 + (\eta - \eta_i)^2]^{\frac{1}{2}}} \quad (5)$$

$$Q(\xi, \eta) = - \sum_{n=1}^N a_n \frac{(\eta - \eta_n)}{|\eta - \eta_n|} e^{-c_n |\eta - \eta_n|} - \sum_{i=1}^I b_i \frac{(\eta - \eta_i)}{|\eta - \eta_i|} e^{-d_i [(\xi - \xi_i)^2 + (\eta - \eta_i)^2]^{\frac{1}{2}}} \quad (6)$$



Domain Decomposition Methods

The idea: Divide the problem into regions and solve them separately instead of dealing with it as a whole.

Advantages:

- Geometrical simplicity
- Application of different modeling equations For example; Navier Stoke's at the objects proximity and Euler at the other regions
- Gain in computation speed with parallel processing

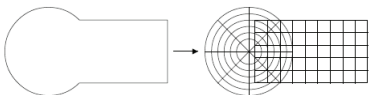


Figure: Splitting of the domain for geometrical simplicity

External Flow: Flow Past an Airfoil

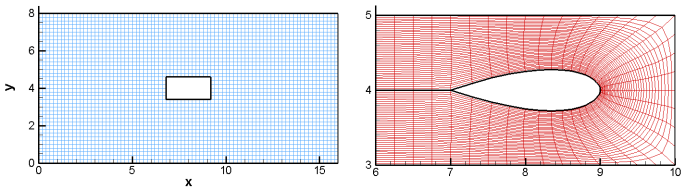


Figure: Separate domains

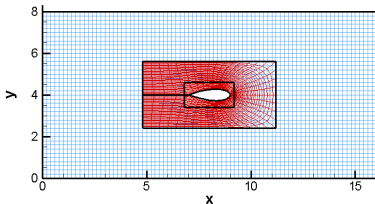


Figure: Overlapping of the domains

Overlapping Domains

- From the cartesian domain to the inner by simple interpolation
- From the inner domain to the cartesian by bilinear interpolation
- Repeat the process until the convergence is achieved

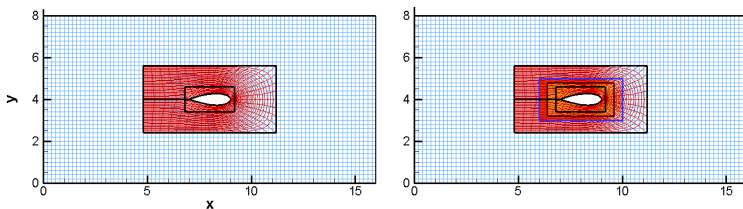


Figure: Adaptive grid

Advantages

- No grid generation in the cartesian domain
- Dealing with complex geometries is easier
- Solving the equations for the cartesian and inner domain separately (parallel computing)
- Calculation of the flow at different angle of attacks

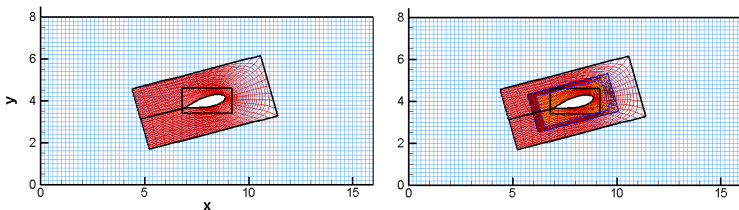


Figure: Different angle of attack

Internal Flow: Flow Between Turbomachinery Blades

- Shock waves
- Turbulent boundary layers and wakes
- Complex geometry

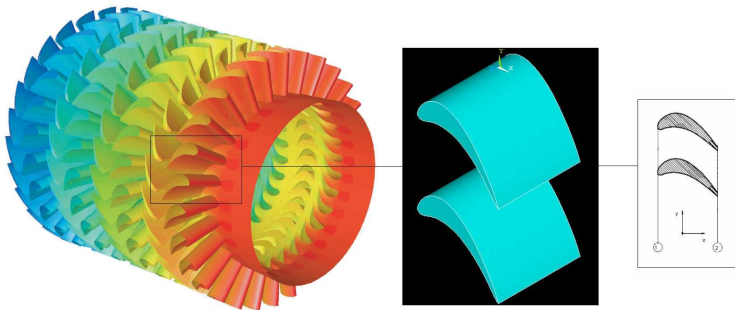


Figure: Problem domains (3D-2D)

Flow Between Turbomachinery Blades

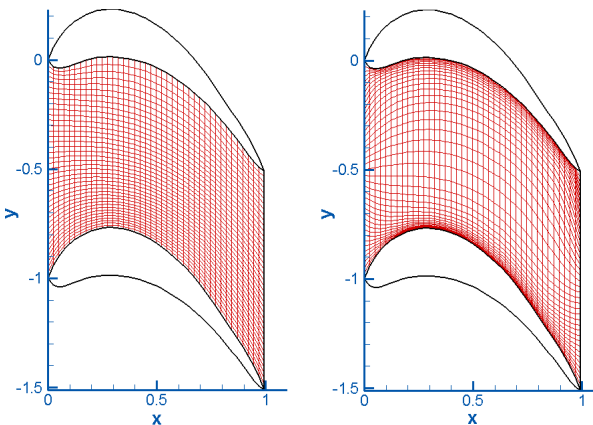


Figure: Controlling grid density

Flow Between Turbomachinery Blades

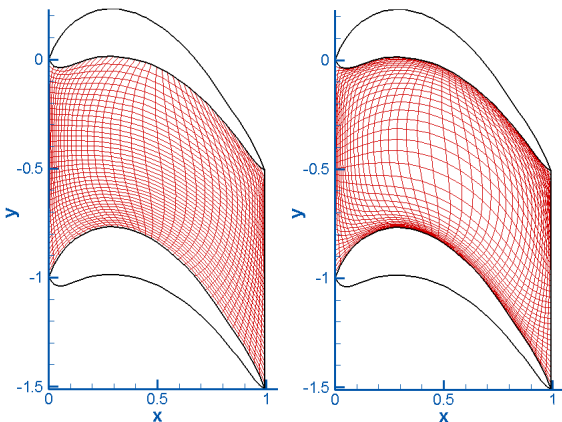


Figure: Controlling grid density

Flow Between Turbomachinery Blades

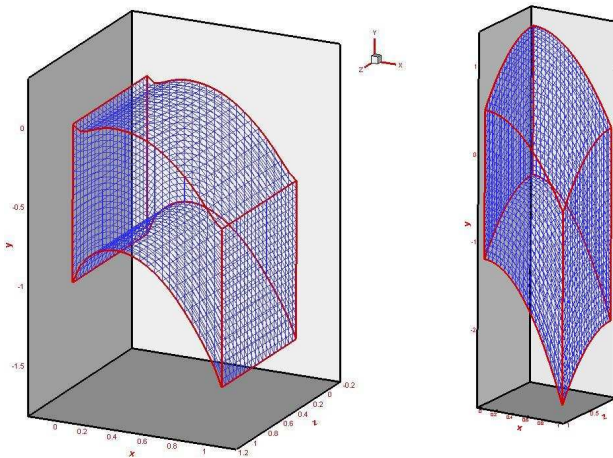


Figure: Problem domains in 3D

Formulations and Assumptions

- Stream Function-Vorticity formulation (2D)
- Velocity-Vorticity formulation (2D - 3D)
- Steady and incompressible (2D - 3D)
- Finite difference discretization (2D - 3D)



Governing Equations in 2D

Stream function - Vorticity:

$$\nabla^2 \Psi + \Omega = 0 \quad (7)$$

$$\nabla^2 \Omega - Re \left[\frac{\partial \Psi}{\partial y} \frac{\partial \Omega}{\partial x} - \frac{\partial \Psi}{\partial x} \frac{\partial \Omega}{\partial y} \right] = 0 \quad (8)$$

Velocity - Vorticity:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = -\frac{\partial \Omega}{\partial y} \quad (9)$$

$$\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} = \frac{\partial \Omega}{\partial x} \quad (10)$$

$$u \frac{\partial \Omega}{\partial x} + v \frac{\partial \Omega}{\partial y} = \frac{1}{Re} \left(\frac{\partial^2 \Omega}{\partial x^2} + \frac{\partial^2 \Omega}{\partial y^2} \right) \quad (11)$$



Coordinate Transformation

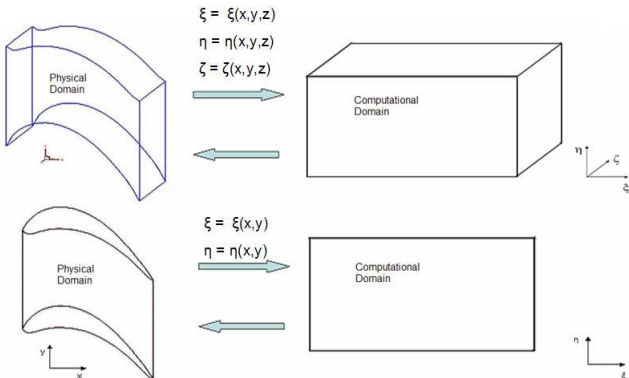


Figure: From Physical to Computational

Transformed Form of The Governing Equations in 2D

Stream Function Equation

$$\begin{aligned} \frac{\partial \psi}{\partial \eta} \left[\frac{\partial y}{\partial \xi} \left(\frac{\partial^2 x}{\partial \xi^2} \alpha + \frac{\partial^2 x}{\partial \eta^2} \beta - 2 \frac{\partial^2 x}{\partial \xi \partial \eta} \gamma \right) - \frac{\partial x}{\partial \xi} \left(\frac{\partial^2 y}{\partial \xi^2} \alpha + \frac{\partial^2 y}{\partial \eta^2} \beta - 2 \frac{\partial^2 y}{\partial \xi \partial \eta} \gamma \right) \right] \\ + \frac{\partial \psi}{\partial \xi} \left[\frac{\partial x}{\partial \eta} \left(\frac{\partial^2 y}{\partial \xi^2} \alpha + \frac{\partial^2 y}{\partial \eta^2} \beta - 2 \frac{\partial^2 y}{\partial \xi \partial \eta} \gamma \right) - \frac{\partial y}{\partial \xi} \left(\frac{\partial^2 x}{\partial \xi^2} \alpha + \frac{\partial^2 x}{\partial \eta^2} \beta - 2 \frac{\partial^2 x}{\partial \xi \partial \eta} \gamma \right) \right] \\ + \frac{\partial^2 \psi}{\partial \xi^2} J \alpha + \frac{\partial^2 \psi}{\partial \eta^2} J \beta - 2 \frac{\partial^2 \psi}{\partial \xi \partial \eta} J \gamma = -J^3 \Omega \end{aligned} \quad (12)$$

Vorticity Transport Equation

$$\begin{aligned} \text{Re} \frac{\partial \psi}{\partial \eta} \frac{\partial \Omega}{\partial \xi} J \left(\frac{\partial x}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial y}{\partial \xi} \frac{\partial x}{\partial \eta} \right) + \text{Re} \frac{\partial \psi}{\partial \xi} \frac{\partial \Omega}{\partial \eta} J \left(\frac{\partial x}{\partial \eta} \frac{\partial y}{\partial \xi} - \frac{\partial y}{\partial \eta} \frac{\partial x}{\partial \xi} \right) = \\ + \frac{\partial \Omega}{\partial \eta} \left[\frac{\partial y}{\partial \xi} \left(\frac{\partial^2 x}{\partial \xi^2} \alpha + \frac{\partial^2 x}{\partial \eta^2} \beta - 2 \frac{\partial^2 x}{\partial \xi \partial \eta} \gamma \right) - \frac{\partial x}{\partial \xi} \left(\frac{\partial^2 y}{\partial \xi^2} \alpha + \frac{\partial^2 y}{\partial \eta^2} \beta - 2 \frac{\partial^2 y}{\partial \xi \partial \eta} \gamma \right) \right] \\ + \frac{\partial \Omega}{\partial \xi} \left[\frac{\partial x}{\partial \eta} \left(\frac{\partial^2 y}{\partial \xi^2} \alpha + \frac{\partial^2 y}{\partial \eta^2} \beta - 2 \frac{\partial^2 y}{\partial \xi \partial \eta} \gamma \right) - \frac{\partial y}{\partial \eta} \left(\frac{\partial^2 x}{\partial \xi^2} \alpha + \frac{\partial^2 x}{\partial \eta^2} \beta - 2 \frac{\partial^2 x}{\partial \xi \partial \eta} \gamma \right) \right] \\ + \frac{\partial^2 \Omega}{\partial \xi^2} J \alpha + \frac{\partial^2 \Omega}{\partial \eta^2} J \beta - 2 \frac{\partial^2 \Omega}{\partial \xi \partial \eta} J \gamma \end{aligned}$$

Governing Equations in 3D

Equations for the velocity components in vector form:

$$\nabla^2 \bar{\mathbf{u}} = -\bar{\nabla} \times \bar{\mathbf{\Omega}} \quad (14)$$

The three-component vorticity transport equation,:

$$(\bar{\mathbf{u}} \cdot \bar{\nabla}) \bar{\mathbf{\Omega}} - (\bar{\mathbf{\Omega}} \cdot \bar{\nabla}) \bar{\mathbf{u}} - \frac{1}{\text{Re}} \nabla^2 \bar{\mathbf{\Omega}} = 0 \quad (15)$$



Boundary conditions

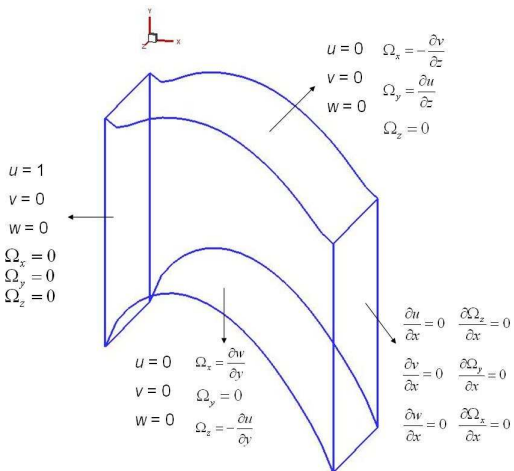
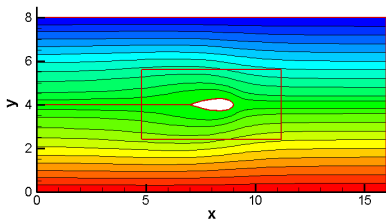
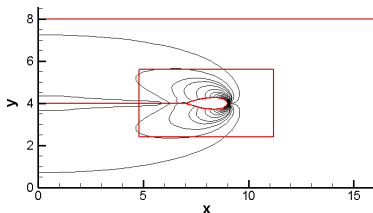


Figure: Boundary conditions

Flow Past an Airfoil



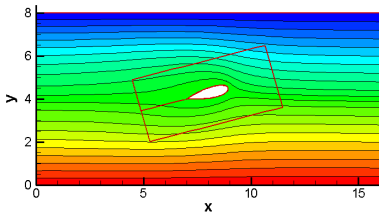
(a) Stream function contours



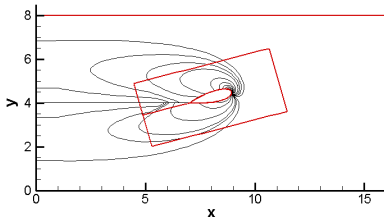
(b) Vorticity contours

Figure: Flow Past an Airfoil ($Re=50$)

Flow Past an Airfoil



(a) Stream function contours



(b) Vorticity contours

Figure: Flow Past an Airfoil ($Re=50$)

Stream function contours

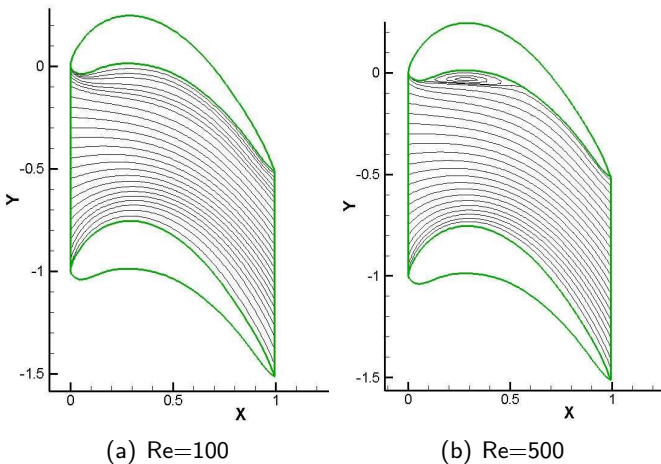


Figure: Stream function contours

Vorticity contours

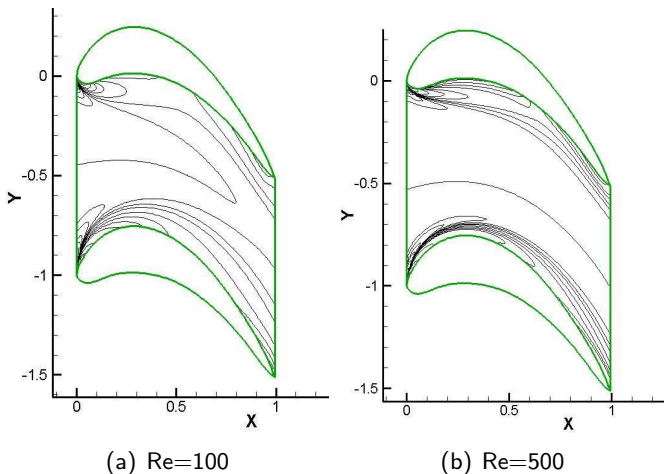
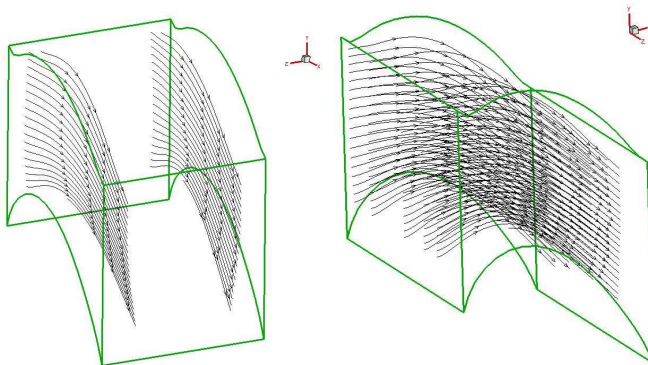


Figure: Vorticity contours

Streamtrace contours

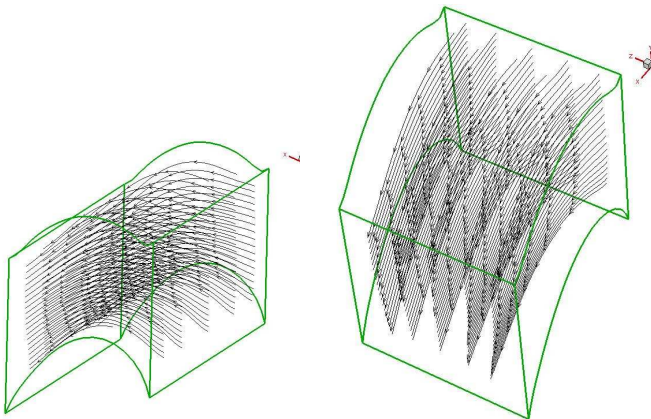


(a) Streamtrace contours

(b) Streamtrace contours

Figure: Velocity-Vorticity Approach

Streamtrace contours



(a) Streamtrace contours

(b) Streamtrace contours

Figure: Velocity-Vorticity Approach

Vorticity contours

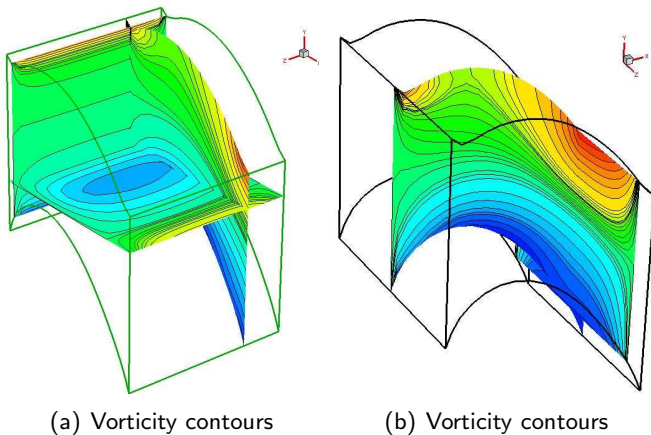
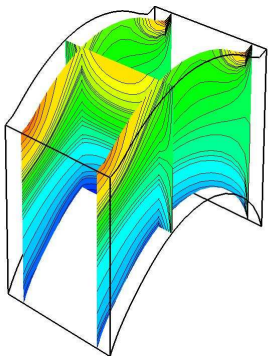
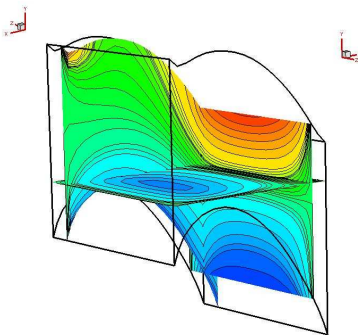


Figure: Velocity-Vorticity Approach

Vorticity contours



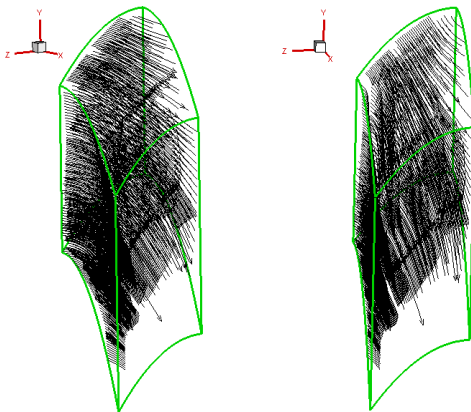
(a) Vorticity contours



(b) Vorticity contours

Figure: Velocity-Vorticity Approach

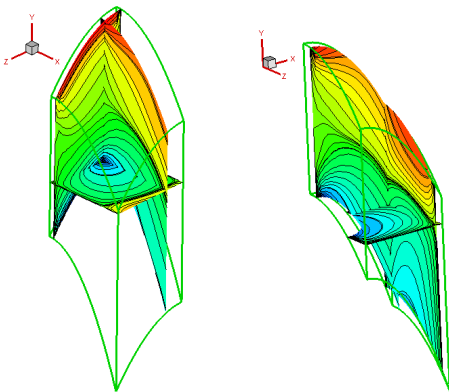
Streamtrace contours



(a) Streamtrace contours (b) Streamtrace contours

Figure: Velocity-Vorticity Approach

Streamtrace contours



(a) Vorticity contours (b) Vorticity contours

Figure: Velocity-Vorticity Approach

Performance

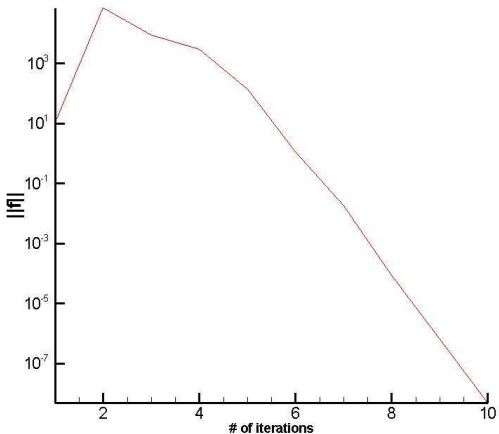


Figure: Typical non-linear convergence history for $Re=100$

Comparison of the Solvers and the Preconditioners

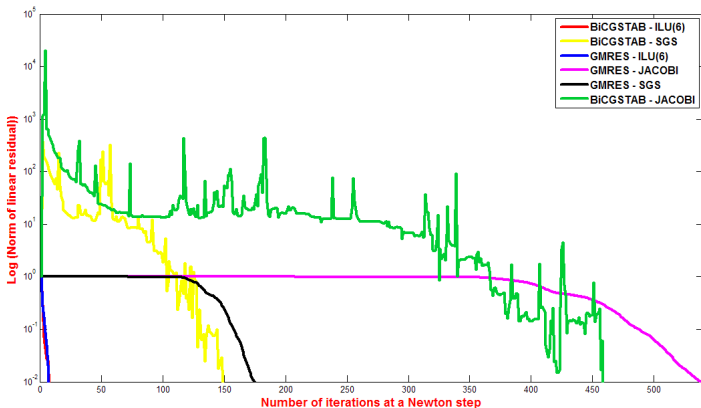


Figure: Comparison of the Solvers and the Preconditioners

Comparison of the Solvers and the Preconditioners

- Implementation of an effective preconditioner is crucial.
- Jacobi does not have a major effect on the convergence pattern.
- BiCGSTAB has a more stable pattern than GMRES and is the fastest solver by means of iteration steps and computation time.
- GMRES is the most stable solver, enables a continuous residual reduction.
- In GMRES the computational work increases linearly with the iteration \Rightarrow Restartable GMRES(m)



Fluid-Structure Interaction

The idea:

- Extract grid points on the blade surface from ANSYS to an aerodynamic code
- Generate grid, execute fluid analysis and achieve the velocity, vorticity and pressure fields
- Transfer pressure values to ANSYS, execute solid analysis and acquire displacements
- With storing displacements from ANSYS, finish first step of FSI
- By using displacements from ANSYS, update geometry and start second step



Future Works

- Parallel Computing
- Nonlinear Preconditioning
- Multigrid/Multilevel Techniques

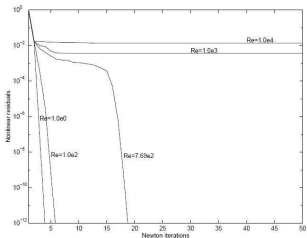


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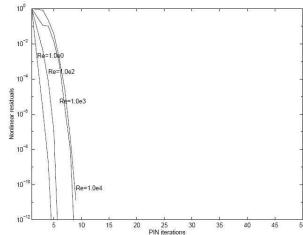
Thank You...



Non-linear & Linear Preconditioning



(a) Newton iterations



(b) PIN iterations

Figure: Nonlinear residual history for the flow problem with different Reynolds number [7]

Solid Analysis

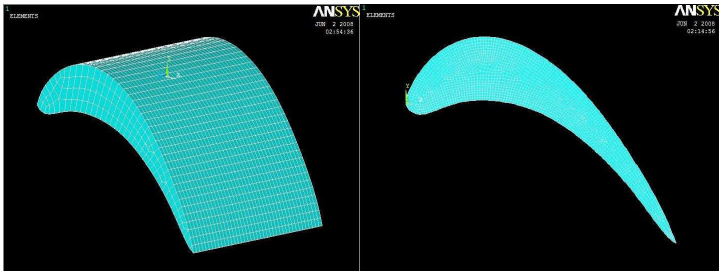
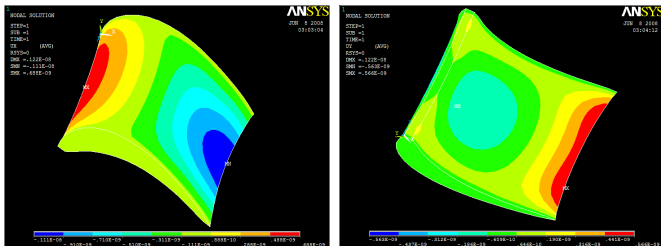


Figure: Blade mesh in ANSYS (3D-2D)

- Hexahedral - Quadrilateral
- MMB: Linear - Elastic - Isotropic

Deformation Under Aerodynamic Loads

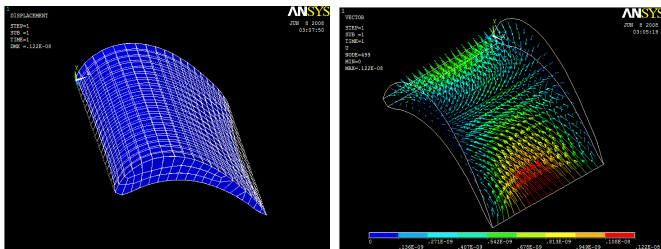


(a) x-component of displacement

(b) y-component of displacement

Figure: Solid Analysis

Deformation Under Aerodynamic Loads



(a) Deformed shape with undeformed edge (b) Vector plot of translation

Figure: Solid Analysis

Transformation Terms

$$\begin{aligned}\alpha &= \left(\frac{\partial x}{\partial \eta}\right)^2 + \left(\frac{\partial y}{\partial \eta}\right)^2 \\ \beta &= \left(\frac{\partial x}{\partial \xi}\right)^2 + \left(\frac{\partial y}{\partial \xi}\right)^2 \\ \gamma &= \frac{\partial x}{\partial \xi} \frac{\partial x}{\partial \eta} + \frac{\partial y}{\partial \xi} \frac{\partial y}{\partial \eta} \\ J &= \frac{\partial x}{\partial \xi} \frac{\partial y}{\partial \eta} - \frac{\partial y}{\partial \xi} \frac{\partial x}{\partial \eta}\end{aligned}\tag{16}$$

Governing Equations in 3D

x-component of velocity :

$$\nabla^2 u = \frac{\partial \Omega_y}{\partial z} - \frac{\partial \Omega_z}{\partial y} \quad (17)$$

y-component of velocity :

$$\nabla^2 v = \frac{\partial \Omega_z}{\partial x} - \frac{\partial \Omega_x}{\partial z} \quad (18)$$

z-component of velocity :

$$\nabla^2 w = \frac{\partial \Omega_x}{\partial y} - \frac{\partial \Omega_y}{\partial x} \quad (19)$$

vorticity transport :

$$\nabla \cdot (u\Omega) - (\Omega \cdot \nabla) u - \frac{1}{\text{Re}} \nabla^2 \Omega = 0 \quad (20)$$

Coordinate Transformation

- The chain rule of partial differentiation

- 1 $\frac{\partial}{\partial x} = \xi_x \frac{\partial}{\partial \xi} + \eta_x \frac{\partial}{\partial \eta} + \zeta_x \frac{\partial}{\partial \zeta}$

- 2 $\frac{\partial}{\partial y} = \xi_y \frac{\partial}{\partial \xi} + \eta_y \frac{\partial}{\partial \eta} + \zeta_y \frac{\partial}{\partial \zeta}$

- 3 $\frac{\partial}{\partial z} = \xi_z \frac{\partial}{\partial \xi} + \eta_z \frac{\partial}{\partial \eta} + \zeta_z \frac{\partial}{\partial \zeta}$

- $J = \frac{\partial(\xi, \eta, \zeta)}{\partial(x, y, z)}$



Streamtrace contours

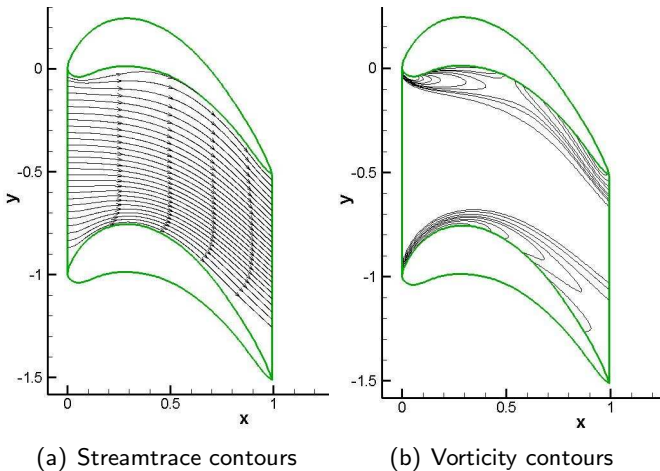


Figure: Velocity-Vorticity Approach

Preconditioning and Matrix-free Approach

$$P_L^{-1} J(\bar{x}) \bar{v} = P_L^{-1} \frac{\bar{f}(\bar{x} + P_R^{-1} \bar{y}) - \bar{f}(\bar{x})}{\varepsilon} = \bar{z} \quad (21)$$

$$P_R \epsilon \bar{v} = \bar{y} \quad (22)$$

$$P_L \bar{z} = J \bar{v} \quad (23)$$



Comparison of the Solvers and the Preconditioners

Table: Comparison of the solvers and preconditioners by means of iteration number and computation time

Method	# of iterations per Newton Step (average)	Time (sec.)
Jacobi-BiCGSTAB	490.5	25.7
SGS-BiCGSTAB	108.8	23.1
ILU(6)-BiCGSTAB	16.2	187.8
Jacobi-GMRES(100)	313.6	26.1
SGS-GMRES(100)	97.7	22.7
ILU(6)-GMRES(100)	12.5	182.14



Coordinate Transformation

Second derivative of an arbitrary variable in the transformed form

$$\begin{aligned} \frac{\partial^2 u}{\partial x^2} &= \frac{\partial u}{\partial \xi} \frac{\partial^2 \xi}{\partial x^2} + \frac{\partial u}{\partial \eta} \frac{\partial^2 \eta}{\partial x^2} + \frac{\partial u}{\partial \zeta} \frac{\partial^2 \zeta}{\partial x^2} + \frac{\partial^2 u}{\partial \xi^2} \left(\frac{\partial \xi}{\partial x} \right)^2 + \frac{\partial^2 u}{\partial \eta^2} \left(\frac{\partial \eta}{\partial x} \right)^2 \\ &+ \frac{\partial^2 u}{\partial \zeta^2} \left(\frac{\partial \zeta}{\partial x} \right)^2 + 2 \frac{\partial^2 u}{\partial \xi \partial \eta} \frac{\partial \xi}{\partial x} \frac{\partial \eta}{\partial x} + 2 \frac{\partial^2 u}{\partial \xi \partial \zeta} \frac{\partial \xi}{\partial x} \frac{\partial \zeta}{\partial x} + 2 \frac{\partial^2 u}{\partial \zeta \partial \eta} \frac{\partial \zeta}{\partial x} \frac{\partial \eta}{\partial x} \end{aligned}$$



Blade Shape

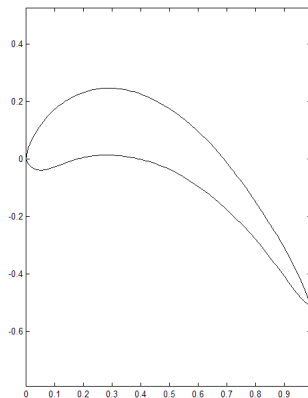
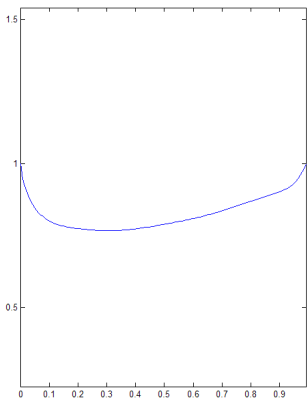


Figure: Blade geometry

Effect of Reynolds Number

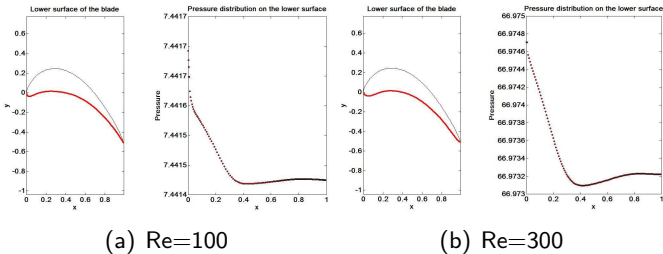
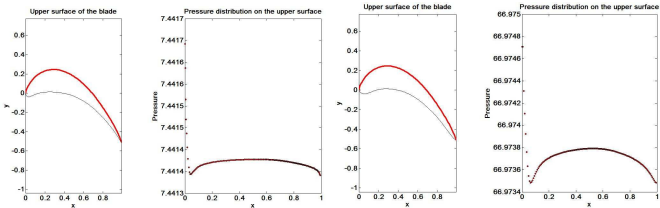


Figure: Pressure distribution

Effects of Reynolds Number



(a) $Re=100$

(b) $Re=300$

Figure: Pressure distribution

Effects of Reynolds Number

- Pressure distribution
distribution character \Rightarrow slightly
values \Rightarrow extensively
- Vortex generation (bottom surface of upper blade)