

Uncovering the limits of uniqueness in sampled Gabor phase retrieval: A dense set of counterexamples in $L_2(\mathbb{R})$

Alaifari, Rima; Bartolucci, Francesca; Wellershoff, Matthias

DOI

[10.1109/SampTA59647.2023.10301382](https://doi.org/10.1109/SampTA59647.2023.10301382)

Publication date

2023

Document Version

Final published version

Published in

Proceedings of the 2023 International Conference on Sampling Theory and Applications (SampTA)

Citation (APA)

Alaifari, R., Bartolucci, F., & Wellershoff, M. (2023). Uncovering the limits of uniqueness in sampled Gabor phase retrieval: A dense set of counterexamples in $L_2(\mathbb{R})$. In *Proceedings of the 2023 International Conference on Sampling Theory and Applications (SampTA)* IEEE.
<https://doi.org/10.1109/SampTA59647.2023.10301382>

Important note

To cite this publication, please use the final published version (if applicable).
Please check the document version above.

Copyright

Other than for strictly personal use, it is not permitted to download, forward or distribute the text or part of it, without the consent of the author(s) and/or copyright holder(s), unless the work is under an open content license such as Creative Commons.

Takedown policy

Please contact us and provide details if you believe this document breaches copyrights.
We will remove access to the work immediately and investigate your claim.

Green Open Access added to TU Delft Institutional Repository

'You share, we take care!' - Taverne project

<https://www.openaccess.nl/en/you-share-we-take-care>

Otherwise as indicated in the copyright section: the publisher is the copyright holder of this work and the author uses the Dutch legislation to make this work public.

Uncovering the limits of uniqueness in sampled Gabor phase retrieval: A dense set of counterexamples in $L^2(\mathbb{R})$

Rima Alaifari
ETH Zürich
Department of Mathematics
Email: rima.alaifari@sam.math.ethz.ch

Francesca Bartolucci
TU Delft
Institute of Applied Mathematics
Email: f.bartolucci@tudelft.nl

Matthias Wellershoff
University of Maryland
Department of Mathematics
Email: wellersm@umd.edu

Abstract—Sampled Gabor phase retrieval — the problem of recovering a square-integrable signal from the magnitude of its Gabor transform sampled on a lattice — is a fundamental problem in signal processing, with important applications in areas such as imaging and audio processing. Recently, a classification of square-integrable signals which are not phase retrievable from Gabor measurements on parallel lines has been presented. This classification was used to exhibit a family of counterexamples to uniqueness in sampled Gabor phase retrieval. Here, we show that the set of counterexamples to uniqueness in sampled Gabor phase retrieval is dense in $L^2(\mathbb{R})$, but is not equal to the whole of $L^2(\mathbb{R})$ in general. Overall, our work contributes to a better understanding of the fundamental limits of sampled Gabor phase retrieval.

Index Terms—Phase retrieval, Gabor transform, sampling result

I. INTRODUCTION

Phase retrieval is a term broadly applied to problems in which information about complex phase needs to be inferred from data. The origins of phase retrieval can be traced back to the years 1915 – 1929 when W. H. Bragg and W. L. Bragg (among others) used X-ray diffraction images of crystals in order to illuminate their atomic structure [1], [2], [3], [4], [5]. Since then phase retrieval has found applications in various fields such as crystallography, electron microscopy and astronomy [6], [7], [8].

Gabor phase retrieval refers to problems in which signals $f \in L^2(\mathbb{R})$ have to be reconstructed from magnitudes of their Gabor transform,

$$\mathcal{G}f(x, \omega) := 2^{1/4} \int_{\mathbb{R}} f(t) e^{-\pi(t-x)^2} e^{-2\pi i t \omega} dt, \quad (x, \omega) \in \mathbb{R}^2.$$

It has been used in a range of audio processing tasks, including the phase vocoder for time-stretching and pitch-shifting of audio signals [9], as well as speech enhancement and source separation [10].

In this contribution, we will specifically focus on the *sampled Gabor phase retrieval* problem which is the recovery of signals f from the magnitude measurements

$(|\mathcal{G}f(x, \omega)|)_{(x, \omega) \in \Lambda}$ where $\Lambda \subset \mathbb{R}^2$ is a lattice¹. We focus on this sampling problem because magnitude information on the entire time-frequency plane $\mathbb{R}^2$ is not available in practice. Instead, only a finite number of measurements are stored and inferences are made based on them. We consider the sampled setup proposed above as a natural and useful compromise between the fully continuous case, where no sampling occurs, and the fully discrete case, where the signals are finite-dimensional vectors.

A. Prior arts: Counterexamples to uniqueness in sampled Gabor phase retrieval

In the following, we will focus on counterexamples to uniqueness in sampled Gabor phase retrieval; i.e. signals whose Gabor transform magnitudes agree on a lattice but which are fundamentally different from each other. Before introducing the concept of a counterexample rigorously, we need to emphasise that there is one ever-present ambiguity in Gabor phase retrieval: the global phase ambiguity. Two signals $f, g \in L^2(\mathbb{R})$ are said to *agree up to global phase* if they are equivalent with respect to the relation

$$f \sim g : \iff f = e^{i\alpha} g, \text{ for some } \alpha \in \mathbb{R}.$$

Notably, $f \sim g$ implies $|\mathcal{G}f| = |\mathcal{G}g|$ such that signals which agree up to global phase cannot be recovered from Gabor transform magnitudes. With this in mind, we define the set of counterexamples.

Definition 1.1 (Counterexamples). *Let $\Lambda \subset \mathbb{R}^2$. The set of counterexamples to uniqueness in Gabor phase retrieval on Λ is defined by*

$$\mathfrak{C}(\Lambda) := \{f \in L^2(\mathbb{R}) \mid |\mathcal{G}f| = |\mathcal{G}g| \text{ on } \Lambda \text{ and } f \not\sim g, \text{ for some } g \in L^2(\mathbb{R})\}.$$

An element $f \in \mathfrak{C}(\Lambda)$ is called a counterexample to uniqueness in sampled Gabor phase retrieval on Λ .

¹A lattice $\Lambda \subset \mathbb{R}^2$ is a discrete subset of the time-frequency plane that can be written as $L\mathbb{Z}^k$ where $L \in \mathbb{R}^{2 \times k}$ is a matrix with linearly independent columns and $k \in \{1, 2\}$.

Counterexamples to uniqueness in sampled Gabor phase retrieval are interesting for two reasons. First, they allow us to better understand the fundamental limits of sampled Gabor phase retrieval, which in turn can guide future research towards achieving uniqueness. In addition, they offer the opportunity to explore the potential relationship between uniqueness and stability in phase retrieval [11].

Let us briefly summarise the recent research on counterexamples in Gabor phase retrieval. The relationship between the Gabor transform and the Bargmann transform (which is described in more detail in Section II) allows for the relation of the Gabor phase retrieval problem to a phase retrieval problem for entire functions. This was realised in [12], [13]. Then, following these ideas, a characterisation of all entire functions of exponential-type whose magnitudes agree on any set of infinitely many equidistant parallel lines was proven in [14].

Using this characterisation and noting that all lattices are a subset of some set of infinitely many equidistant parallel lines, it becomes possible to construct various types of counterexamples. This idea has been applied in [15] to construct explicit counterexamples to uniqueness in sampled Gabor phase retrieval on any lattice. (See [14] for a more in-depth explanation.) An extension of the results in [15] has appeared in [16].

B. Our contributions

In this contribution, we show that the set of counterexamples $\mathfrak{C}(\Lambda)$ is dense in $L^2(\mathbb{R})$ when $\Lambda \subset \mathbb{R}^2$ is a lattice or a set of equidistant parallel lines. We also show that the Gaussian is *not* a counterexample for quadratic lattices, $\Lambda = a\mathbb{Z}^2$, with $a \in (0, 1)$. Therefore, *the set of counterexamples is dense but not equal to the whole of $L^2(\mathbb{R})$ in general.* We prove these two results by using the connection between the Bargmann transform and the Gabor transform as well as some classical results from complex analysis.

Note that this contribution is a condensed and modified version of the section on the fragility of uniqueness in sampled Gabor phase retrieval in the larger manuscript [11]. Apart from a comprehensive treatment of counterexamples, the larger manuscript also discusses the stability of Gabor phase retrieval as well as its potential connection with uniqueness in sampled Gabor phase retrieval. Here, we focus specifically on showing that the counterexamples are dense.

Notation

Rotation by $\theta \in \mathbb{R}$ on $\mathbb{R}^2$ is denoted by $R_\theta : \mathbb{R}^2 \rightarrow \mathbb{R}^2$; in matrix notation, we have

$$R_\theta = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}.$$

Translation by $x \in \mathbb{R}$ on $L^p(\mathbb{R})$, where $p \in [1, \infty]$, is denoted by $T_x : L^p(\mathbb{R}) \rightarrow L^p(\mathbb{R})$; i.e.

$$T_x f(t) = f(t - x), \quad t \in \mathbb{R},$$

for $f \in L^p(\mathbb{R})$. Finally, the normalised Gaussian is denoted by

$$\varphi(t) = 2^{1/4} e^{-\pi t^2}, \quad t \in \mathbb{R}.$$

II. THE RELATION BETWEEN THE BARGMANN AND GABOR TRANSFORM

As mentioned before, we will make use of the well-known connection between the Bargmann transform and the Gabor transform [17]. The *Fock space* $\mathcal{F}^2(\mathbb{C})$ is the Hilbert space of all entire functions for which the norm induced by the inner product

$$(F, G)_{\mathcal{F}} := \int_{\mathbb{C}} F(z) \overline{G(z)} e^{-\pi |z|^2} dz$$

is finite. The *Bargmann transform* $\mathcal{B} : L^2(\mathbb{R}) \rightarrow \mathcal{F}^2(\mathbb{C})$,

$$\mathcal{B}f(z) := 2^{1/4} \int_{\mathbb{R}} f(t) e^{2\pi i t z - \pi t^2 - \frac{\pi}{2} z^2} dt, \quad z \in \mathbb{C},$$

acts as an isomorphism between $L^2(\mathbb{R})$ and $\mathcal{F}^2(\mathbb{C})$. It is related to the Gabor transform by the formula

$$\mathcal{G}f(x, -\omega) = e^{\pi i x \omega} \mathcal{B}f(x + i\omega) e^{-\frac{\pi}{2}(x^2 + \omega^2)}, \quad (1)$$

for $(x, \omega) \in \mathbb{R}^2$. It is the formula above that allows us to relate the Gabor phase retrieval problem to a phase retrieval problem for entire functions. We are going to use this relation in the sequel.

III. THE SET OF COUNTEREXAMPLES IS DENSE

In this section, we show that the set of counterexamples $\mathfrak{C}(\Lambda)$ is dense in $L^2(\mathbb{R})$ when $\Lambda \subset \mathbb{R}^2$ is a lattice or a set of equidistant parallel lines. Our strategy will be to design entire functions $H_\delta^\pm \in \mathcal{F}^2(\mathbb{C})$ which converge to a constant as $\delta \rightarrow 0$, do not agree up to global phase and still satisfy $|H_\delta^+| = |H_\delta^-|$ on $\mathbb{R} + ia\mathbb{Z}$. Once these entire functions have been designed, we can multiply them with $\mathcal{B}f$ and take the inverse Bargmann transform in order to get counterexamples on $\mathbb{R} \times a\mathbb{Z}$ that are close to $f \in L^2(\mathbb{R})$.

We design the functions $H_\delta^\pm \in \mathcal{F}^2(\mathbb{C})$ by modifying the counterexamples on $\mathbb{R} \times a\mathbb{Z}$,

$$h^\pm(t) := \varphi(t) \left(\cosh\left(\frac{\pi t}{a}\right) \pm i \sinh\left(\frac{\pi t}{a}\right) \right), \quad t \in \mathbb{R},$$

presented in [15]. To accomplish this, we compute the Bargmann transforms of $h^\pm$ which are given by

$$z \mapsto (1 \mp i + (1 \pm i)e^{\frac{\pi z}{a}}) e^{-\frac{\pi z}{2a}}$$

up to a constant depending on a . Next, we note that time-shifting $h^\pm$ by $u \in \mathbb{R}$ will produce additional counterexamples on $\mathbb{R} \times a\mathbb{Z}$ according to the covariance property of the Gabor transform [17]. For $f \in L^2(\mathbb{R})$, it holds that

$$\mathcal{B}T_u f(z) = \mathcal{B}f(z - u) e^{\pi i u z} e^{-\frac{\pi}{2} u^2}, \quad z \in \mathbb{C},$$

such that the Bargmann transforms of the counterexamples $T_u h^\pm$, with $u = -\frac{a}{\pi} \log \delta$, are given by

$$z \mapsto (1 \mp i + (1 \pm i)\delta \cdot e^{\frac{\pi z}{a}}) \delta^{-az} e^{-\frac{\pi z}{2a}}$$

up to a constant depending on a and δ . After multiplying by $\delta^{az} e^{\frac{\pi z}{2a}}$, it follows that

$$z \mapsto 1 \mp i + (1 \pm i)\delta \cdot e^{\frac{\pi z}{a}}$$

are entire functions whose magnitudes agree on $\mathbb{R} + ia\mathbb{Z}$ (which do not agree up to global phase) and we define

$$H_\delta^\pm(z) := 1 \pm i\delta \cdot e^{\frac{\pi z}{a}}, \quad z \in \mathbb{C}, \quad (2)$$

after multiplying by $(1 \pm i)/2$.

Theorem III.1. *Let $a > 0$. Then, $\mathfrak{C}(\mathbb{R} \times a\mathbb{Z})$ is dense in $L^2(\mathbb{R})$.*

Proof. Let $\epsilon > 0$ and $f \in L^2(\mathbb{R})$. We want to show that there exist $g^\pm \in L^2(\mathbb{R})$ which do not agree up to global phase, are ϵ -close to f in $L^2(\mathbb{R})$ and satisfy

$$|\mathcal{G}g^+| = |\mathcal{G}g^-| \text{ on } \mathbb{R} \times a\mathbb{Z}.$$

To do so, we note that the monomials

$$e_n(z) := \left(\frac{\pi^n}{n!}\right)^{1/2} z^n, \quad n \in \mathbb{N}_0, \quad z \in \mathbb{C},$$

form an orthonormal basis for the Fock space $\mathcal{F}^2(\mathbb{C})$ [17]. Therefore, the space of complex polynomials is dense in $\mathcal{F}^2(\mathbb{C})$ and we can find $P \in \mathbb{C}[z]$ such that

$$\|\mathcal{B}f - P\|_{\mathcal{F}} < \frac{\epsilon}{2}.$$

Let us now consider the entire functions $H_\delta^\pm$ defined in equation (2) and note that $G_\delta^\pm := H_\delta^\pm \cdot P \in \mathcal{F}^2(\mathbb{C})$ since $G_\delta^\pm$ are entire functions of exponential-type. Hence, we can define the signals $g_\delta^\pm := \mathcal{B}^{-1}G_\delta^\pm \in L^2(\mathbb{R})$. To establish the desired properties of $g_\delta^\pm$, we will work with their Bargmann transforms $G_\delta^\pm$. First, we note that $|H_\delta^+| = |H_\delta^-|$ on $\mathbb{R} + ia\mathbb{Z}$ implies $|G_\delta^+| = |G_\delta^-|$ on $\mathbb{R} + ia\mathbb{Z}$ and thus $|\mathcal{G}g_\delta^+| = |\mathcal{G}g_\delta^-|$ on $\mathbb{R} \times a\mathbb{Z}$ by equation (1). Secondly, we note that the entire functions $G_\delta^\pm$ do not agree up to global phase: indeed, both entire functions $H_\delta^\pm$ have infinitely many roots but no root of H_δ^+ is a root of H_δ^- and vice versa. At the same time, P is a polynomial and has only finitely many roots. It follows that G_δ^+ does have roots which are not roots of G_δ^- (and vice versa) and thus $G_\delta^+ \not\sim G_\delta^-$. By the linearity of the Bargmann transform, we can conclude that $g_\delta^+ \not\sim g_\delta^-$. Finally, we note that the definition of $H_\delta^\pm$ in equation (2) directly implies that

$$\|P - P \cdot H_\delta^\pm\|_{\mathcal{F}} = \delta \|z \mapsto P(z) \cdot e^{\pi z/a}\|_{\mathcal{F}},$$

and so there exists a $\delta > 0$ depending on a , ϵ and P (which in turn depends on f and ϵ) such that

$$\|P - P \cdot H_\delta^\pm\|_{\mathcal{F}} < \frac{\epsilon}{2}.$$

We conclude

$$\begin{aligned} \|f - g_\delta^\pm\|_2 &= \|\mathcal{B}f - H_\delta^\pm \cdot P\|_{\mathcal{F}} \\ &\leq \|\mathcal{B}f - P\|_{\mathcal{F}} + \|P - H_\delta^\pm \cdot P\|_{\mathcal{F}} < \epsilon. \end{aligned}$$

□

Remark III.2 (Some explanations on the proof). As $\mathcal{B}f \in \mathcal{F}^2(\mathbb{C})$, for $f \in L^2(\mathbb{R})$, we know that $\mathcal{B}f$ is either an entire function of exponential-type or an entire function of second order. If $\mathcal{B}f$ is of second order, then its type is either strictly

smaller than $\pi/2$ or exactly $\pi/2$. In most of these cases, it holds that $\mathcal{B}f \cdot H_\delta^\pm \in \mathcal{F}^2(\mathbb{C})$ and thus we can define

$$g_\delta^\pm := \mathcal{B}^{-1}(\mathcal{B}f \cdot H_\delta^\pm) \in L^2(\mathbb{R}),$$

with

$$\delta < \frac{\epsilon}{\|z \mapsto \mathcal{B}f(z)e^{\pi z/a}\|_{\mathcal{F}}},$$

to obtain counterexamples which are ϵ -close to f in $L^2(\mathbb{R})$. We note that $g_\delta^\pm$ are small additive perturbations of our original signals f .

Unfortunately, there is one case in which this simple strategy does not work: the one in which $\mathcal{B}f$ is a second-order entire function of type $\pi/2$. Indeed, in this case, it is not guaranteed that $\mathcal{B}f \cdot H_\delta^\pm$ is in the Fock space. — Two striking examples for why this can fail can be found in [18]. — As the only situation in which $\mathcal{B}f \cdot H_\delta^\pm$ is not in the Fock space occurs when $\mathcal{B}f$ is exactly of order two and of type $\pi/2$, it seems obvious that the functions f for which $\mathcal{B}f \cdot H_\delta^\pm \in \mathcal{F}^2(\mathbb{C})$ holds must be dense in $L^2(\mathbb{R})$. We can prove this by realising that the complex polynomials are dense in $\mathcal{F}^2(\mathbb{C})$.

Theorem III.1 continues to hold for any set of infinitely many equidistant parallel lines. We can show this by considering the entire functions

$$H_\delta^\pm(z) := 1 \pm i\delta \exp\left(\frac{\pi e^{i\theta}}{a}(z - \bar{\lambda}_0)\right)$$

and realising that the corresponding signals $g_\delta^\pm \in L^2(\mathbb{R})$ satisfy

$$|\mathcal{G}g_\delta^+| = |\mathcal{G}g_\delta^-| \text{ on } \mathbb{R}_\theta(\mathbb{R} \times a\mathbb{Z}) + \lambda_0,$$

where $a > 0$, $\lambda_0 \in \mathbb{R}^2 \simeq \mathbb{C}$. The statement for general lattices follows from the same consideration because all lattices are subsets of some set of infinitely many equidistant parallel lines. We therefore arrive at the following result.

Theorem III.3. *Let $\Lambda \subset \mathbb{R}^2$ be a set of equidistant parallel lines or a lattice. Then, $\mathfrak{C}(\Lambda)$ is dense in $L^2(\mathbb{R})$.*

To illustrate our main results, we construct counterexamples that are close to the Hermite functions and plot their spectrograms.

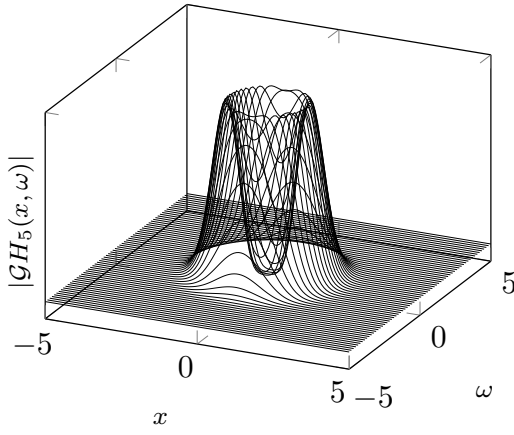
Example III.4. *Consider the n -th Hermite function $H_n \in L^2(\mathbb{R})$ given by*

$$\mathcal{B}H_n(z) = e_n(z) = \left(\frac{\pi^n}{n!}\right)^{1/2} z^n, \quad z \in \mathbb{C}.$$

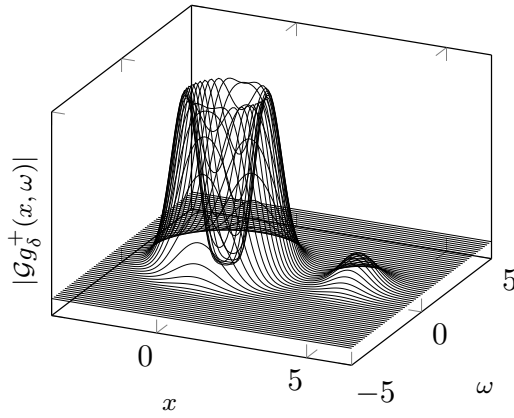
By equation (1), the Gabor transform of the Hermite function is

$$\begin{aligned} \mathcal{G}H_n(x, \omega) &= e^{-\pi i x \omega} \mathcal{B}H_n(x - i\omega) e^{-\frac{\pi}{2}(x^2 + \omega^2)} \\ &= \left(\frac{\pi^n}{n!}\right)^{1/2} e^{-\pi i x \omega} (x - i\omega)^n e^{-\frac{\pi}{2}(x^2 + \omega^2)}, \end{aligned}$$

for $(x, \omega) \in \mathbb{R}^2$. If we plot the magnitude of the above (for $n = 5$), we obtain Figure 1a. Next, we want to find counterexamples which are close to H_n . According to Remark III.2, we can



(a) $|\mathcal{G}H_5|$



(b) $|\mathcal{G}g_\delta^+|$, for $\delta = \frac{1}{50} \exp(-10\pi)$.

Fig. 1. The Gabor magnitude of the fifth Hermite function (Figure 1a) and of a counterexample g_δ^+ to sampled Gabor phase retrieval on $\mathbb{R} \times \frac{1}{4}\mathbb{Z}$ (Figure 1b).

define $g_\delta^+ := \mathcal{B}^{-1}(\mathcal{B}H_n \cdot H_\delta^+)$. Let us visualise the spectrogram of g_δ^+ , i.e.

$$\mathcal{G}g_\delta^+(x, \omega) = \mathcal{G}H_n(x, \omega) \cdot H_\delta^+(x - i\omega), \quad (x, \omega) \in \mathbb{R}^2,$$

in Figure 1b (for $n = 5$, $a = \frac{1}{4}$ and $\delta = \frac{1}{50} \exp(-10\pi)$).

IV. THE GAUSSIAN IS NOT A COUNTEREXAMPLE

Finally, we can show that the Gaussian is not a counterexample on $a\mathbb{Z}^2$ if $a \in (0, 1)$. Specifically, we prove the following result.

Theorem IV.1. *Let $0 < a < 1$ and $f \in L^2(\mathbb{R})$ be such that*

$$|\mathcal{G}f(x, \omega)|^2 = |\mathcal{G}\varphi(x, \omega)|^2, \quad (x, \omega) \in a\mathbb{Z}^2.$$

Then, there exists an $\alpha \in \mathbb{R}$ such that $f = e^{i\alpha}\varphi$.

Since the Bargmann transform of the Gaussian is one, equation (1) implies that the theorem above is equivalent to the following lemma.

Lemma IV.2. *Let $0 < a < 1$ and let $F \in \mathcal{F}^2(\mathbb{C})$ be such that*

$$|F(z)| = 1 = |\mathcal{B}\varphi(z)|, \quad z \in a\mathbb{Z} + ia\mathbb{Z}.$$

Then, there exists an $\alpha \in \mathbb{R}$ such that $F = e^{i\alpha}$.

The intuition for the proof of this lemma comes from the maximum modulus principle: we note that we are considering a second order entire function F which is bounded on all lattice points; this suggests that F should be constant in the entire complex plane as long as the lattice is dense enough. This intuition is indeed correct, as evidenced by the following result independently discovered by V. G. Iyer [19] and A. Pfluger [20].

Theorem IV.3. *Let h be an entire function such that*

$$\limsup_{r \rightarrow \infty} \frac{\log M_h(r)}{r^2} < \frac{\pi}{2},$$

where $M_h(r) := \max_{|z|=r} |h(z)|$. If there exists a constant $\kappa > 0$ such that

$$|h(m + in)| \leq \kappa, \quad m, n \in \mathbb{Z},$$

then h is constant.

Proof of Lemma IV.2. Consider the function $h(z) := F(az)$, for $z \in \mathbb{C}$. It holds that

$$|h(z)| = |F(az)| \leq \|F\|_{\mathcal{F}} \cdot e^{\frac{\pi}{2}|az|^2} = \|F\|_{\mathcal{F}} \cdot e^{\frac{\pi a^2}{2}|z|^2},$$

for $z \in \mathbb{C}$, such that

$$\begin{aligned} \limsup_{r \rightarrow \infty} \frac{\log M_h(r)}{r^2} &\leq \limsup_{r \rightarrow \infty} \left(\frac{\log \|F\|_{\mathcal{F}}}{r^2} + \frac{\pi a^2}{2} \right) \\ &= \frac{\pi a^2}{2} < \frac{\pi}{2}. \end{aligned}$$

Additionally,

$$|h(m + in)| = |F(am + ian)| = 1, \quad m, n \in \mathbb{Z},$$

holds such that the assumptions of Theorem IV.3 are met and we can conclude that h is constant. As $|h(0)| = 1$, it follows that there must exist an $\alpha \in \mathbb{R}$ such that $h = e^{i\alpha}$ which implies $F = e^{i\alpha}$. $\square$

We have therefore shown that the set of counterexamples is not equal to the whole of $L^2(\mathbb{R})$ when Λ is a sufficiently dense quadratic lattice.

Remark IV.4. *A natural confusion that might arise in connection with Theorem IV.1 is in how far it is different from the result in [21] on shift-invariant spaces with Gaussian generator $V_\beta^1(\varphi)$. While Theorem IV.1 implies that the Gaussian can be distinguished from all other functions in $L^2(\mathbb{R})$ based on its sampled Gabor magnitude measurements, the result in [21] only implies that it can be distinguished from the functions in $V_\beta^1(\varphi) \subset L^2(\mathbb{R})$.*

ACKNOWLEDGEMENTS

The authors would like to extend their heartfelt thanks to Stefan Steinerberger for his insightful discussions and acknowledge funding through the SNSF grant 200021_184698.

REFERENCES

- [1] W. H. Bragg, "Bakerian lecture: X-rays and crystal structure," *Philosophical Transactions of the Royal Society of London. Series A, Containing Papers of a Mathematical or Physical Character*, vol. 215, pp. 253–274, 1915, <https://www.jstor.org/stable/91108>.
- [2] W. L. Bragg, "The determination of parameters in crystal structures by means of Fourier series," *Proceedings of the Royal Society of London. Series A, Containing Papers of a Mathematical and Physical Character*, vol. 123, no. 792, pp. 537–559, 1929, <https://doi.org/10.1098/rspa.1929.0083>.
- [3] J. M. Robertson, "X-ray analysis and application of Fourier series methods to molecular structures," *Reports on Progress in Physics*, vol. 4, no. 1, pp. 332–367, January 1937, <https://doi.org/10.1088/0034-4885/4/1/324>.
- [4] W. L. Bragg, "The Rutherford memorial lecture, 1960 — the development of X-ray analysis," *Proceedings of the Royal Society of London. Series A. Mathematical and Physical Sciences*, vol. 262, no. 1309, pp. 145–158, 1961, <https://doi.org/10.1098/rspa.1961.0109>.
- [5] D. Sayre, "X-ray crystallography: The past and present of the phase problem," *Structural Chemistry*, vol. 13, no. 1, pp. 81–96, February 2002, <https://doi.org/10.1023/A:1013477415486>.
- [6] J. R. Fienup, "Phase retrieval algorithms: A comparison," *Applied Optics*, vol. 21, no. 15, pp. 2758–2769, August 1982, <https://doi.org/10.1364/AO.21.002758>.
- [7] J. C. Dainty and J. R. Fienup, *Phase retrieval and image reconstruction for astronomy*, 1st ed. Academic Press, March 1987, ch. 7, pp. 231–275.
- [8] Y. Shechtman, Y. C. Eldar, O. Cohen, H. N. Chapman, J. Miao, and M. Segev, "Phase retrieval with application to optical imaging: A contemporary overview," *IEEE Signal Processing Magazine*, vol. 32, no. 3, pp. 87–109, May 2015, <https://doi.org/10.1109/MSP.2014.2352673>.
- [9] Z. Průša and N. Holighaus, "Phase vocoder done right," in *2017 25th European Signal Processing Conference (EUSIPCO)*. Kos, Greece: IEEE, August 2017, pp. 976–980, <https://doi.org/10.23919/EUSIPCO.2017.8081353>.
- [10] R. Abdelmalek, Z. Mnasri, and F. Benzarti, "Audio signal reconstruction using phase retrieval: Implementation and evaluation," *Multimedia Tools and Applications*, vol. 81, no. 11, pp. 15 919–15 946, May 2022, <https://doi.org/10.1007/s11042-022-12421-1>.
- [11] R. Alaifari, F. Bartolucci, S. Steinerberger, and M. Wellershoff, "On the connection between uniqueness from samples and stability in Gabor phase retrieval," arXiv:2205.07013 [math.FA], May 2022, <https://doi.org/10.48550/arXiv.2205.07013>.
- [12] J. N. Mc Donald, "Phase retrieval and magnitude retrieval of entire functions," *Journal of Fourier Analysis and Applications*, vol. 10, no. 3, pp. 259–267, May 2004, <https://doi.org/10.1007/s00041-004-0973-9>.
- [13] P. Jaming, "Uniqueness results in an extension of Pauli's phase retrieval problem," *Applied and Computational Harmonic Analysis*, vol. 37, no. 3, pp. 413–441, November 2014, <https://doi.org/10.1016/j.acha.2014.01.003>.
- [14] M. Wellershoff, "Phase retrieval of entire functions and its implications for Gabor phase retrieval," arXiv:2202.03733 [math.CV], February 2022, <https://doi.org/10.48550/arXiv.2202.03733>.
- [15] R. Alaifari and M. Wellershoff, "Phase retrieval from sampled Gabor transform magnitudes: Counterexamples," *Journal of Fourier Analysis and Applications*, vol. 28, no. 1, pp. 1–9, December 2021, <https://doi.org/10.1007/s00041-021-09901-7>.
- [16] P. Grohs and L. Liehr, "On foundational discretization barriers in STFT phase retrieval," *Journal of Fourier Analysis and Applications*, vol. 28, no. 2, pp. 1–39, April 2022, <https://doi.org/10.1007/s00041-022-09935-5>.
- [17] K. Gröchenig, *Foundations of Time-Frequency Analysis*, 1st ed., ser. Applied and Numerical Harmonic Analysis (ANHA). Boston, MA: Birkhäuser, December 2000, <https://doi.org/10.1007/978-1-4612-0003-1>.
- [18] C. Bénéteau, B. J. Carswell, and S. Kouckekian, "Extremal problems in the Fock space," *Computational Methods and Function Theory*, vol. 10, no. 1, pp. 189–206, June 2010, <https://doi.org/10.1007/BF03321762>.
- [19] V. G. Iyer, "A note on integral functions of order 2 bounded at the lattice points," *Journal of the London Mathematical Society*, vol. s1-11, no. 4, pp. 247–249, October 1936, <https://doi.org/10.1112/jlms/s1-11.4.247>.
- [20] A. Pfluger, "On analytic functions bounded at the lattice points," *Proceedings of the London Mathematical Society*, vol. s2-42, no. 1, pp. 305–315, 1937, <https://doi.org/10.1112/plms/s2-42.1.305>.
- [21] P. Grohs and L. Liehr, "Injectivity of Gabor phase retrieval from lattice measurements," *Applied and Computational Harmonic Analysis*, vol. 62, pp. 173–193, January 2023, <https://doi.org/10.1016/j.acha.2022.09.001>.