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Alignment and optical verification of DESHIMA 2.0 at ASTE

Arend Moerman^{a,*}, Kenichi Karatsu^{a,b}, Jochem J. A. Baselmans^{a,b,c},
Shahab O. Dabironezare^{a,b}, Shinji Fujita^d, Robert Huiting^b, Kotaro Kohno^{e,f},
Yuri Nishimura^e, Fenno Steenvoorde^g, Tatsuya Takekoshi^h, Yoichi Tamuraⁱ,
Akio Taniguchi^h, Stephen J. C. Yates^j, Bernhard R. Brandl^{k,l}, and Akira Endo^a

^aDelft University of Technology, Faculty of Electrical Engineering, Mathematics and Computer Science, Delft, The Netherlands

^bSRON-Netherlands Institute for Space Research, Leiden, The Netherlands

^cUniversität zu Köln, Physikalisches Institut, Cologne, Germany

^dThe Institute of Statistical Mathematics, Tachikawa, Japan

^eThe University of Tokyo, Institute of Astronomy, Graduate School of Science, Mitaka, Japan

^fThe University of Tokyo, Graduate School of Science, Research Center for the Early Universe, Bunkyo-ku, Japan

^gDelft University of Technology, DEMO: Electronic and Mechanical Support Division, Delft, The Netherlands

^hKitami Institute of Technology, Kitami, Japan

ⁱNagoya University, Graduate School of Science, Department of Physics, Nagoya, Japan

^jSRON-Netherlands Institute for Space Research, Groningen, The Netherlands

^kLeiden University, Leiden Observatory, Leiden, The Netherlands

^lDelft University of Technology, Faculty of Aerospace Engineering, Delft, The Netherlands

ABSTRACT. We developed, characterized, and verified an alignment procedure for the DESHIMA 2.0 instrument, an ultra-wide-band spectrometer operating between 200 and 400 GHz, at the ASTE telescope. To this end, we mounted the warm optics, consisting of a modified Dragonian dual reflector system, on a motor-controlled hexapod. Crucial in the alignment procedure is our sky chopper, which allows fast beam switching. It has a small entrance and exit aperture coupling to the (cold) sky, which creates a measurable signal with respect to the warm cabin environment. By scanning the instrument beam across the entrance aperture of the sky chopper using the hexapod, we found the hexapod configuration that produced the lowest signal on our detectors, implying that the beam is coupled fully to the cold sky and not the warm cabin. We first characterized the alignment procedure in the laboratory, where we used a vat containing liquid nitrogen as the cold source behind the sky chopper. Then, we applied the alignment procedure to DESHIMA 2.0 at ASTE. We found that the alignment procedure significantly improved the aperture efficiency compared with previously reported values of the aperture efficiency of DESHIMA at ASTE, which indicates the veracity of the alignment procedure.

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1 Introduction

The Deep Spectroscopic High-redshift Mapper (DESHIMA) is an ultra-wide-band (UWB) sub-millimeter (submm) integrated superconducting spectrometer (ISS) using MKIDs.¹ The primary

*Address all correspondence to Arend Moerman, A.Moerman-1@tudelft.nl

design (DESHIMA 1.0²) has instantaneously covered 332 to 377 GHz by 49 frequency channels, and the demonstration instrument has successfully seen first light in 2017³ at the 10-m Atacama Submillimeter Telescope Experiment (ASTE)⁴ telescope in Chile. DESHIMA 2.0 is a successor to DESHIMA 1.0 with the upgrades, including an increase in the bandwidth (200 to 400 GHz), the number of frequency channels (347 channels), the optical efficiency, and the addition of a fast sky-position chopper for atmospheric subtraction, among others.⁵ The science cases range from multiline spectroscopy of dusty star-forming galaxies to calculate spectroscopic redshifts and study their physical conditions,⁶ broad-band continuum observations of asymptotic giant branch stars, to galaxy cluster diagnostics through the Sunyaev-Zel'dovich effect.^{7,8}

The DESHIMA 2.0 design is based on a superconducting on-chip filter bank,^{9,10} fed by a leaky-lens antenna.^{11–13} The filter bank is connected to an NbTiN-Al hybrid MKID¹⁴ array, which is read out using frequency-domain multiplexing.¹⁵ The integrated nature of the DESHIMA architecture and the multiplexed readout makes it inherently scalable, and as such, it plays an important role in future integral field unit (IFU) designs,¹⁶ such as the Terahertz Integral Field Unit with Universal Nanotechnology (TIFUUN) instrument, which is currently in its early stages of development.¹⁷ The cryostat housing DESHIMA 2.0 is an adiabatic demagnetization refrigerator (ADR), where it operates at a temperature of 120 mK.

The submillimeter beam out of the vacuum window of DESHIMA can be misaligned from the design by as much as ~ 4 deg,² which can lead to significant degradation of the aperture efficiency when mounted on the ASTE telescope without sufficient correction.² Furthermore, because the cold optics are partially disassembled during transportation, there is no guarantee that the misalignment will be persistent from the laboratory to the telescope, requiring measurement of the misalignment in the telescope cabin. Fortunately, DESHIMA 2.0 employs a modified Dragonian dual reflector¹⁸ at room temperature for magnification and coupling to the ASTE subreflector, and the positions of these mirrors can be adjusted to correct for potential misalignment.

Such a correction for misalignment by one or more room-temperature optical components is common in mm/submm astronomy although the exact methodology can vary depending on the system. A common method employs an optical laser to align the optics at the telescope. Examples include CONCERTO at APEX^{19,20} and NIKA2 at IRAM.²¹ However, this is an indirect method that assumes the submillimeter beam to be aligned to the laser ray, which is not the case for DESHIMA 2.0. The SOFIA instrument has used radiometric alignment by introducing an additional component that intercepts the beam.²² An alternative approach is to measure the mechanical positions of the optical components by means of a Faro measurement arm and correct for any misalignment²³ although the position of the cryogenic mirrors cannot be measured.

Recently, Moerman et al.²⁴ demonstrated a method for aligning the DESHIMA 2.0 instrument in the laboratory. They measured the complex (phase and amplitude) submm beam to extract misalignment information^{25,26} and used this, in combination with a ray-tracing algorithm, to align the instrument using a modified Dragonian dual reflector and a mechanical hexapod. They argued that this method can correct for lateral shifts and tilt, but that a separate subreflector alignment should be performed to correct for the vertical shift or defocus. However, applying this method at ASTE would require installing a coherent submm source and an XY-scanning stage in the small ASTE receiver cabin. Moreover, the presence of the aforementioned sky-position chopper significantly complicates the application of the alignment method of Moerman et al.²⁴

In this paper, we describe an alignment method that directly uses the instrument beam for alignment, without the need for quantification of misalignment using a visible light laser or a complex beam pattern measurement. The new method makes use of the temperature contrast between the room temperature entrance and exit apertures of the sky chopper and a cold background source, in combination with the modified Dragonian dual reflector and mechanical hexapod also used by Moerman et al.²⁴ The main advantage of the proposed method is the lack of need for elaborate measurement setups, making it simpler to execute than the procedure in Moerman et al.²⁴ Instead, only a cold source, such as the atmosphere, is necessary. This makes the proposed method attractive for implementation at telescopes with limited receiver cabin space, such as the ASTE telescope. The method only has to be performed once after installation of the dual reflector, hexapod, and sky chopper. In addition, the method can be executed fully

remotely over an internet connection, eliminating the need to be physically present at the telescope site. This allows for remote checking of the alignment whenever it is desired.

The paper is structured as follows. We start by briefly discussing the optical path of DESHIMA 2.0 at ASTE in Sec. 2.1. Particular focus here will be given to the sky chopper. In Sec. 2.2, we discuss the general methodology behind the alignment method. In addition, we discuss the alignment of the ASTE subreflector in Sec. 2.3. In Sec. 2.4, we show the laboratory setup used for the characterization of the alignment method. We also show the telescope setup at ASTE that was actually aligned using the proposed method. In Sec. 3.1, we show the results from the laboratory characterization, and in Sec. 3.2, we show results from the actual alignment at ASTE.

2 Methods

2.1 Optical Path at ASTE

We start by briefly describing relevant elements in the optical path of DESHIMA 2.0 at ASTE. A sketch of the setup and global coordinate system can be found in Fig. 1(a) and a CAD render in Fig. 1(b). A photograph of the ASTE telescope can be found in Fig. 1(c).

The first optical components seen from the DESHIMA leaky-lens antenna are two off-axis paraboloidal mirrors, M_6 and M_5 , which are placed in a parabolic relay²⁷ configuration. This compensates for the inherent aberration introduced by a single off-axis paraboloidal mirror.²⁸ One focus of the parabolic relay overlaps the leaky-lens antenna, and the other is located in the right of the relay in Fig. 1(a) and denoted the cold focus (shown as p_{CF}). Both M_6 and M_5 are located inside the cryostat and are therefore referred to as the cold optics. The 30-deg rotation of the cryostat around the x -axis [see Fig. 1(b)] is applied so that the cryostat is upright when the telescope is tilted to an elevation of 60 deg. This ensures that the cryostat is mostly upright during regular observations.

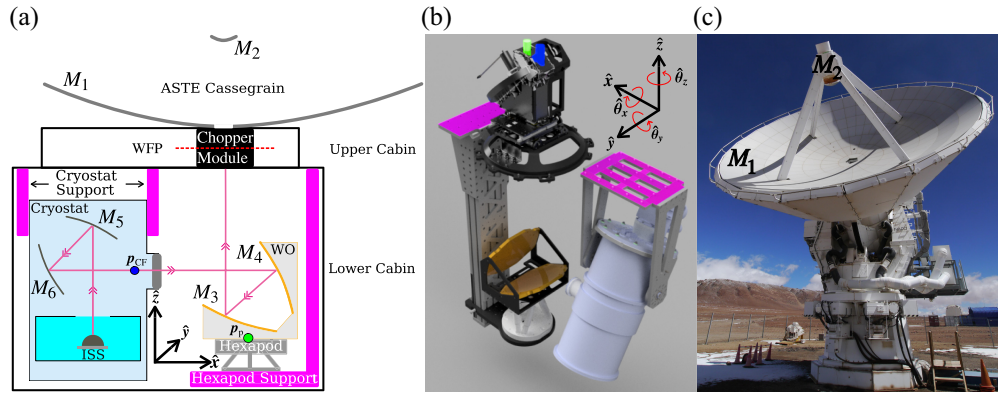


Fig. 1 Overview of the optical setup of DESHIMA 2.0 at ASTE including the sky chopper. (a) Sketch of the setup at the ASTE telescope. Illustrated are the integrated superconducting spectrometer (ISS), cryostat and cryostat optics M_6 and M_5 , the cold focus p_{CF} (blue dot), the warm optics M_4 and M_3 , the hexapod, support structures for the hexapod and cryostat, and the chopper module. The chopper module is located in the upper cabin of ASTE, and the rest in the lower cabin. The warm focal plane, denoted WFP in the sketch, is illustrated with a red dashed line and runs through the sky chopper. The coordinate axes shown are used throughout the paper as the basis coordinates. The origin of the coordinate system is the cold focus, unless stated otherwise. The hexapod pivot p_p , denoted by the green dot, is located in the center of the hexapod stage. The Cassegrain setup at ASTE, after the chopper module, consists of a hyperboloid subreflector M_2 and a paraboloid primary M_1 . (b) CAD render of the setup in panel (a), showing the sky chopper and the two optical paths (in green and blue) coming out of the top of the chopper. We show the coordinate axes again to emphasize the orientation of the system. In addition, we show (right-handed) rotation axes in this coordinate system. (c) Photograph of the ASTE telescope in the Atacama desert, Pampa la Bola, Chile. Both M_2 and M_1 are illustrated in the figure. The lower cabin is visible below M_1 . The hexagonal hole in the vertex of M_1 is the entrance of light from the sky into the upper cabin.

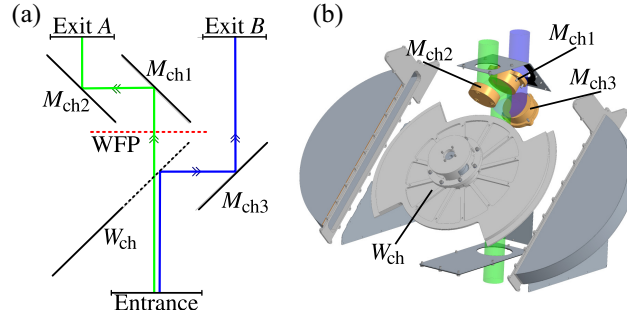


Fig. 2 Overview and sketch of the sky chopper. (a) Schematic sketch showing the location of the chopper wheel W_{ch} and the three planar mirrors M_{ch1} , M_{ch2} , and M_{ch3} . At the bottom of the sketch, we denote the entrance aperture. Both exit apertures A and B are also shown. Note also that the color coding of the beam through A and B matches the colors of the beams coming out of the sky chopper in Fig. 1(b). The warm focal plane, denoted WFP, runs through the middle of the sky chopper and is denoted by the horizontal red dashed line. (b) Exploded view of a CAD render of the chopper module, showcasing the contents of the sky chopper. The illustrated green and blue tubes do not represent the actual shape of the DESHIMA 2.0 beam but are for illustrative purposes only.

After the cold optics, the beam of the ISS encounters the warm optics. It consists of a hyperboloidal mirror M_4 and an ellipsoidal mirror M_3 , which are placed in a modified Dragonian¹⁸ setup and monolithically mounted onto a motor-controlled hexapod. One focus of the warm optics overlaps the cold focus, the other is situated above the warm optics and is called the warm focus. The warm focus is enclosed within the sky chopper and lies in the warm focal plane (WFP). The WFP normal is oriented along the z -axis.

The sky chopper is a device containing three planar mirrors, M_{ch1} , M_{ch2} , and M_{ch3} , and a slotted, reflective wheel, which we will denote W_{ch} . The beam enters the sky chopper through the entrance aperture, after which it encounters W_{ch} . As W_{ch} rotates, half of the time the beam passes through the slots in the wheel, onto M_{ch1} and M_{ch2} , leaving the sky chopper through the exit aperture A . We will denote the beam coming out of the exit aperture A beam A . If W_{ch} is closed, the beam reflects off W_{ch} , onto M_{ch3} , and out through the exit aperture B . This beam we denote beam B . For a detailed overview of the sky chopper and its contents, see Fig. 2. The entrance and exit apertures fulfill a crucial role in the alignment procedure by constraining the beam location and tilt as it passes through the sky chopper, essentially coupling the beam to the warm sky chopper aperture/interior and cabin environment if the alignment is poor, and coupling it to the cold sky if the alignment is good. In Sec. 2.2, we will explain this in more detail.

After exiting the chopper module, both beams encounter M_2 , the secondary (hyperboloidal) reflector of the ASTE Cassegrain. After M_2 , the beams encounter M_1 , the paraboloidal primary of the ASTE Cassegrain. Beam A (B) illuminates M_5 50 (−50) mm off-center along the x -axis. This causes beam A (B) to have a 117 (−117) arcsec pointing offset along the azimuthal axis with respect to the broadside, which is fixed by design.

2.2 Warm Optics Alignment

The alignment technique relies on the warm/cold load contrast on the MKIDs as the instrument beam is scanned across the sky chopper entrance aperture using the warm optics mounted on the hexapod. Both paths of the chopper (see Fig. 2) essentially represent a cylinder, each with a unique exit aperture A and B and a shared entrance aperture. The sky chopper interior and apertures have a physical temperature T_0 , which we assume to be constant throughout the alignment.

The atmosphere has an effective physical temperature T_{atm} , which can be converted into the line-of-sight (LOS) brightness temperature T_{sky} using

$$T_{sky}(\nu, PWV) = (1 - \eta_{atm}^{csc(e)}(\nu, PWV))T_{atm}, \quad (1)$$

where $\eta_{atm}(\nu, PWV)$ is the atmospheric transmission coefficient toward the zenith, ν is the frequency at which the atmosphere is observed, PWV is the precipitable water vapor toward the zenith, and e is the elevation of the telescope at which the atmosphere is observed. We assume for

the current treatment that there are only two relevant brightness temperatures, T_0 for the sky chopper interior, apertures, and the ASTE receiver cabin in general, and T_{sky} for the atmosphere.

Let \mathbf{p}_{hex} and $\boldsymbol{\theta}_{\text{hex}}$ denote the hexapod position and orientation, respectively. We define them as follows:

$$\mathbf{p}_{\text{hex}} = X\hat{\mathbf{x}} + Y\hat{\mathbf{y}} + Z\hat{\mathbf{z}}, \quad (2a)$$

$$\boldsymbol{\theta}_{\text{hex}} = U\hat{\boldsymbol{\theta}}_x + V\hat{\boldsymbol{\theta}}_y + W\hat{\boldsymbol{\theta}}_z. \quad (2b)$$

Note that Eqs. (2a) and (2b) represent translations and rotations performed solely by the motion of the hexapod and not by how it is mounted or placed. The coordinate system used in Eqs. (2a) and (2b) is the same as in Fig. 1(b), with the origin placed in the pivot \mathbf{p}_p . We define the hexapod configuration to be the pair $(\mathbf{p}_{\text{hex}}, \boldsymbol{\theta}_{\text{hex}})$ and the hexapod home position to be the configuration $(\mathbf{0}, \mathbf{0})$, i.e., $X = Y = Z = 0$ mm and $U = V = W = 0$ deg. The home configuration is the configuration of the hexapod after installation in the ASTE receiver cabin, before the alignment procedure, and hence has an arbitrary alignment. It therefore serves as an important reference point for comparison of the instrument performance after the alignment procedure.

Let $\delta x(\mathbf{p}_{\text{hex}}, \boldsymbol{\theta}_{\text{hex}})$ denote the response of a single MKID on DESHIMA 2.0 as a function of the hexapod configuration. We adopt the definition given by Takekoshi et al.²⁹

$$\delta x(\mathbf{p}_{\text{hex}}, \boldsymbol{\theta}_{\text{hex}}) = \frac{f_0 - f(T_b)}{f_0}, \quad (3)$$

where $f(T_b)$ represents the resonance frequency of the MKID, loaded at brightness temperature $T_b(\mathbf{p}_{\text{hex}}, \boldsymbol{\theta}_{\text{hex}})$, and f_0 the resonance frequency of the MKID, loaded at some reference brightness temperature $T_b = T_0$, which we take as room temperature.² We explicitly include the dependency of T_b on the hexapod configuration to stress the fact that, as the hexapod changes configuration, the beam will be scanned across the sky chopper entrance aperture and hence see a varying brightness temperature distribution.

Consider the case where the optical system is so misaligned that the entire beam terminates on the entrance aperture or on the chopper interior. In this case, $T_b \sim T_0$ and $\delta x \sim 0$. When $T_{\text{sky}} < T_b < T_0$, the beam partly goes through the chopper module and couples to the atmosphere but also partly couples to the chopper aperture/interior, and $\delta x < 0$. In this case, the beam is less misaligned, but the MKID is still partly loaded by the sky chopper aperture/interior. When $T_b \sim T_{\text{sky}}$, the entire beam couples to the atmosphere. Here, δx reaches a minimum, and we can conclude that the beam is well-aligned because it goes through the chopper completely. The goal of the warm optics alignment then is to find a configuration $(\mathbf{p}_{\text{hex}}, \boldsymbol{\theta}_{\text{hex}})$ such that δx is minimized.

The alignment can be done with the chopper in either position. This is because the beam from DESHIMA 2.0 scans the entrance aperture, which is shared between beam *A* and *B*. In addition, the distances between the entrance aperture and exit apertures *A* and *B* are roughly the same, and therefore, a misaligned beam, entering through the entrance aperture, will experience the same warm spillover in both paths. Most important is the presence of two small co-aligned apertures. With only a single aperture, the beam position can be constrained, but the beam tilt can still be nonzero. Only upon including the second aperture can the beam tilt also be constrained.

2.3 Subreflector Alignment

Because the warm optics cannot correct for defocus and lateral misalignment simultaneously,²⁴ the subreflector is used for refocusing the optical system after the warm optics alignment. The subreflector alignment consists of two steps, which must be repeated at different elevations because of the elevation-dependent gravitational deformation of ASTE.³⁰ The first step consists of scanning the telescope pointing in azimuth and elevation across a small, bright source to measure the peak brightness and repeating this at various subreflector (x, y) combinations to overlap the Cassegrain focus with the warm focus in the x - y plane. The second step consists of scanning the subreflector along the z -axis while observing a small, bright source, to overlap the Cassegrain focus with the warm focus along the z -axis. When the foci overlap, the response of the bright source will be maximal, and the x , y , and z values corresponding to the maximum are selected as

the optimal subreflector (x, y, z) position. The focal optimization of the subreflector is done using DESHIMA channels at around 350 GHz.

2.4 Experimental Setup

We use two different experimental setups for the characterization, a laboratory and telescope setup. The laboratory setup is designed in such a way to emulate the telescope setup, which can be seen in Fig. 1(b). In this way, we can develop and characterize the alignment method in the laboratory before going to the ASTE telescope to apply the method. We omitted the 30-deg rotation of the cryostat around the x -axis and added a -90 -deg rotation of the hexapod, support structure, and warm optics as the beam measurement setup is too large to be placed upright in the laboratory, similar to Moerman et al.²⁴ For the definition of the hexapod configuration in the laboratory setup, we also rotate the coordinate system used in Eqs. (2a) and (2b) by -90 deg around the x -axis so that it is easier to compare the hexapod characterization in the laboratory with the characterization at ASTE. Because it is infeasible to use the atmosphere as the cold source in the laboratory setup, we place a cold absorber, immersed in a vat containing liquid nitrogen (N_2), behind the sky chopper. Because of the temperature difference between N_2 (77 K) and room temperature (~ 300 K), we expect ample contrast between the absorber and sky chopper and hence, an accurate alignment. For a photograph of the laboratory setup, see Fig. 3(a). In Fig. 3(b), we show the telescope setup.

3 Results

3.1 Laboratory Characterization

Here, we describe the laboratory results of scanning the DESHIMA 2.0 beam across the sky chopper entrance aperture using the hexapod and warm optics with the liquid N_2 behind the chopper entrance aperture. We use the DESHIMA channel at 250 GHz for alignment.

We started by scanning along the X and Y degrees-of-freedom (DoFs) of the hexapod, centered at $X = Y = 0$ mm. We chose this combination because these are nondegenerate, meaning that a change in X or in Y tilts/displaces the beam around/along separate axes, as was shown by the hexapod characterization performed by Moerman et al.²⁴ The result of this initial scan can be

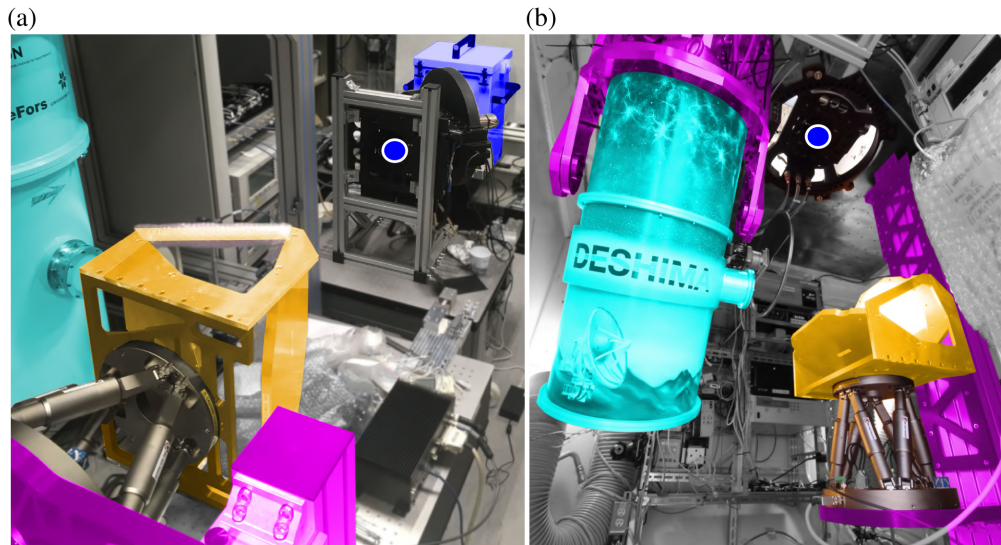


Fig. 3 Overview of the two experimental setups used in this work. (a) The laboratory setup used for the initial characterization of the alignment method. To the left, in light blue, is the cryostat. The warm optics (orange) are attached to the hexapod, which in turn is attached to the support structure (purple). The cold source, an absorber immersed in a vat containing liquid N_2 (dark blue), is placed behind the chopper module (black). The chopper itself is supported by a rectangular frame. The blue disk on the chopper indicates the entrance aperture. (b) The setup used inside the ASTE lower cabin. The color coding is the same as for panel (a), but because we now use the atmosphere instead of the cold absorber, there is no blue-colored source behind the chopper.

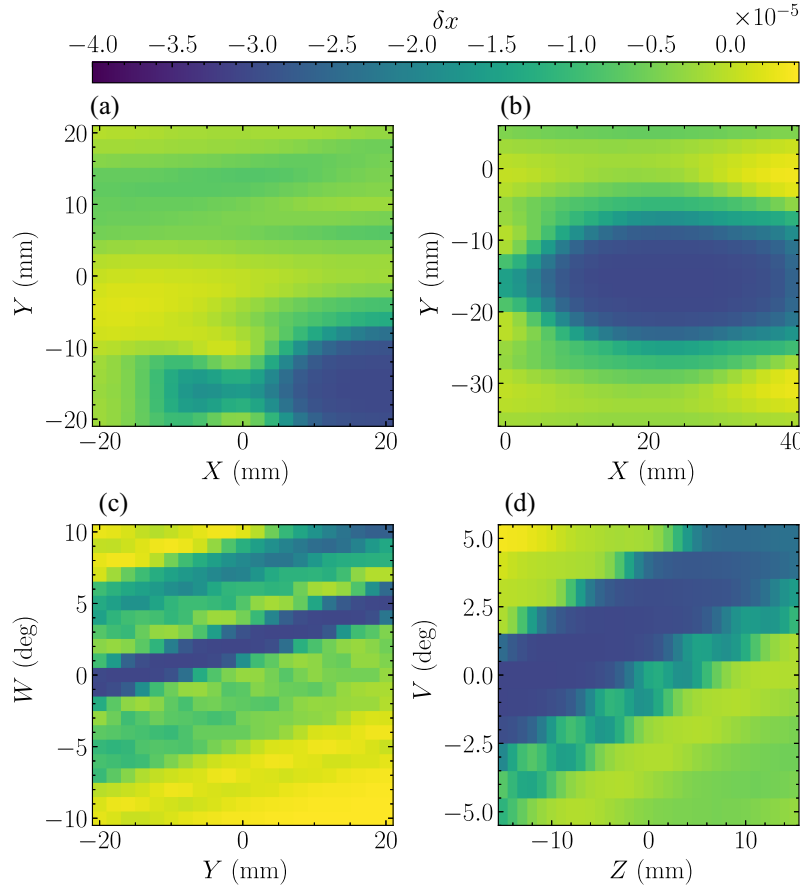


Fig. 4 Instrument responses δx as a function of hexapod scanning DoFs, for the four hexapod scans used for the laboratory characterization. Shown responses are at a channel frequency $\nu = 250$ GHz. We extended the color bar to $\delta x = -4 \cdot 10^{-5}$ to enhance contrast inside the cold spots. (a) The X–Y scan centered at $X = Y = 0$ mm. (b) The X–Y scan centered at $X = 20$ mm and $Y = -15$ mm. (c) The Y–W scan centered at $Y = 0$ mm and $W = 0$ deg. (d) The Z–V scan centered at $Z = 0$ mm and $V = 0$ deg.

found in Fig. 4(a). We see that δx varies significantly over the scan, and it appears that the minimum of δx is located in the lower right of the scan. Because the DoFs are nondegenerate, we see that there is a single X–Y combination for which δx reaches a minimum. Smoothly varying either X or Y around the minimum X–Y combination increases δx , giving rise to a localized δx dip, which we call a cold spot.

For the second X–Y scan, we moved the hexapod to $X = 20$ mm and $Y = -15$ mm, as suggested by Fig. 4(a), and scanned the DoFs around these values. This result can be found in Fig. 4(b). It is immediately visible that the minimum of δx is now centered on the scan, indicating that this combination of X and Y already produces a better alignment compared with $X = Y = 0$ mm.

After the second X–Y scan, we were interested in whether we can use Z in the alignment because the hexapod DoF characterization performed by Moerman et al.²⁴ showed that X and Z are degenerate. We did a single Z sweep for the hexapod, setting $X = Y = 0$. We found that $Z = -12$ gives the lowest δx for this X–Y combination. Because we have no a priori knowledge on whether it is best to use X or Z, we pick the DoF that needs the smallest displacement with respect to the home configuration, which is Z in this case.

For the third scan, we scan the Y and W DoFs, setting $Z = -12$ mm and $X = 0$ mm. We center Y at 0 mm and W at 0 deg. The results are shown in Fig. 4(c). These DoFs are degenerate; therefore, we do not expect to see a single Y–W combination producing the smallest response. Rather, we expect a range of Y–W combinations to produce a small δx , creating a cold valley of minimal δx values. We set $W = 2$ deg and $Y = 0$ mm to see how a rotational degree of freedom

performs. However, from Fig. 4(c), it is clear that any Y – W combination lying inside the cold valley is a valid configuration. Had we picked $Y = -15$, W would have been set to 0 deg.

For the final scan, we verify whether $X = 0$ and $Z = -12$ is indeed a good choice. We do this by scanning Z and V . In this way, we see if $X = 0$ and $Z = -12$ gives the lowest δx for $V = 0$. If so, we can be sure that the X – Z combination is good. The result of this is illustrated in Fig. 4(d).

From the characterization, we conclude that a good prescription is to first scan two non-degenerate DoFs, to see where the cold spot is located. We chose X and Y , but we showed that one can also change from X to Z as the scanning variable as these two are degenerate. Then, we perform two scans: for each nondegenerate DoF, we perform a scan with that DoF and a DoF that is degenerate. In this way, we can fine-tune two degenerate DoFs and ensure that the beam is aligned in the axis affected by the two DoFs.

After finding the hexapod configuration that corresponded to the cold spot, we measured the complex beam patterns to assess the improvement in performance due to the alignment. Using the harmonic phase and amplitude measurement technique described by Moerman et al.,²⁴ we obtained the complex beam patterns of the home setup and aligned setup, for both beams A and B , across the DESHIMA 2.0 bandwidth. By fitting a Gaussian beam to the measured beam patterns, we obtained the tilt and focus location of the Gaussian beam.^{26,31} Using the optical simulation software PyPO,³² we propagated the measured beam patterns through a model of the ASTE telescope by placing the fitted focus in the cabin focus of the Cassegrain setup at ASTE and calculated the spillover efficiency η_s on M_2 , taper efficiency η_t on the aperture subtended by M_1 , and the aperture efficiency $\eta_{ap} \approx \eta_s \eta_t$. Because we want to compare the laboratory values to the on-site values, we multiplied η_{ap} by the Ruze efficiency³³ η_ρ of the ASTE telescope to take into account surface roughness, which is given by

$$\eta_\rho = \exp\left(-\left(\frac{4\pi\rho}{\lambda}\right)^2\right), \quad (4)$$

where ρ is the root-mean-square (RMS) surface roughness of the ASTE reflectors, which is estimated to be 42 μm , and λ the wavelength of incoming radiation. We did not take into account the aperture blockage of the primary aperture due to the support struts for M_5 , as this was not taken into account for the design values. We compared the aligned η_{ap} to design values. In addition, we calculated the η_{ap} of the hexapod in the home configuration. For this calculation of η_{ap} , we used a complex beam pattern measured after removing the entire sky chopper as the beam was too misaligned and did not pass through the chopper. These results can be found in Fig. 5.

It is evident from Fig. 5 that the alignment works successfully in the laboratory. The η_{ap} is at the design level, with an average increase of a factor of 1.80 for beam A and 1.81 for beam B , both with respect to the misaligned case.

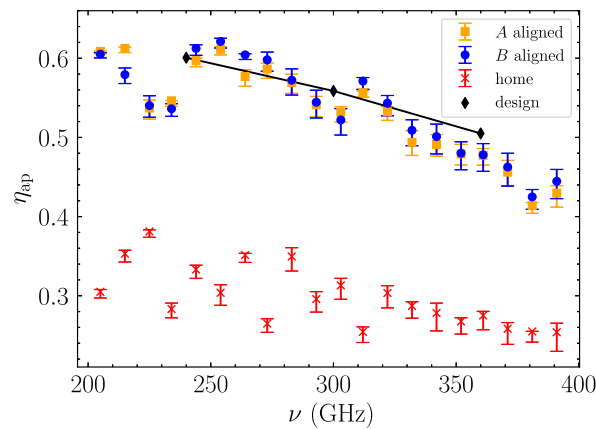


Fig. 5 Calculated η_{ap} as function of channel frequency ν in GHz for the laboratory characterization. Red crosses represent the hexapod in the home position. The orange boxes and blue dots represent the aligned hexapod position for beams A and B , respectively. Black diamonds represent design values. All η_{ap} are multiplied by the Ruze efficiency η_ρ . The errors were calculated by propagating the uncertainties in the focus location from Gaussian fitting.

3.2 Telescope Results

We report on the alignment for DESHIMA 2.0 at ASTE that was performed on November 7, 2023, between 01:15 and 04:50 UTC, fully remotely from the ASTE base camp in San Pedro de Atacama. The telescope elevation during the alignment was 60 deg, which is in the center of the elevation range of ASTE (30 to 90 deg). We used beam A to align DESHIMA 2.0. Again, we use the DESHIMA channel at 250 GHz for alignment. Using the Atacama Pathfinder EXperiment (APEX)³⁴ radiometer data, we find that the mean PWV during the alignment was ~ 5.1 mm. At 250 GHz, this corresponds to a T_{sky} of ~ 70 K. Combined with the telescope interior temperature, which is ~ 293 K, this gives a temperature contrast comparable to the temperature contrast during the laboratory characterization. Hence, we expect ample contrast for alignment at the telescope. The standard deviation of the mean PWV over the entire alignment was around 0.13 mm, which had a negligible effect on the T_{sky} variation at 250 GHz during the alignment. We also note that the PWV during the alignment was high compared with median values recorded for the Chajnantor plateau,³⁵ where ASTE is situated. As T_{sky} increases monotonically with PWV, a lower PWV enhances the temperature contrast and hence improves the alignment quality.

We first repeated the series of scans taken in Sec. 3.1 to verify the usability of this method at ASTE. The results can be found in Fig. 6. We have repeated the first two X – Y scans to localize the cold spots, as per Sec. 3.1. We did not perform a scan along the Z DoF. The Y – W scan was taken with $X = 5$ mm and centered at $Y = 4$ mm. The final X – V scan was to check if the previously found parameters indeed create a minimum δx value. The difference in δx contrast

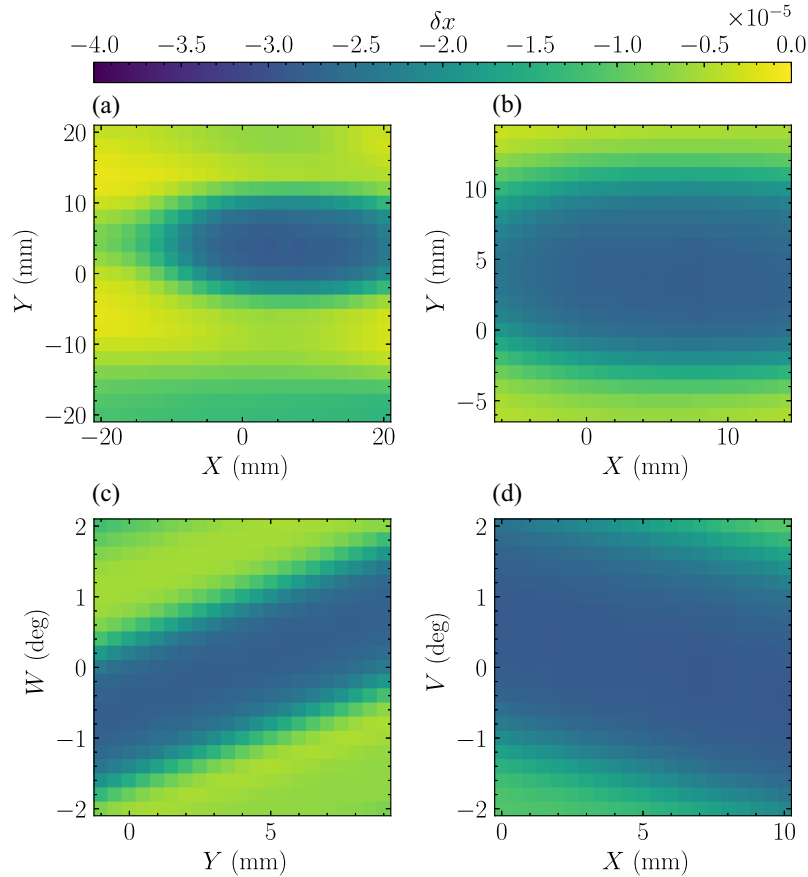


Fig. 6 Instrument responses δx as a function of hexapod scanning DoFs, for the four hexapod scans used for the measurements at ASTE. Shown responses are at a channel frequency $\nu = 250$ GHz. The scans were taken at a telescope elevation of 60 deg. Again, we extended the color bar to $\delta x = -4 \cdot 10^{-5}$ to enhance contrast inside the cold spots. (a) The X – Y scan centered at $X = Y = 0$ mm. (b) The X – Y scan centered at $X = 5$ and $Y = 4$ mm. (c) The Y – W scan centered at $Y = 4$ mm and $W = 0$ deg. (d) The X – V scan centered at $X = 5$ mm and $V = 0$ deg.

between Figs. 4 and 6 is due to the difference in temperatures of the warm surroundings and cold source in the laboratory and telescope setups.

To verify that the ISS beam is indeed coupled to the atmosphere, we performed four observations of the atmosphere without moving the telescope. Two observations are taken with the hexapod in the home position and two with the hexapod in the aligned position. For each hexapod position, one observation is performed using beam A, the other using beam B. We use the skydip method described by Takekoshi et al.²⁹ to convert δx to brightness temperature T_b . From the APEX radiometer data, we obtained the average PWV during these observations, which turned out to be 1.22 mm, which we then used to calculate η_{atm} using the Atmospheric Transmission at Microwaves (ATM) software.³⁶ Then, we calculated the theoretical sky brightness temperature during each observation using Eq. (1). A comparison between the measured brightness temperatures T_b in the home and aligned hexapod positions and calculated T_{sky} is shown in Fig. 7.

From Fig. 7, we can see that the measured T_b with the aligned hexapod is consistent with the calculated T_{sky} , indicating that the instrument beam couples to the atmosphere but also that the calibration method described by Takekoshi et al.²⁹ is applicable to a wide-band spectrometer. We find that for lower frequencies, below ~ 325 GHz, the alignment lowers T_b for both beams A and B by ~ 20 K compared with the home position, which indicates that the alignment results in less warm spillover and hence a lower background level. Around 325 GHz, there is a strong water emission line that dominates T_b for both the aligned and home positions. Between 325 and 370 GHz, the alignment again shows a lower T_b , indicating that also at these frequencies, the alignment decreases the warm spillover. Above this frequency, the difference in T_b approaches zero again, as around 370 GHz, another strong water emission line is present.

As a final stability check, we have also performed the X–Y scan at an elevation of 90 and 30 deg and found the same shape of the scan, as seen in Fig. 6(b). We have also calculated the

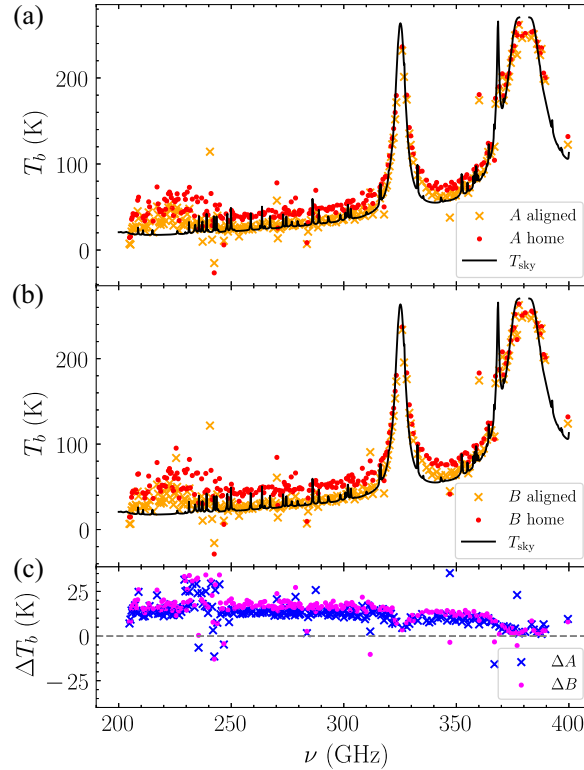


Fig. 7 Observations of the atmosphere with beams A and B, for both the home and aligned hexapod position. (a) Measured T_b in beam A as a function of filter frequency ν for both the home and aligned hexapod configuration. The black line is the theoretical atmosphere T_{sky} calculated using ATM,³⁶ assuming a PWV of 1.22 mm and an elevation of 60 deg. (b) Same as panel (a) but then for beam B. (c) The differences between T_b measured in the home and aligned position, for both beam A and B. For reference, the grey, horizontal dashed line is plotted at $\Delta T_b = 0$ K.

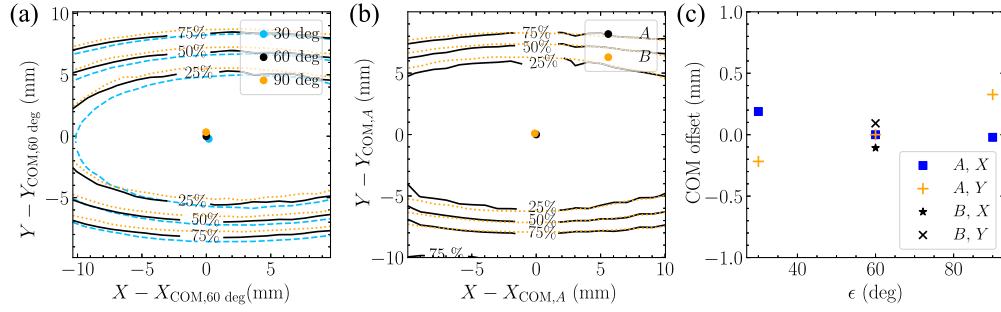


Fig. 8 Several hexapod X–Y scans to assess alignment stability for different telescope elevations and sky chopper beams. Each scan is normalized such that the minimum δx is at 0 and the maximum δx at 1. Contours for each scan are shown at 25%, 50%, and 75% levels. (a) X–Y hexapod scan contours and center of mass (COM) for three different telescope elevation angles. The COM positions are denoted by the dots. The X and Y coordinates of the scans are shifted to the COM of the scan at 60 deg. (b) X–Y hexapod scan contours and COM for beam A and B, both at an elevation of 60 deg. To faithfully compare beam A and B, the scans were taken shortly after each other and therefore the beam A scan at an elevation of 60 deg in panel (a) is different from the one here, explaining the slightly different cold spot shapes. The X and Y coordinates of the scans are shifted to the COM of the scan taken in beam A. (c) COM offsets along the X and Y coordinates with respect to the COM of the scan in beam A at an elevation angle of 60 deg.

center of mass (COM) for each scan by first normalizing the scan so that the maximum δx is set to 1 and the minimum to 0. Then, we changed the sign of the normalized scan and added 1 so that the maximum δx is at 0 and the minimum at 1. From there, we calculated the COM and found that the COM of the cold spot does not significantly move for different elevations. This indicates that the alignment is stable over an elevation angle and could be performed at any elevation angle allowed by the telescope housing the instrument. In addition, we have performed the X–Y scan using beam B at an elevation of 60 deg. We found again a similar shape to a beam A scan at an elevation of 60 deg and no significant movement of the COM. The results for the different elevation scans and the beam B comparison can be found in Fig. 8.

3.3 Measured Aperture Efficiency

To calculate the aperture efficiency η_{ap} , we took raster scans of Mars with ASTE, remotely from San Pedro de Atacama using the COSMOS3³⁷ telescope control software. Separate scans were taken for both beam A and B. The scans were taken on July 17, 2024. We determined the Martian apparent angular diameter using NASA JPL’s On-Line Solar System Data Service³⁸ and adopted $\theta_{\text{Mars}} = 5.62$ arcsec. For visualization, the scans at $\nu = 300$ GHz can be found in Fig. 9.

Visually, the two scans look very similar, indicating that the alignment has achieved similar quality for beam A and B.

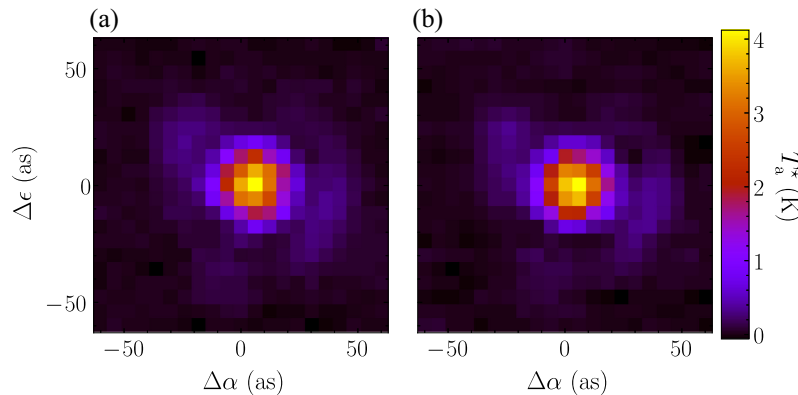


Fig. 9 Beam maps for beam A and B at 300 GHz taken on Mars. The azimuth and elevation offsets for both panels are with respect to the center of the scan. (a) Beam map taken in beam A. (b) Beam map taken in beam B.

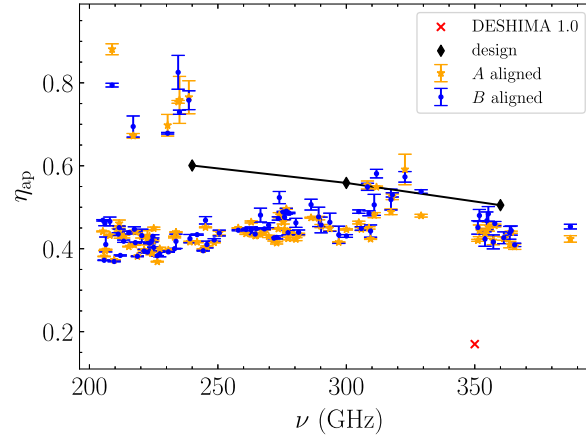


Fig. 10 Measured η_{ap} as a function of frequency, obtained from observations of Mars. The red cross denotes the η_{ap} measured during the DESHIMA 1.0 campaign in 2017, which was 0.17 at 350 GHz. The error bars denote the variance in the measured η_{ap} . We only show η_{ap} for which the signal-to-noise ratio is larger than 3.

We then used the “Mars brightness model” by Lellouch and Amri³⁹ to calculate the Martian disk-averaged brightness temperature as a function of frequency. Together with the raster scans, η_{ap} was calculated for both beam A and B. The results can be found in Fig. 10. We calculated the error in η_{ap} , σ_{ap} from the errors of the fitted parameters in the fitting procedure. In Fig. 10, we only show values for which $\eta_{\text{ap}}/\sigma_{\text{ap}} > 3$.

In Fig. 10, it is visible that the measured η_{ap} for beam A and B are consistent, indicating that the alignment works for both beams. In addition, the measured efficiency is larger than the measured η_{ap} for DESHIMA 1.0 during the 2017 campaign, which was 0.17 at 350 GHz.² Using the alignment procedure described in this work, we have increased η_{ap} by a factor of 2.5 at 350 GHz. It should be noted that ASTE has not received any upgrades or improvements to the reflector surfaces between the DESHIMA 1.0 and DESHIMA 2.0 campaigns, implying that the surface RMS has not changed. The main difference between the optical systems of DESHIMA 1.0 and DESHIMA 2.0 at ASTE, excluding the sky-position chopper, is the upgrade from a double-slot antenna⁴⁰ to a leaky-lens antenna⁵ and a modification of the cold optics to accommodate the change of antenna.²⁷ This change increased the bandwidth of DESHIMA 2.0 compared with DESHIMA 1.0 but did not change the illumination pattern and edge taper on the subreflector, and hence η_{ap} , of DESHIMA 2.0 compared with DESHIMA 1.0 at 350 GHz.

Most of the measured η_{ap} are consistently lower than the laboratory values, especially for lower frequencies. The cause of this is a lower than expected T_{a}^* , which we suspect is due to an increased sidelobe level. The laboratory values were calculated without accounting for the sub-reflector support struts, which are known to cause an elevated sidelobe level.^{41,42} Error beams arising from the surface roughness of ASTE could also lower the aperture efficiency. They have been proven to be a significant source for the pick-up of the IRAM telescope⁴³ and could also be significant for ASTE. Methods that take error beams into account exist⁴⁴ but were not considered given the scope of this work. Another cause for elevated sidelobes could be a slightly elevated edge taper on the ASTE secondary.

We confirmed the elevated sidelobe level by observing the Moon, which is of sufficient angular diameter to fill the sidelobes of the DESHIMA beam pattern and any error beams. This observation was compared with a brightness temperature model of the Moon given by Mangum⁴⁵ and showed that T_{a}^* at 240, 300, and 360 GHz was recovered to a level within 95% of the expected value. In Fig. 11, we show this result, together with the beam A observation of Mars, which shows the lower than expected T_{a}^* .

Conversely, there are six channels between 210 and 240 GHz that have η_{ap} values about twice as large as the rest of the η_{ap} in this range. Upon inspection, we found that the fitted T_{a}^* for these channels was about twice as large as the surrounding channels, explaining the discrepancy. We also checked the central and maximal values of the Mars maps directly and found

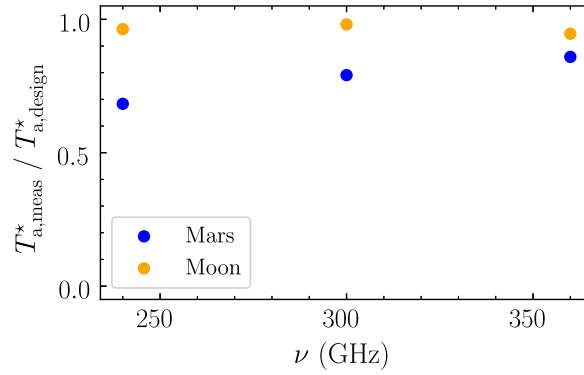


Fig. 11 Comparison of the measured $T_{a,meas}^*$ to the expected $T_{a,design}^*$ for both the Moon and Mars beam A observations. For the Moon, we used the brightness temperature model from Mangum⁴⁵ and assumed this to be equal to $T_{a,design}^*$. For Mars, we used the brightness temperature model by Lellouch and Amri³⁹ and the design η_{ap} and DESHIMA beam sizes to calculate $T_{a,design}^*$.

the same discrepancy, which indicated that the fitting process was not the cause. In addition, we inspected the off-source atmospheric spectrum from the outskirts of the Mars maps and found elevated T_a^* values for these channels. As the atmosphere is large enough to fill the entire DESHIMA beam at all channel frequencies, these observed T_a^* are independent of η_{ap} and hence could indicate an issue in the responsivity calibration for these particular channels.

4 Conclusion

We have developed an alignment method for the DESHIMA 2.0 spectrometer at ASTE, which uses the contrast of the cold atmosphere against the warm interior of the sky chopper together with a modified Dragonian dual reflector and mechanical hexapod, and can be remotely performed over the internet.

Before we applied the alignment method to DESHIMA 2.0 at ASTE, we characterized the alignment method in the laboratory, using the same DESHIMA 2.0 chip as in Chile. We placed an absorber, immersed in a vat containing liquid N₂, behind the sky chopper to emulate the cold sky. The laboratory measurements show that the alignment method successfully recovers the designed performance, starting from a misaligned system.

The alignment was then applied to the optical chain of DESHIMA 2.0 at ASTE. We showed that the characterization and alignment procedure, which we tested in the laboratory, is identical at the telescope, using the atmosphere instead of N₂. This proves that the atmosphere is a suitable cold background against the sky chopper interior and can be used for the described alignment method. In addition, we showed observations of the atmosphere performed with DESHIMA 2.0 at ASTE, compared these to the theoretical atmosphere brightness temperature curve, and found that the alignment causes less warm spillover for both beams A and B compared with the hexapod home position. As a final test, we measured the aperture efficiency using Mars and showed that we have increased the efficiency by a factor of 2.5 compared with DESHIMA 1.0 during its 2017 campaign. These observations provide evidence for the successful alignment using our method. However, we find that the measured efficiency did not reach design values, and we suspect a raised sidelobe level or error beams to cause this.

The alignment method can in principle be applied to any single-pixel submm spectrometer. In this work, we used the sky chopper, which acted as a cylinder with a higher temperature than the atmosphere. The role of the cylinder can be played by any cylinder-like object, provided that its brightness temperature is sufficiently higher than that of the atmosphere.

5 Appendix: Determination of Aperture Efficiency from Planet Map

To obtain the aperture efficiency η_{ap} from a planet map, a set of parameters needs to be extracted. To obtain these map parameters, a disk-convolved 2D Gaussian fit was performed to each

channel map. Throughout this explanation, we will denote the corrected antenna temperature map $T_a^*(\alpha, \epsilon)$, where α and ϵ represent the horizontal (azimuth) and vertical (elevation) coordinates, respectively. In our case, we bin the full time-ordered data (TOD) in azimuth and elevation pixels of a size of 6 by 6 arcsec. Each pixel in $T_a^*(\alpha, \epsilon)$ has an associated standard deviation denoted $\sigma_a(\alpha, \epsilon)$. The conversion of measured antenna temperature T_a to T_a^* was carried out using the chopper wheel calibration method.⁴⁶ As initial parameters for the Gaussian to be fitted, we used the statistical moments of the map as described by Burger and Burge.⁴⁷ The raw moments M_{ij} and central moments μ_{ij} are given by

$$M_{ij} = \sum_{\alpha} \sum_{\epsilon} \alpha^i \epsilon^j T_a^*(\alpha, \epsilon), \quad (5)$$

$$\mu_{ij} = \sum_{\alpha} \sum_{\epsilon} (\alpha - \bar{\alpha})^i (\epsilon - \bar{\epsilon})^j T_a^*(\alpha, \epsilon), \quad (6)$$

where $i, j \in \mathbb{N}$ and $\bar{\alpha}$ and $\bar{\epsilon}$ are the central α and ϵ values

$$(\bar{\alpha}, \bar{\epsilon}) = \left(\frac{M_{10}}{M_{00}}, \frac{M_{01}}{M_{00}} \right). \quad (7)$$

We estimate the floor level $T_{a,\text{floor}}^*$ by taking the average of the off-source sections of the $T_a^*(\alpha, \epsilon)$ map. The amplitude $T_{a,\text{amp}}^*$ is then estimated by taking the maximum of $T_a^*(\alpha, \epsilon) - T_{a,\text{floor}}^*$.

The position angle Ψ of the 2D Gaussian, with respect to the positive α -axis, is given by

$$\Psi = \frac{1}{2} \arctan \left(\frac{2\mu_{11}}{\mu_{20} - \mu_{02}} \right), \quad (8)$$

and the eccentricity e is given by

$$e = \frac{(\mu_{20} - \mu_{02})^2 + 4\mu_{11}^2}{(\mu_{20} + \mu_{02})^2}. \quad (9)$$

For the semi-minor axis w_y , we use the diffraction-limited full-width-half-maximum (FWHM) for cosine-tapered illumination,⁴⁸ calculated using

$$w_y \approx 1.2 \frac{\lambda}{2R_{\text{tel}}}, \quad (10)$$

where λ is the wavelength of the radiation at which the map was taken and R_{tel} is the radius of the primary aperture of the telescope. For ASTE, $R_{\text{tel}} = 5$ m. Because we define the semi-major axis w_x to lie along the x -axis (without rotation applied), we calculate the initial estimate to be

$$w_x = w_y (1 - e^2)^{-1/2}. \quad (11)$$

We then generate a Gaussian $\mathcal{G}(\alpha, \epsilon)$ using the following expression:

$$a = \frac{\cos^2 \Psi}{2\sigma_x^2} + \frac{\sin^2 \Psi}{2\sigma_y^2}, \quad (12)$$

$$b = -\frac{\sin \Psi \cos \Psi}{2\sigma_x^2} + \frac{\sin \Psi \cos \Psi}{2\sigma_y^2}, \quad (13)$$

$$c = \frac{\sin^2 \Psi}{2\sigma_x^2} + \frac{\cos^2 \Psi}{2\sigma_y^2}, \quad (14)$$

$$\mathcal{G}(\alpha, \epsilon) = \exp(-(a(\alpha - \bar{\alpha})^2 + b(\alpha - \bar{\alpha})(\epsilon - \bar{\epsilon}) + c(\epsilon - \bar{\epsilon})^2)). \quad (15)$$

Note that we use the standard deviation in Eq. (12) instead of the FWHM. We can convert from some FWHM w to standard deviation σ using

$$\sigma = \frac{w}{2\sqrt{2 \ln 2}}. \quad (16)$$

We then define a disk $\mathcal{D}(\alpha, \epsilon)$ with angular diameter θ equal to the apparent angular diameter of the planet that is scanned. Each point on the disk is assigned a brightness temperature T_b , which we take as the disk-averaged brightness temperature for the planet in question

$$\mathcal{D}(\alpha, \epsilon) = \begin{cases} T_b & \text{if } \sqrt{\alpha^2 + \epsilon^2} \leq \theta, \\ 0 & \text{otherwise.} \end{cases} \quad (17)$$

The Gaussian is normalized and convolved with the disk to obtain the disk-convolved Gaussian $\mathcal{GD}(\alpha, \epsilon)$

$$\mathcal{GD}(\alpha, \epsilon) = \frac{\sum_{\alpha'} \sum_{\epsilon'} \mathcal{D}(\alpha', \epsilon') \mathcal{G}(\alpha - \alpha', \epsilon - \epsilon') \Delta \alpha' \Delta \epsilon'}{\sum_{\alpha} \sum_{\epsilon} \mathcal{G}(\alpha, \epsilon) \Delta \alpha \Delta \epsilon}. \quad (18)$$

For the initial estimate for the fit, we set $T_b = 1$, normalize $\mathcal{GD}(\alpha, \epsilon)$ to unit amplitude, multiply by $T_{a, \text{amp}}^*$, and add $T_{a, \text{floor}}^*$. For the subsequent fitting procedure, however, we do not normalize to unit amplitude and just add $T_{a, \text{floor}}^*$ to $\mathcal{GD}(\alpha, \epsilon)$. We weigh each point in $T_a^*(\alpha, \epsilon)$ by $\sigma_a^{-2}(\alpha, \epsilon)$.

Using the `curve_fit` function in the `scipy`⁴⁹ package, we fit $\mathcal{GD}(\alpha, \epsilon)$ to $T_a^*(\alpha, \epsilon)$. The free parameters are T_b , w_y , $\bar{\alpha}$, $\bar{\epsilon}$, $T_{a, \text{floor}}^*$, e , and Ψ . After fitting, the relevant parameters for the aperture efficiency calculation are denoted \hat{T}_b , \hat{e} and \hat{w}_y , from which we can calculate \hat{w}_x using Eq. (11). With \hat{T}_b and a brightness temperature model $T_{b, \text{model}}$ of the observed planet, we calculate the main beam efficiency

$$\eta_{\text{mb}} = \frac{\hat{T}_b}{T_{b, \text{model}}}. \quad (19)$$

From \hat{w}_x and \hat{w}_y , we calculate the main beam solid angle

$$\Omega_{\text{mb}} = \frac{\pi \hat{w}_x \hat{w}_y}{4 \ln 2}, \quad (20)$$

and using η_{mb} and Ω_{mb} , we finally calculate the aperture efficiency

$$\eta_{\text{ap}} = \frac{\lambda^2 \eta_{\text{mb}}}{A_p \Omega_{\text{mb}}}, \quad (21)$$

where $A_p = \pi R_{\text{tel}}^2$ is the physical area of the primary aperture.

Disclosures

The authors declare that there are no financial interests, commercial affiliations, or other potential conflicts of interest that could have influenced the objectivity of this research or the writing of this paper.

Code and Data Availability

The processed data and code used for generating the figures in this paper are available on Zenodo: <https://zenodo.org/records/15229117>. The raw, unprocessed data are ~31.5 GB in size and are not publicly available due to their large size. The raw data are available upon request from the corresponding author.

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Arend Moerman received his MSc degree in astronomy and astrophysics from Leiden University in 2022. He is currently with Delft University of Technology, where he is pursuing a doctoral degree. His current research interest is focused on galaxy cluster dynamics and instrumentation for the mm/submm wavelength range.

Kenichi Karatsu is an Instrument Scientist with SRON. He received his PhD in science in 2011 from Kyoto University with the study of proton spin structure at the RHIC-PHENIX experiment. In 2015, he joined the DESHIMA project at TU Delft/SRON. His main role is to lead laboratory evaluation and telescope campaign of the instrument. His research interest is to develop an experimental instrument for revealing mysteries of the Universe.

Jochem J. A. Baselmans is full professor in experimental astronomy at the Delft University of Technology, senior instrument scientist at the SRON Netherlands Institute for Space Research and visiting professor at the University of Cologne. He is an expert in the development of instrumentation for far-infrared astronomy based upon Microwave Kinetic Inductance Detectors. He invented the concept of the on-chip spectrometer for far-infrared radiation and received an ERC-COG to develop this technology. He is the lead system engineer for the first on-chip spectrometer instrument, DESHIMA, fielded on the 10m ASTE telescope in 2019 and its follow-up instrument

DESHIMA 2.0. He is the detector lead for the new AMKID camera on APEX and has pushed MKID technology to sensitivity levels suitable for future cryogenically cooled space-based observatories.

Shahab O. Dabironezare received his PhD in electrical engineering from Delft University of Technology (TU Delft), Delft, The Netherlands, in 2020. He is currently an assistant professor in the THz Sensing Group at TU Delft and an instrument scientist at the Netherlands Institute for Space Research (SRON). His research interests include wideband antennas at millimeter and submillimeter-wave applications, wide field-of-view imaging systems, quasi-optical systems, lens antennas, and absorber-based passive cameras.

Shinji Fujita is a project assistant professor at The Institute of Statistical Mathematics, Japan. His research focuses on molecular clouds and star formation in the Milky Way Galaxy, using radio telescopes such as ASTE and the Nobeyama 45m telescope. He employs data science approaches including machine learning and deep learning to analyze large-scale astronomical data.

Robert Huiting is a mechanical design engineer. He received his BS degree in precision engineering (2001) and technology management (2003) at the Hogeschool van Utrecht. He has been working as a design engineer for SRON Netherlands Institute for Space Research since 2007. He has worked as a research engineer at the FOM Institute for Plasma Physics for the XUV optics group in the years before. His experience from the last 13 years lies in designing instrumentation for ground-based space research.

Kotaro Kohno is a professor and director at the Institute of Astronomy, the University of Tokyo. He received his PhD in 1998 for studies of dense molecular gas in Seyfert and starburst galaxies. His research focuses on spatially resolved studies of galaxies using ALMA and JWST, and on surveys at submm/mm/radio wavelengths using techniques such as gravitational lensing, radio-optical crossmatching, and line intensity mapping. He is also actively involved in instrumentation development for these studies.

Yuri Nishimura is a project assistant professor at the Institute of Astronomy, the University of Tokyo, focusing on integrated superconducting spectrometer development and analysis. She received her PhD in physics from the University of Tokyo in 2017. Her research interests include the physics and chemistry of the interstellar medium, particularly on scales ranging from molecular clouds to entire galaxies, in the context of galaxy evolution.

Fenno Steenvoorde is an instrument maker at DEMO, TU Delft. He graduated from the Leidse instrumentmakers School in 2020, and worked for SRON Netherlands Institute for Space Research during 2020–2022. As a member of the DESHIMA team, he developed cryogenic and optical components, and participated in the installation on the ASTE telescope in 2023. His current interest includes the development of the cryogenic components for TIFUUN.

Tatsuya Takekoshi is an assistant professor at Kitami Institute of Technology, specializing in submillimeter and millimeter wave instrumentation for astronomy. His research focuses on star formation and observational studies using the ASTE telescope. He is an expert in optical and mechanical design for astronomical instrumentation, contributing to the development of advanced observing systems, including cryogenic and superconducting detector technologies.

Yoichi Tamura is a professor of astronomy at Nagoya University. He received his PhD from the University of Tokyo and has worked for the National Astronomical Observatory of Japan and the University of Tokyo before joining Nagoya University. He brings a 20-year experience in millimeter and submillimeter astronomy and instrumentation. His main interests include a conceptual design of a future large-aperture submillimeter telescope, commissioning, and signal processing of an integrated superconductor spectrometer, and a coherent receiver system to existing millimeter and submillimeter telescopes. These things facilitate his astronomical interests in galaxy formation in the early Universe.

Akio Taniguchi is a project assistant professor at Kitami Institute of Technology. He received his PhD from the University of Tokyo in 2018 and was a postdoc and a project assistant professor at Nagoya University before joining Kitami Institute of Technology. His research interests focus

on the development of efficient spectroscopy for next-generation submillimeter single-dish telescopes based on data scientific approaches. He is contributing to the development of data analysis software for the submillimeter spectroscopic instruments such as DESHIMA, to the data analysis methods implemented on top of them, and to the commissioning and science verification using them.

Stephen J. C. Yates is an instrument scientist at SRON, the Netherlands Institute for Space Research (since 2006). He received his PhD from the University of Bristol (2003), followed by work at the CNRS-CRTBT (now Institut Néel) Grenoble on experimental low-temperature physics (2003 to 2006). His current interests concentrate on MKID development for submillimeter to FIR astronomical applications but also include a wider interest in device physics and superconductivity, optics, and full end-to-end instrument characterization and performance. He has published over 80 papers.

Bernhard R. Brandl is a professor of infrared astronomy, at Leiden University and Technical University Delft, who works on starburst galaxies, HII regions, and astronomical instrumentation. He has authored over 130 papers in refereed astronomical journals and over 150 technical publications, with more than 15,000 citations. He has been involved in the design and construction of the JWST-MIRI spectrometer, and he is the principal investigator of the first-generation instrument METIS for ESO's Extremely Large Telescope.

Akira Endo is an experimental astronomer and associate professor at TU Delft. He is interested in mapping the large-scale Universe in 3D, and the required development of (sub)millimeter wave imaging spectrometers. He leads the DESHIMA collaboration as the Dutch principal investigator, and currently leads an ERC Consolidator grant collaboration to develop TIFUUN: an integral field spectrometer based on plug-and-play, open-hardware integral field units (IFUs) that are individually designed to address specific astronomical questions.