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ABSTRACT

We present the first experimental application of coherent Raman spectroscopy (CRS) on the ro-vibrational ν_2 mode spectrum of methane (CH_4). Ultrabroadband femtosecond/picosecond (fs/ps) CRS is performed in the molecular fingerprint region from 1100 to 2000 cm^{-1} , employing fs laser-induced filamentation as the supercontinuum generation mechanism to provide the ultrabroadband excitation pulses. We introduce a time-domain model of the CH_4 ν_2 CRS spectrum, including all five ro-vibrational branches allowed by the selection rules $\Delta v = 1$, $\Delta J = 0, \pm 1, \pm 2$; the model includes collisional linewidths, computed according to a modified exponential gap scaling law and validated experimentally. The use of ultrabroadband CRS for *in situ* monitoring of the CH_4 chemistry is demonstrated in a laboratory CH_4 /air diffusion flame: CRS measurements in the fingerprint region, performed across the laminar flame front, allow the simultaneous detection of molecular oxygen (O_2), carbon dioxide (CO_2), and molecular hydrogen (H_2), along with CH_4 . Fundamental physicochemical processes, such as H_2 production via CH_4 pyrolysis, are observed through the Raman spectra of these chemical species. In addition, we demonstrate ro-vibrational CH_4 ν_2 CRS thermometry, and we validate it against CO_2 CRS measurements. The present technique offers an interesting diagnostics approach to *in situ* measurement of CH_4 -rich environments, e.g., in plasma reactors for CH_4 pyrolysis and H_2 production.

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I. INTRODUCTION

Methane (CH_4) is the simplest hydrocarbon molecule and one of the most abundant chemical species in the universe, with its ubiquitous chemistry having important implications in planetary^{1,2} and life³ science. In recent years, the challenges posed by anthropogenic climate change have highlighted the critical role played by CH_4 in the oil and gas industry and in combustion technologies. CH_4 is one of the most powerful greenhouse gases and, despite being short-lived in the Earth's atmosphere, it has the second largest effective climate forcing potential after carbon dioxide (CO_2):⁴ CH_4

is the major component of natural gas and it is associated with all hydrocarbon fuels. On the other hand, natural gas (compressed or liquefied) is the main alternative fuel for internal combustion engines,⁵ reducing the emission of transport vehicles with respect to typical diesel and gas oil engines, and CH_4 can be produced in a sustainable fashion as a bio-⁶ and e-fuel.⁷ In this respect, CH_4 can also be used as a chemical energy storage to enable the development of the hydrogen economy:⁷ CH_4 can be transported within the existing infrastructure for natural gas and valorized via chemical reforming to molecular hydrogen⁸ (H_2) or other commodity hydrocarbons.^{9,10}

In view of this, CH₄ has been extensively investigated by laser diagnostics with linear optical techniques, based on absorption or scattering, providing *in situ* detection and ro-vibrational spectroscopy in energy systems. Optical probes based on laser absorption have been developed using, e.g., near-¹¹ and mid-infrared (IR)¹² diode lasers, quantum-cascade lasers,¹³ and frequency combs,¹⁴ and have been vastly employed for CH₄ detection in atmospheric environments¹⁵ and for combustion diagnostics.¹⁶ Absorption techniques have the advantage of simplicity and robustness, as well as the availability of large spectroscopic databases for the interpretation of experimental spectra,¹⁷ they can realize ultrafast spectroscopy employing femtosecond (fs) laser sources,¹⁸ and the capability for imaging measurements has also been demonstrated.^{19,20} On the other hand, they suffer from the lack of spatial resolution in the longitudinal direction as they are line-of-sight techniques,¹⁶ which can severely limit the accuracy of quantitative spectroscopy in inhomogeneous gas-phase media, such as turbulent combustion environments.

The use of spontaneous Raman scattering in these contexts allows for spatially resolved measurements, with the possibility for imaging and high repetition rates.²² Spontaneous Raman spectroscopy has been thoroughly applied to the study of CH₄,^{23,24} both in combustion^{25–29} and plasma^{21,30} environments, demonstrating its suitability as a probe molecule for the direct measurement of the rotational and vibrational temperature even in non-equilibrium systems, with important implications for CH₄ chemistry.³¹ The main drawback of spontaneous Raman is the typically low cross section for this incoherent scattering process, which complicates its application in luminous environments, requiring large laser pulse energies to achieve single-shot measurements, with the risk of inducing the optical breakdown of the gas-phase medium or molecular photofragmentation.²⁷ In addition, further incoherent emission processes (e.g., chemiluminescence and fluorescence) can shadow the Raman signal, and the use of nanosecond (ns) laser pulses limits the temporal resolution so that spontaneous Raman spectroscopy cannot resolve ultrafast dynamics such as energy redistribution and relaxation processes.³²

The limitations due to the incoherent nature of the Raman scattering process can be overcome by resorting to non-linear optics, particularly in the form of stimulated (SRS) and coherent Raman scattering (CRS). Despite the different nature of their emission,³³ SRS and CRS share the use of the driving pump and Stokes laser fields with frequency differences matching the ro-vibrational transitions of the Raman-active molecules, resulting in the emission of a coherent signal, which can be remotely detected even in harsh environments. Since the very beginning of their development, continuous-wave SRS^{34,35} and CRS^{36–39} have been employed to perform high-resolution measurements of the ro-vibrational Raman spectrum of CH₄ in the so-called pentad region, allowing for the assignment of individual rotational lines in the isotropic Q-branch of the symmetric C–H stretch [ν_1 mode, see Fig. 1(a)]. They were further employed to investigate the rotational energy transfer in molecular collisions^{40–42} between the vibrationally coherent CH₄ molecules and other collisional partners, e.g., N₂ and argon (Ar). In the context of combustion diagnostics, the CH₄ CRS signal was recorded at temperatures as high as 1273 K in a furnace,⁴³ and CH₄ ns CRS was used to measure the temperature in a

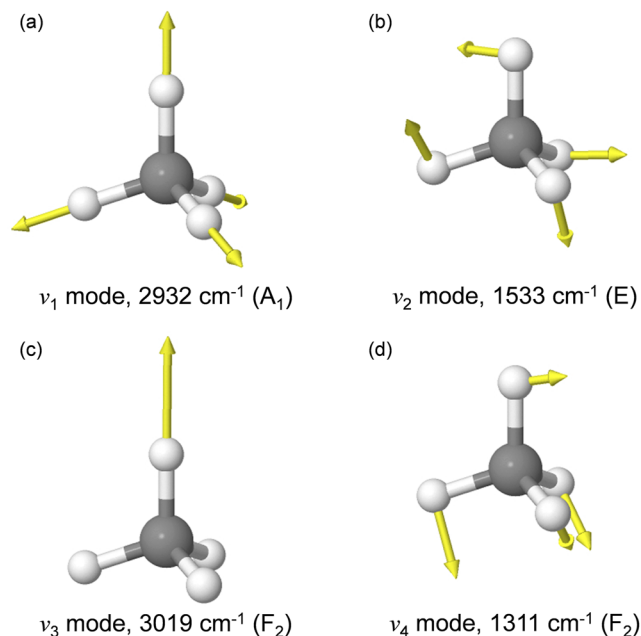


FIG. 1. The normal vibrational modes of methane.²¹ (a) Symmetric C–H stretch mode. (b) Doubly degenerate H–C–H bend mode. (c) Triply degenerate asymmetric C–H stretch mode. (d) Triply degenerate H–C–H bend mode (Jmol: an open-source Java viewer for chemical structures in 3D. <http://www.jmol.org/>).

laminar CH₄/air diffusion flame⁴⁴ and in supercritical liquid oxygen (LOX)/CH₄ combustion.⁴⁵

In recent years, the introduction of ultrafast laser sources, providing pulses of the duration of pico- (ps) and femtosecond (fs), has led to the development of various kinds of time-resolved CRS techniques, able to directly measure the rotational and vibrational wave packets of Raman-active molecules.⁴⁶ In particular, hybrid fs/ps CRS,⁴⁷ employing a combination of broadband fs pump and Stokes and narrowband ps probe pulses, achieved simultaneous time and frequency resolution, measuring the CRS spectrum below the relaxation timescale of the ro-vibrational energy. Different kinds of time-resolved CRS have been applied to CH₄. Relative CH₄ concentrations measurements in binary mixtures with N₂ were demonstrated via hybrid CRS by Engel *et al.*⁴⁸ and via chirped-probe-pulse (CPP) CRS by Dennis *et al.*⁴⁹ Bohlin and Klierer at Sandia National Laboratories (SNL), Livermore, CA, employed a hollow-core fiber to compress a 45 fs pulse to <7 fs supercontinuum to excite the ro-vibrational Raman modes up to the pentad region of the CH₄ Raman spectrum, realizing ultrabroadband fs/ps CRS imaging.⁵⁰ In a later work,⁵¹ the same group employed ultrabroadband CRS imaging to investigate CH₄/air flame-wall interaction, realizing the simultaneous detection of N₂, H₂, molecular oxygen (O₂), carbon monoxide (CO), CO₂, and CH₄. Recently, researchers at SNL Livermore and at the Plasma Physics Laboratory, Princeton, NJ, have presented a thorough investigation of the ν_1 mode spectrum of by hybrid fs/ps CRS:⁵² they demonstrated the viability of CH₄ CRS thermometry up to 1000 K, realized time-resolved measurements of the collisional dephasing of the Raman coherence of CH₄ perturbed by itself, by N₂

and by Ar, and developed a time-domain model of the CH₄ ν_1 CRS spectrum.

All the applications of CRS to CH₄ have been focused on the pentad region of its spectrum (~ 2600 – 3300 cm⁻¹), which is dominated by the fundamental bands of the ν_1 mode and of the asymmetric C–H stretch (ν_3) mode [see Fig. 1(c)]. The so-called dyad region (~ 1200 – 1900 cm⁻¹), due to the doubly and triply degenerate bending of the H–C–H bonds²³ [ν_2 and ν_4 modes, in Figs. 1(b) and 1(d) respectively], has been thus far neglected as the Raman cross-section of the bending modes is more than one order of magnitude smaller than that of the stretch modes. On the other hand, this represents an important region for combustion and plasma diagnostics, as it overlaps with the ro-vibrational spectra of CO₂⁵³ and O₂,⁵⁴ as well as with the pure-rotational spectrum of H₂, thus offering an interesting window for the study of CH₄ chemistry with high spatiotemporal resolution. In view of this, we here present the first CRS investigation of the ν_2 mode of CH₄, dominating the dyad region of its Raman spectrum. We employ a two-beam hybrid fs/ps CRS system⁵⁵ to perform ultrabroadband CRS spectroscopy in the region 1100 – 1200 cm⁻¹ of the Raman spectrum using *in situ* fs-laser-induced filamentation as the compressed supercontinuum generation mechanism.^{54,56,57} The use of hybrid fs/ps allows us to simultaneously access the frequency and time-domain to investigate the coupling of the rotational and vibrational energy degrees of freedom of the CH₄ molecule, as well as the rotational energy transfer in molecular collisions. In order to demonstrate the potential of CH₄ ν_2 CRS as a diagnostic tool for combustion and plasma technologies, we performed spatially resolved multiplexed measurements across a laminar CH₄/air diffusion flame.

II. METHODS

A. Optical setup

The experimental setup for two-beam ultrabroadband fs/ps CRS is similar to the one detailed in a previous study (Fig. 2);⁵⁴ only a brief summary is here provided. A regenerative chirped-pulse amplifier (CPA) system (Astrella, Coherent) is employed as the single light source, providing fs laser pulses at 800 nm, with

a pulse energy of ~ 7.5 mJ, at a repetition rate of 1 kHz. A 65% split of the total pulse energy is temporally compressed to 35 fs and fed to a second-harmonic bandwidth compressor (SHBC, Light Conversion) to generate a frequency-doubled ~ 5 ps pulse, which serves as the narrow-band probe pulse. Its duration and spectral linewidth are tuned by a 4f-pulse shaper to 7.3 ps and 4.1 cm⁻¹ (time-bandwidth product: 0.89 for a top-hat pulse⁵⁸) with a resulting probe pulse energy of 200 μ J. The remainder of the CPA output is employed as a single degenerate fs pump/Stokes pulse in a two-beam phase CRS matching configuration:^{55,59} an external compressor unit provides independent control of the spectral phase of the fs pulse so as to compensate for the dispersion due to the optical elements on the pump/Stokes beam path. As demonstrated in previous studies,^{54,60} fs laser-induced filamentation can be used to generate a compressed supercontinuum, which acts as a single ultrabroadband pump/Stokes pulse, coherently exciting all the Raman-active modes whose characteristic period is longer than the supercontinuum pulse. Here, this *in situ* generation/*in situ* use scheme is employed to perform ultrabroadband CRS in the vibrational fingerprint region: a spherical lens ($f = 500$ mm) focuses the ~ 1.2 mJ pump/Stokes pulse, with a theoretical irradiance of ~ 357 TW/cm² at the focal point, thus resulting in the formation of a laser-induced plasma filament, which propagates for ~ 13 mm before collapsing. The laser focusing lens is mounted on a linear stage, whose position is adjusted to maintain the filament ~ 4 mm before the probe volume. As the fs pulse experiences self-phase modulation and nonlinear propagation in the plasma medium inside the filament, these result in a compressed supercontinuum output, which serves as the ultrabroadband pump/Stokes pulse.

A crossed beam quasi-phase-matched configuration is adopted with the probe beam focused by a spherical lens ($f = 300$ mm) to the measurement location and crossing the pump/Stokes at a $\sim 3^\circ$ angle, resulting in an estimated 20 μ m [width, full width at half maximum (FWHM)] $\times 2.5$ mm (length, FWHM) $\times 20$ μ m (height, FWHM) probe volume. The relative delay of the probe pulse is tuned by controlling an automated delay line on the pump/Stokes beam path with a <10 fs resolution (Thorlabs). A combination of 800 nm half-wave plate and thin-film polarizer is employed to split a ~ 50 μ J portion of the beam: a single reflection from a wedge plate is focused by a cylindrical lens ($f = 300$ mm) onto a sCMOS beam profiler (WinCamD, Dataray) so as to monitor the (most sensitive) vertical alignment of the pump/Stokes beam over the whole range of motion of the automated stage. The ultrabroadband CRS signal is collimated through a spherical lens ($f = 400$ mm) and a band-pass filter (20 nm bandwidth FWHM, Semrock) is employed to suppress the probe, almost co-propagating with the signal in this geometry. The signal is finally dispersed through a high-resolution transmission grating (3039 lines/mm, Ibsen Photonics) and imaged by a relatively fast-focusing lens ($f = 200$ mm) onto the sCMOS detector (Zyla, Andor), resulting in a detection bandwidth of 900 cm⁻¹, with a dispersion of 0.43 cm⁻¹/pixel. The 4.1 cm⁻¹ spectral linewidth of the ps probe pulse is the limiting factor for the resolution of the ultrabroadband CRS signal; the instrumental broadening due to the spectrometer is evaluated by fitting a Voigt profile with ~ 0.1 cm⁻¹ FWHM Lorentzian and ~ 1.7 cm⁻¹ FWHM Gaussian contributions. Half-wave plates are employed to maximize the efficiency of the transmission grating (measured to $>90\%$ for S-polarization at 400 nm) in the 4f-filter and to control the relative polarization

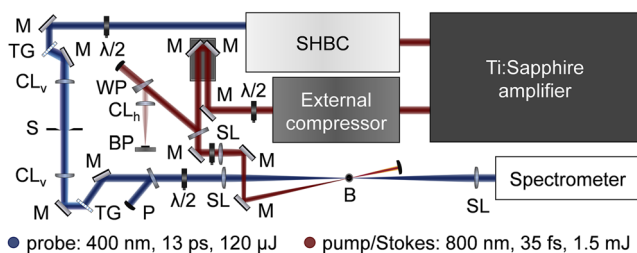


FIG. 2. Schematic of the ultrabroadband two-beam fs/ps CRS instrument. A single Ti:sapphire regenerative amplifier is employed to generate the 35 fs pump/Stokes and the 7 ps probe pulses. *In-situ* fs filamentation is employed as the supercontinuum generation mechanism: the ultrabroadband pump/Stokes beam is crossed by the probe beam at a distance of ~ 4 mm after exiting the filament. TG, transmission grating; S, slit; TS, translation stage; CL, cylindrical lens with horizontal (h) and vertical (v) symmetry axes; $\lambda/2$, half-wave plate; WP, wedge plate; BP, beam profiler; B, burner; P, polarizer; SL, spherical lens; M, mirror.

angle between the pump/Stokes and the probe field: in this work, the beams were set to have same polarization so as to maximize the signal generation efficiency, and a 400 nm half-wave plate is employed to turn the polarization of the (completely depolarized) CH₄ ν_2 CRS signal and optimize its diffraction efficiency in the spectrometer.

B. Time-resolved CRS measurements

The ν_2 mode Raman spectrum of CH₄ is here investigated under a number of different experimental conditions by varying the collisional environment to measure its dephasing coefficients and by performing ultrabroadband spectroscopy in a CH₄/air flame, demonstrating single-shot detection of the CH₄ ν_2 CRS spectrum at temperatures as high as ~800 K.

The collisional dephasing of the CH₄ ν_2 CARs signal was measured by performing time-resolved measurements of its spectrum and delaying the probe pulse relative to the pump/Stokes up to 220 ps, in steps of 5.1 ps. Samples of 1000 single-shot CH₄ ν_2 CRS spectra were recorded up to ~80–100 ps, where the collisional dephasing results in a reduction of the signal-to-background ratio (SBR) to ~3.3, and ten-shot-averaged spectra were recorded from that point onward: at the transition point, both single-shot and shot-averaged datasets were acquired and compared, computing a scaling factor over the different dynamic ranges of the dephasing experiments. Collisional dephasing measurements were performed at atmospheric pressure in different mixtures of combustion-relevant gases. The measurements were performed in an open flow, regulated by digital flow controllers (Bronkhorst), provided inside a stainless steel T-junction: two set of measurements for binary mixtures of 75% and 50% CH₄ in combination with N₂, H₂, and argon were performed. The demonstration of the single-shot detection of the CH₄ ν_2 CRS spectrum at high temperatures in a chemically reacting flow was realized in a laminar CH₄/air flame, provided on a Bunsen burner. The burner is a seamless steel pipe with an inner diameter of 19 mm through which the CH₄ flow was delivered with an exit plane velocity of 2.65 cm/s, ensuring a laminar flow (Reynolds number ~35); a steel mesh was placed ~15 mm above the burner to stabilize the flame, while the measurements were performed ~1 mm above its nozzle. Point-wise ultrabroadband CRS measurements were performed at 25 locations across the flame front, moving from the center of the burner (location $y = 0$ mm) with a 0.5 mm step size. In order to extend the detection limit at lower concentrations and higher temperatures, a set of 1000 ten-shot-averaged CRS spectra was acquired at each measurement location.

C. Time-domain CRS model

The modeling of the CH₄ ν_2 coherent Raman spectrum follows the modeling approach of the CH₄ ν_1 spectrum presented by Chen *et al.*,⁵² based on the assumption of impulsive excitation of the Raman coherence, and on a phenomenological description of the resulting third-order nonlinear optical susceptibility following Prince *et al.*⁴⁷ This approach has been successfully demonstrated for both pure-rotational^{61–63} and ro-vibrational^{54,64} CRS of multiple combustion-relevant species. The main assumptions of the time-domain model can be outlined as follows: the nonlinear interaction of the electronically off-resonant pump, Stokes probe laser fields (at angular frequencies ω_i with $i = 1, 2$, and 3 for the pump, Stokes,

and probe, respectively) with the gas-phase medium induces the macroscopic third-order polarization:

$$P^{(3)}(t, \tau_{12}, \tau_{23}) = \left(-\frac{i}{\hbar}\right)^3 \int_{-\infty}^t dt_3 \int_{-\infty}^t dt_2 \int_{-\infty}^t dt_1 \times [E_3(t-t_3) e^{i(\omega_3 - \omega_2 + \omega_1)} \times \delta(t_3) \chi^{(3)}(t_2) E_2^*(t-t_3 + \tau_{23} - t_2) e^{i(-\omega_2 + \omega_1)} \times E_1(t-t_3 + \tau_{23} - t_2 + \tau_{12} - t_1) e^{i\omega_1} \delta(t_1)], \quad (1)$$

where E_i and E_i^* are the envelope of the i th field and its complex conjugate, respectively, t_i is the coherence timescale for the molecules after the interaction with the i th field, τ_{ij} is the delay of the interaction with the i th and j th fields, δ is the Dirac function, representing the instantaneous dephasing of the (virtual) electronic coherence, and $\chi^{(3)}$ is the third-order susceptibility of the gas-phase medium. Under the assumption of impulsive excitation of the ro-vibrational Raman coherence, we can disregard the temporal envelope of the combined pump/Stokes field (in case of two-beam CRS⁵⁵) and write

$$P^{(3)}(t, \tau_{23}) = \left(-\frac{i}{\hbar}\right)^3 E_3(t + \tau_{23} - t_2) \chi^{(3)}(t_2). \quad (2)$$

Hence, the CRS intensity spectrum is obtained, upon Fourier transforming the polarization field, as

$$I_{\text{CRS}}(\omega, \tau_{23}) = \left| \mathcal{F}\{P^{(3)}(t, \tau_{23})\} \right|^2. \quad (3)$$

$\chi^{(3)}$ represents the core of the model outlined so far and is typically treated phenomenologically as

$$\chi^{(3)}(t_2) = \sum_{(v_i, J_i)} W_{(v_i, J_i) \rightarrow (v_f, J_f)} \times \exp[i(i\omega_{(v_i, J_i) \rightarrow (v_f, J_f)} - \Gamma_{(v_i, J_i) \rightarrow (v_f, J_f)})t], \quad (4)$$

with summation running over all the possible initial and final ro-vibrational states (represented by the total angular momentum and vibrational quantum numbers, v and J) allowed by the Raman selection rules for the molecular species considered. Thus, the temporal evolution of $\chi^{(3)}$ is computed as the interferogram of damped harmonic oscillations, at the Raman frequencies $\omega_{(v_i, J_i) \rightarrow (v_f, J_f)}$ and with damping coefficients $\Gamma_{(v_i, J_i) \rightarrow (v_f, J_f)}$ dominated by the collisional rotational energy transfer (RET) at atmospheric pressure,⁶⁵ weighted by the semi-classical transition probabilities.⁶⁶ Thus, Eq. (4) introduces in the time-domain CRS model the spectroscopic properties of the Raman-active species, which in turn depend on the chemical structure of the molecule considered.

The case of CH₄ is considerably more complicated than simple diatomic molecules, as early studies of its absorption spectrum revealed.⁶⁷ CH₄ is a polyatomic spherical top molecule, possessing three threefold and twofold axes of symmetry, and six planes of symmetry: in terms of these symmetry elements, CH₄ belongs to the tetrahedral point group (T_d).⁶⁸ Since the molecule has five nuclei, there are nine possible vibrational modes that, according to group theory, can be clustered in the four normal modes as illustrated in Fig. 1, called by the Greek letter “ ν ”, and identified by the vibrational quantum numbers (ν_1 , ν_2 , ν_3 , and ν_4) and the symmetry labels A, E, and F for non-degenerate, doubly degenerate, and

triply degenerate modes, respectively. The first normal mode (ν_1 at 2932.4 cm^{-1}), as already mentioned, is the non-degenerate symmetric stretch of the C–H bonds, which preserves the symmetry of the molecule in the vibrational ground state (0000) so that the instantaneous polarizability tensor has zero anisotropic invariant, and the associated Raman spectrum consists only of the isotropic Q-branch (with selection rules: $\Delta\nu = 1$, $\Delta J = 0$). The ν_2 mode (at 1533.3 cm^{-1}) is a doubly degenerate bend of the H–C–H bonds: this mode like all the other vibrational modes of CH_4 is Raman active, with selection rules $\Delta J = 0, \pm 1, \pm 2$, so that the Raman spectrum is composed of O-, P-, Q-, R-, and S-branch lines. The asymmetric C–H stretch mode (ν_3 at 3018.5 cm^{-1}) and the ν_4 bend mode (at 1310.8 cm^{-1}), being the least symmetric modes, are triply degenerate and the reduced symmetry results in the selection rules: $\Delta J = 0, \pm 1, \pm 2$. Experimental measurements on the depolarization ratio of the fundamental ν_2 and ν_3 Raman spectra have shown that these modes are completely depolarized (i.e., $\rho = 0.75$), with the Raman activity entirely due to the anisotropic part of the polarizability tensor. The normal frequencies of the stretch modes are roughly double than the bending modes (i.e., $\nu_1 \approx \nu_3 \approx 2\nu_2 \approx 2\nu_4$) so that vibrational bands can be grouped into polyads, denoted P_n , with polyad number $n = (2\nu_1 + 2\nu_3 + \nu_2 + \nu_4)$. P_0 is the vibrational ground state; P_1 has two vibrational levels, and Raman transitions $P_0 \leftarrow P_1$ involving the fundamental bands of ν_2 and ν_4 constitute the dyad region of the Raman spectrum; similarly, P_2 has five vibrational levels, and transitions $P_1 \leftarrow P_2$ involve the fundamental ν_1 and ν_3 modes, as well as the overtones ($2\nu_2$ and $2\nu_4$) and combination ($\nu_2 + \nu_4$) of the bending modes, defining the pentad region.

The concept of symmetry plays an essential role in molecular spectroscopy as it constrains the form of the molecular wave function according to Pauli's exclusion principle. Under the assumption of the Born–Oppenheimer approximation, the total wave function of the CH_4 molecule is written as the product of separable functions,

$$\Psi = \Psi_T \Psi_E \Psi_V \Psi_R \Psi_S, \quad (5)$$

where Ψ_T is the translational wave function, Ψ_E is the electronic wave function (including the electron spin), Ψ_V is the vibrational wave function, Ψ_R is the rotational wave function, and Ψ_S is the nuclear spin wave function. As the translational motion does not involve any internal degree of freedom, Ψ_T is completely symmetrical (A symmetry), and so is Ψ_E in the electronic ground state. The symmetry of the vibrational wave function depends on the symmetry of the irreducible representations of the T_d point group corresponding to the normal vibrational modes in Fig. 1. The rotational state of a CH_4 molecule can be specified by the quantum numbers J and M , where M represents the projection of the total angular momentum with respect to an inertial frame of reference: the associated wave function is a solution of the spherical top wave equation. Wilson, Jr. showed that this solution is a linear combination of irreducible representations of symmetry A, E, and F.⁶⁹ his results are summarized in Table I. The rotational wave function has $(2J + 1)$ degenerated projections with respect to an arbitrary inertial axis (i.e., $M = 0, \pm 1, \dots, \pm J$), hence the common factor in the table: this degeneracy cannot be lifted by intramolecular interaction, but only by the application of an external field (e.g., in Stark effect).⁶⁸

TABLE I. Symmetry of the rotational wave function of CH_4 (with $p = 0, 1, 2, 3, 4$, or 5).

$J = 6p$	$(2J + 1) [(p + 1)A + 2pE + 3pF]$
$J = 6p + 1$	$(2J + 1) [pA + 2pE + 3(3p + 1)F]$
$J = 6p + 2$	$(2J + 1) [pA + 2(p + 1)E + 3(3p + 1)F]$
$J = 6p + 3$	$(2J + 1) [(p + 1)A + 2pE + 3(3p + 2)F]$
$J = 6p + 4$	$(2J + 1) [(p + 1)A + 2(3p + 1)E + 3(3p + 2)F]$
$J = 6p + 5$	$(2J + 1) [pA + 2(p + 1)E + 3(3p + 3)F]$

Finally, CH_4 has 16 possible nuclear spin wave functions: two are singlet states (with the nuclear spin quantum number $I = 0$, and associated projection $M_I = 0$), three are triplet states ($I = 1$, $M_I = 0, \pm 1$), and one is a quintet state ($I = 2$, $M_I = 0, \pm 1, \pm 2$). In terms of symmetry, Ψ_S is the combination:⁶⁹ $5A + E + 3F$. The CH_4 molecule has four identical ^1H nuclei, which obey the Fermi–Dirac statistics: according to Pauli's principle, the total wave function must be symmetric with respect to any proper rotation of the T_d point group so that its overall symmetry must be A.

Under the multiplication in Eq. (5), the symmetry label obeys the rules in Table II, and only wave functions with resulting symmetry A are allowed. This has important implications for the statistical weights of the ro-vibrational energy levels included in the calculation of Boltzmann distribution

$$\rho = g_V g_R g_S \frac{\exp(-hcE_{v,J}/k_B T)}{Z}, \quad (6)$$

where $E_{v,J}$ is the ro-vibrational energy (expressed in cm^{-1}) of the molecule (v is a short-hand notation for the complete set of vibrational quantum numbers), g_V is the vibrational degeneracy, g_R is the degeneracy of the rotational wave function according to the entries in Table I, and g_S is the nuclear spin degeneracy, being 5, 2, and 3 for A, E, and F symmetry labels, respectively. According to the symmetry multiplication rules in Table II, only the product of identical symmetry labels produces a representation with A symmetry: hence, the often statement, in the literature, that ro-vibrational states (i.e., $\Psi_V \Psi_R$) of symmetry A, E, and F are fivefold, twofold, and threefold degenerate, respectively.

In order to model the temporal evolution of the nonlinear susceptibility according to Eq. (4), one must know the frequencies of the allowed Raman transitions, and the corresponding polarizabilities of the CH_4 molecule, which enter the weighting factors as⁴⁴

$$W_{(v_i, J_i) \rightarrow (v_f, J_f)} = |\langle v_i, J_i | \hat{\alpha} | v_f, J_f \rangle|^2 (\rho_{(v_f, J_f)} - \rho_{(v_i, J_i)}). \quad (7)$$

These quantities can be computed adopting a tetrahedral formalism developed for the analysis of spherical top molecules,⁷⁰ which removes inter-polyad contributions to the higher order, thanks to

TABLE II. Multiplication table for the irreducible representations.

	A	E	F
A	$A = A$	$E \not\equiv A$	$F \not\equiv A$
E	$E \not\equiv A$	$(2A + E) \supseteq A$	$F \not\equiv A$
F	$F \not\equiv A$	$F \not\equiv A$	$(A + E + 2F) \supseteq A$

implicit contact transformations,⁷¹ while explicitly retaining interpolyad contributions due to the strong coupling between vibrational states within individual polyads. This approach was developed at the University of Burgundy and allowed the derivation of effective ro-vibrational Hamiltonians for polyads up to the tetradecad (i.e., P_4 , at $\sim 6200\text{ cm}^{-1}$), from which the position and intensities of the spontaneous Raman spectrum of CH_4 are computed.^{72–74} The resulting spectral database has been employed by multiple research groups to simulate the spontaneous CH_4 Raman spectrum in both the dyad⁷⁵ and pentad region,^{21,30,76} and recently, its application has been extended to the ν_1 CH_4 CRS spectrum.⁵² According to the notation of Jourdanneau *et al.*, the spontaneous Raman line intensities in the spectral database are computed as⁷⁶

$$I_{(\nu_i J_i) \rightarrow (\nu_f J_f)} \propto \rho(\nu_i J_i) [S_{0,2} \mathcal{J}_{(\nu_i J_i) \rightarrow (\nu_f J_f)}^2 + S_{2,2} \mathcal{A}_{(\nu_i J_i) \rightarrow (\nu_f J_f)}^2], \quad (8)$$

where $S_{0,2}$ are the Stone coefficients, depending on the observation geometry in a spontaneous Raman experiment, while $\mathcal{J}$ and $\mathcal{A}$ are the transition values of the isotropic and anisotropic components of the polarizability tensor, respectively. As already mentioned, the ν_2 mode Raman spectrum is completely depolarized so that its transition polarizability has only an anisotropic component, and the corresponding Stone coefficient becomes an irrelevant scaling factor. In contrast to the incoherent spontaneous Raman scattering, the probability amplitude for the coherent Raman scattering process depends on the differential of the Boltzmann population in the initial and final ro-vibrational states, as shown in Eq. (7). Thus, the weight factors in Eq. (4) are computed from the line intensities in the spectral database as

$$W_{(\nu_i J_i) \rightarrow (\nu_f J_f)} = I_{(\nu_i J_i) \rightarrow (\nu_f J_f)} \left(\frac{\rho(\nu_f J_f) - \rho(\nu_i J_i)}{\rho(\nu_i J_i)|_{T_0}} \right), \quad (9)$$

with T_0 ($\sim 1500\text{ K}$) being the reference temperature for the calculation of the line intensities in the spectral database.²¹ The spectral database contains $\sim 16 \times 10^6$ lines in the dyad region, involving vibrational states up to the tetradecad and rotational states up to $J = 23$: of these, $\sim 11 \times 10^6$ correspond to ν_2 mode Raman transitions (i.e., $\Delta \nu_2 = 1$, $\Delta \nu_4 = 0$) and $\sim 5 \times 10^6$ are ν_4 mode transitions. The ν_4 mode band is the weakest of the ro-vibrational Raman spectrum of CH_4 and could not be observed in our CRS experiments: thus, the time-domain CH_4 CRS model was limited to the ν_2 mode spectrum. Each of the ν_2 mode transitions gives rise to a damped harmonic contribution that is computationally evaluated on a temporal grid up to 1 ns, with 8 fs step size, corresponding to a resolution of 0.03 cm^{-1} for the Fourier-transformed CRS signal. As pointed out by Chen *et al.*,⁵² the main challenge in the development of an accurate time-domain model for the CH_4 CRS spectrum is the staggering number of Raman transitions included in the MeCaSDa calculated spectral database computed at the University of Burgundy^{72,73} (and accessible at the following url: <http://vamdc.icb.cnrs.fr/PHP/methane.php>) and the resulting computational cost. Implementing the calculation of $\chi^{(3)}$ as a running sum for every transition input avoids exceeding memory limits of the available computational resources; each spectral calculation still needs to be iterated for different input values of temperature and, if the effects of molecular collisions are not negligible,

species concentrations in the probe volume. Hence, the calculation of spectral libraries for quantitative spectroscopy can be a cumbersome task, especially when the possibility of CH_4 CRS thermometry in non-equilibrium environments is considered.²¹ The computational cost can be reduced by implementing a filter to exclude transitions whose intensity is lesser than the cut-off value of the strongest transition in the spectrum,⁵² depending on both the strength of the associated instantaneous dipole and on the Boltzmann population.

Figure 3(a) shows the transition polarizability for a small fraction of the Raman transitions considered in the database, corresponding to the fundamental band of the ν_2 mode spectrum ($0000 \leftarrow 0100$) and its first two hot bands ($0001 \leftarrow 0101$ and $0100 \leftarrow 0002$): the overall database spans more than 18 orders of magnitude in the Raman polarizability. These are combined with the differential Boltzmann population between the initial and final transition states at different temperatures, between 300 and 1500 K in Fig. 3(b), to compute the overall line intensities of the CRS spectrum. Implementing different cut-off values the number of Raman transitions taken into account in the calculation of $\chi^{(3)}$ can be significantly reduced, at cost of considering only a fraction of the Boltzmann population. In particular, adopting a tolerance of 1% or 0.1% of the strongest line in each branch most of the cumulative distribution function (CDF) is neglected. Increasing the tolerance to a millionth (billionth) of the strongest line more than 40% (90%) of the CDF is taken into account at all temperatures considered in the theoretical CRS library. The non-monotonic behavior shown by the corresponding curve in Fig. 3(b) depends on the combination of the different Raman polarizabilities and of the spreading Boltzmann distribution at higher temperatures. A cut-off value of

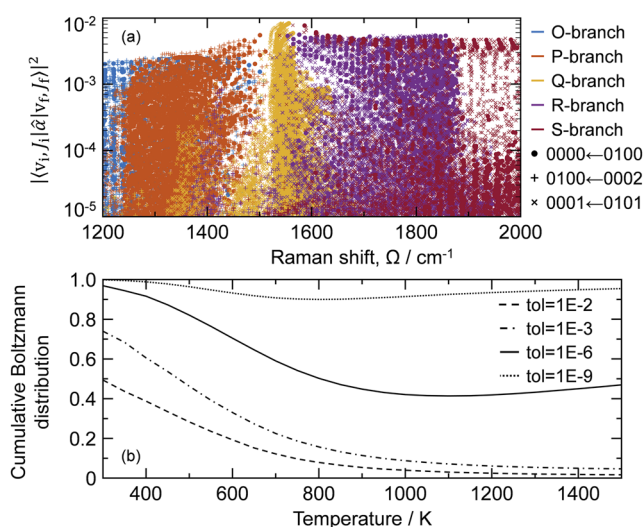


FIG. 3. The spontaneous Raman cross-section for CH_4 ν_2 mode transitions. (a) Transition moment of the polarizability anisotropy for the fundamental band ($0000 \leftarrow 0100$) and the first two hot bands ($0100 \leftarrow 0200$ and $0001 \leftarrow 0101$) of the ν_2 mode, according to the different Raman selection rules. (b) The cumulative distribution function, representing the fraction of the Boltzmann population included in the calculation of $\chi^{(3)}$ as a function of the input temperature and cut-off value selected for the filtering of the spectral database.

1×10^{-4} was found to be sufficient to model the time-domain behavior of the CH_4 ν_2 CRS spectrum at room temperature, where the effect of vibrational hot bands is measured to be negligible. In the flame experiment, on the other hand, the cut-off value needs to be increased to guarantee the independence of CH_4 ν_2 CRS thermometry from spectral filtering. An unfiltered spectral library was used to fit the flame spectra up to ~ 800 K, and a cut-off value of 1×10^{-6} was then found to perfectly reproduce these results while including only half of the cumulative Boltzmann distribution as shown in Fig. 3(b).

The time-domain CH_4 ν_2 CRS model is complemented by the inclusion of Raman linewidths, representing the collisional dephasing of the signal. The collisional dephasing coefficients of the Q-branch spectrum $\Gamma_{(v_i, J_i) \rightarrow (v_f, J_f)}$ are modeled according to the modified exponential gap (MEG) law^{77,78} as

$$\Gamma_{(v_i, J_i) \rightarrow (v_f, J_f)} = \sum_{k \neq j} \gamma_{kj}, \quad (10)$$

with upward and downward collisional transition rates, between the i th and j th rotational energy states, given by

$$\gamma_{ji} = \alpha p \left(\frac{T_0}{T} \right)^n \left(\frac{1 + a E_i / k_B T}{1 + a E_j / k_B T} \right)^2 \exp \left[\frac{-\beta (E_j - E_i)}{k_B T} \right], \quad (11)$$

$$\gamma_{ij} = \frac{2J_i + 1}{2J_j + 1} \gamma_{ji} \exp \left(\frac{E_j - E_i}{k_B T} \right), \quad (12)$$

where p is the pressure, $T_0 = 296$ K is the reference temperature, E_i and E_j represent the ro-vibrational energy in the upper and lower vibrational states (so that the energy gap is always positive), respectively, and α , β , δ , a , and n are fitting parameters for the scaling law. In the present linewidth model, the effect of the symmetry of ro-vibrational wave function is neglected⁴² so that the linewidth computed for a specific value of the rotational quantum number J is applied to all its symmetry components in Table I; similarly, the same RET rate is assumed for all the vibrational hot bands. Following the work of Chen *et al.*,⁵² the species-specific constant a is set to 2, and the number of fitting parameters to be simultaneously determined is reduced by independently fitting α and β at room temperature, assuming $\delta = 1$ and $n = 0$. While the MEG model has been successfully applied to isotropic ν_1 Q-branch CRS spectrum of CH_4 in multiple studies,^{44,45,52} the sum rule in Eq. (10) is not rigorously satisfied in the case of the anisotropic CH_4 ν_2 lines, where molecular reorientation and degenerate levels are present.⁷⁹ To a first approximation, this fact can be neglected, and the corresponding Raman lines can be assumed to satisfy Eq. (10): this approximation, though coarse, was employed with satisfactory results to model the collisional line broadening of the P- and R-branch lines in the absorption spectra of CO_2 .⁸⁰ We justify the present use of the MEG scaling law as a first-approximation model of the collisional dephasing of the anisotropic CH_4 ν_2 Q-branch spectrum in view of the fact that only atmospheric pressure experiments are here reported so that the molecular reorientation in inelastic collisions and the intra-branch coupling can be reasonably neglected. On the other hand, the sum rule in Eq. (10) is not directly used to compute the dephasing coefficients of the O-, P-, R-, and S-branch lines, as the inter-branch coupling is assumed to dominate over the intra-branch contribution.⁷⁹ The collisional dephasing rates of these lines are computed

assuming the “random phase approximation” (RPA), whereby the width of an anisotropic line involving a change in the rotational quantum number ($J_i \rightarrow J_f$) depends only on the relaxation of the rotational energy levels labeled by the quantum numbers J_i and J_f , as⁸¹

$$\Gamma_{(v_i, J_i) \rightarrow (v_f, J_f)} = \frac{1}{2} (\Gamma_{(v_i, J_i) \rightarrow (v_f, J_i)} + \Gamma_{(v_i, J_f) \rightarrow (v_f, J_f)}). \quad (13)$$

III. RESULTS AND DISCUSSION

A. Time-resolved CH_4 ν_2 spectrum

An example of a single-shot CH_4 ν_2 coherent Raman Stokes spectrum (CSRS), acquired in a room-temperature methane flow in the spectral range 1100 – 2000 cm^{-1} is shown in Fig. 4. The CH_4 ν_2 spectrum presents all five branches corresponding to the selection rules $\Delta J = 0, \pm 1$, and ± 2 . The spectral resolution of our CRS instrument (limited by the 4.1 cm^{-1} width of the probe spectrum) is insufficient to resolve the rotational structure of the Q-branch, which is then a single convolved feature at ~ 1535 cm^{-1} . The ordering of the branch labels with the Raman shift on the Stokes of the probe line is opposite to the one on the anti-Stokes side (CARS), as shown in Fig. 4: the O- and P-branch appearing at larger Raman shifts than the R- and S-branch.

The spectroscopic data for the CH_4 Raman spectrum discussed in Sec. II can be directly applied to the CSRS spectrum, with the only caveat that the branch labels for negative and positive changes in J need to be interchanged. A comparison of the experimental measurements and model prediction of the dynamic behavior of the CH_4 ν_2 CRS spectrum is provided in Fig. 5.

The coherence beating in time-resolved CRS spectra is a well-known phenomenon due to the presence of unresolved spectral lines at the resolution of the CRS instrument. Examples of this are found in hybrid fs/ps CRS spectra of air, where some N_2 and O_2 lines cannot be resolved for probe pulse durations lesser than ~ 60 ps and in high-temperature N_2 spectra due to the presence of unresolved hot bands.⁸² The presence of significant vibrational hot bands in room-temperature CRS spectra is uncommon for diatomic molecules, with relatively high vibrational constant (e.g., ~ 2330 cm^{-1} for N_2) but

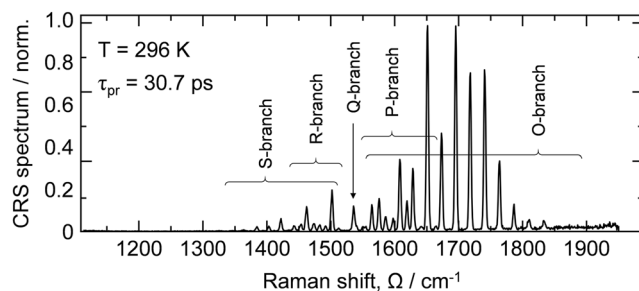


FIG. 4. Single-shot CH_4 ν_2 CSRS spectrum at room-temperature. The ν_2 mode Raman-activity is completely anisotropic and the selection rules allow for O-, P-, Q-, R-, and S-branch transitions. Note that, as the CRS signal is acquired in the Stokes side, the usual branch ordering is reverted with negative changes in the total angular momentum quantum number ($\Delta J = -1$ and -2 for the P- and O-branches, respectively) determining a larger frequency transition.

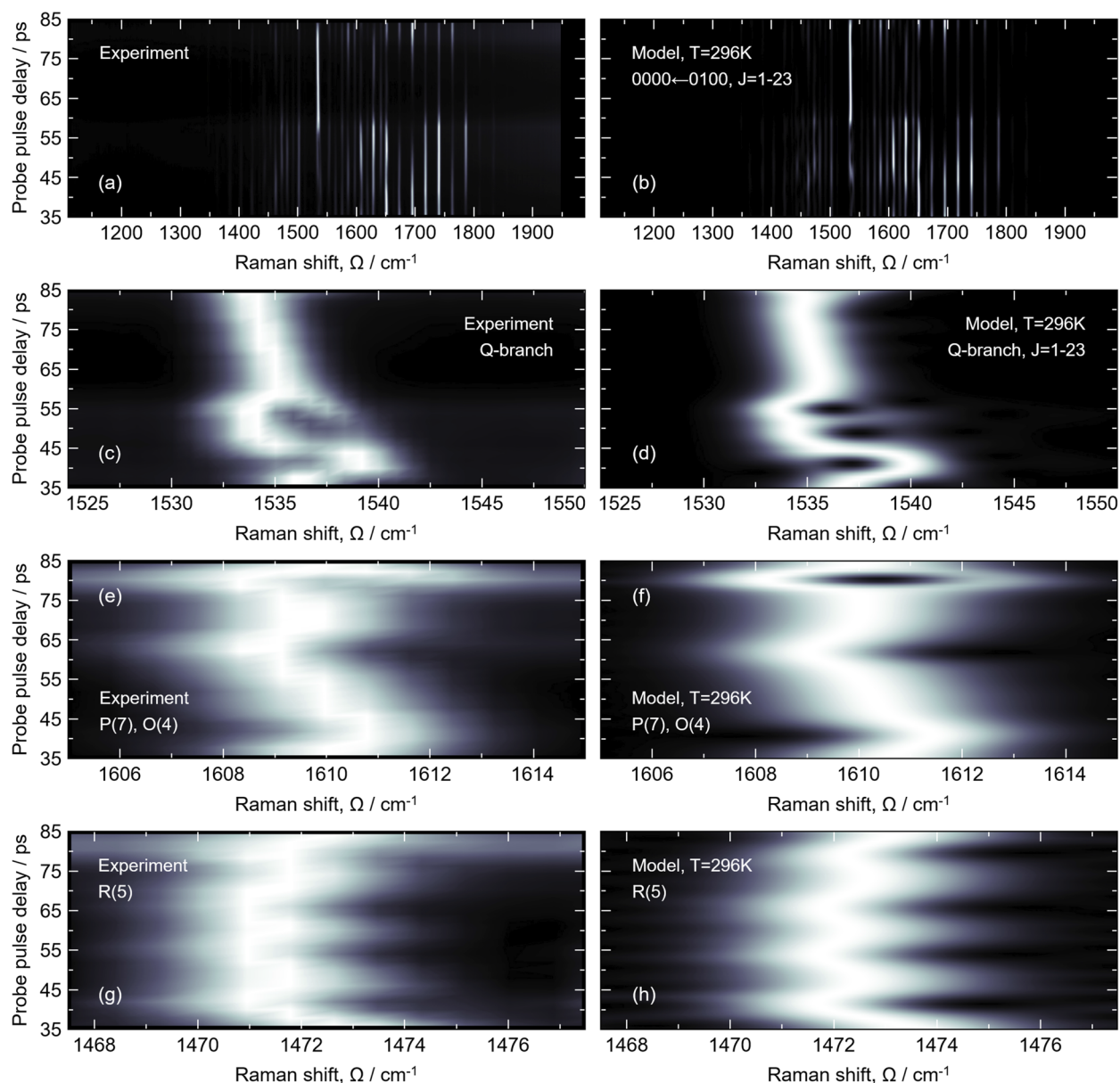


FIG. 5. Time-resolved CH_4 ν_2 CRS spectra: comparison between experiments and time-domain model. (a) Experimental spectrochronogram of the ro-vibrational CH_4 ν_2 mode, acquired in a room-temperature CH_4 flow over for values of the probe pulse delay spanning over 50 ps. A 1000-shot-averaged CRS spectrum is shown at each probe delay. (b) Modeled spectrochronogram including only the fundamental band of the ν_2 mode (i.e., $0000 \leftarrow 0100$) and rotational states up to $J = 23$. (c) Experimental spectrochronogram of the ν_2 Q-branch spectrum. (d) Modeled spectrochronogram of the ν_2 Q-branch spectrum. (e) Experimental spectrochronogram of the convolved P(7) and O(4) ro-vibrational lines. (f) Modeled spectrochronogram of the convolved P(7) and O(4) ro-vibrational lines. (g) Experimental spectrochronogram of the R(5) line of the P-branch spectrum. (h) Modeled spectrochronogram of the P(6) line of the P-branch spectrum.

more complex polyatomic molecules can have vibrational modes at lower frequencies (e.g., $\sim 667\text{ cm}^{-1}$ for the bending mode of CO_2) so that excited vibrational states can have a significant Boltzmann population.⁵⁴ In order to assess the impact of vibrational hot bands on the temporal beating of the room-temperature CH_4 ν_2 CRS spectrum, the time-domain model was here limited to the inclusion of

fundamental transitions between the ground and first vibrationally excited states: as the model correctly reproduces the beating pattern in the experimental spectrogram, vibrational hot bands must have a negligible impact on this behavior.

It is useful to distinguish the dynamics of the different branches in the CH_4 ν_2 CRS spectrum to understand the origin of its beating.

As shown in Fig. 5(c), the Q-branch lines appear as a single unresolved spectral feature centered approximately at 1535 cm^{-1} and shifting by $\sim 5\text{ cm}^{-1}$ on either side, depending on the probe pulse delay. Its beating is captured by the model in Fig. 5(d), including rotational lines up to state $J = 23$: it is reasonable to interpret this as the interferogram of these unresolved lines. It is then interesting to analyze the other spectral branches, which are fully resolved in Fig. 5. As an example, the dynamic behavior O(4) at $\sim 1610\text{ cm}^{-1}$ is represented in Fig. 5(e) and modeled in Fig. 5(f): a severe beating is still observed, which even leads to the line splitting at a probe delay of $\sim 80\text{ ps}$. While the beating can be partly attributed to the fact that the O- (and S-) branch lines are as a matter of fact not isolated, but overlap to the P- (R-) branch lines, the interference of only two harmonic contributions is expected to give rise to a simple sinusoidal pattern. A non-trivial beating pattern is also observed for line R(5) in Figs. 5(g) and 5(h), which does not overlap to any line of the S-branch and is perfectly isolated. Besides unresolved lines due to different chemical species, and to vibrational hot bands, coherence beating in time-resolved CRS spectra can arise from intramolecular interaction, whereby the coupling of two energy degrees of freedom breaks some implicit symmetry of the molecular wave function, thus splitting otherwise degenerate energy states. This is the case, e.g., in the pure-rotational CRS spectrum of O_2 , where the spin-orbit coupling gives rise to the beating of unresolved triplet transitions,⁸³ and for the ν_2 mode Raman spectrum of CH_4 . In this case, the gyroscopic coupling between the ro-vibrational levels of the ν_2 and ν_4 modes in the CH_4 dyad lifts the degeneracy of the ro-vibrational energy associated with the total angular momentum quantum number J . This effect can be thought of in classical terms as

the introduction of a non-inertial Coriolis force^{84,85} in the molecule-fixed frame of reference, which introduces an orientation-dependent contribution to the rotational energy of the molecule. The ν_2 - ν_4 Coriolis coupling is explicitly taken into account in the tetrahedral formalism employed to compute the spectral database⁸⁶ so that the rotational sub-states associated with the same J are treated as non-degenerate, and the time-domain CRS model correctly reproduces their coherence beating.

B. Linewidth measurements

The collisional dephasing of the CH_4 ν_2 CRS spectrum is presented in Fig. 6.

First, the dephasing of the spectrally integrated Q-branch signal in a room-temperature, atmospheric flow of CH_4 is shown in Fig. 6(a): the beating of the unresolved J -lines in the Q-branch spectrum has a significant impact on the CRS signal intensity. Thus, the temporal behavior of the CRS signal is determined by both the coherence beating and the collisional dephasing of the rotational lines. The collisional dephasing is modeled according to the MEG scaling law in Eqs. (10)–(12) by fitting simultaneously the α and β parameters. In order to highlight the effect of the molecular collisions, the theoretical behavior in their absence is also shown (dashed line) in Fig. 6(a). The experimental dephasing of the CH_4 ν_2 Q-branch spectrum and the prediction of the collisional model show a satisfactory agreement over the measured probe delay range of 20–180 ps and prove that, at least to a first approximation and under ambient conditions, the collisional dephasing of the anisotropic Q-branch lines can be well approximated by the sum rule with a

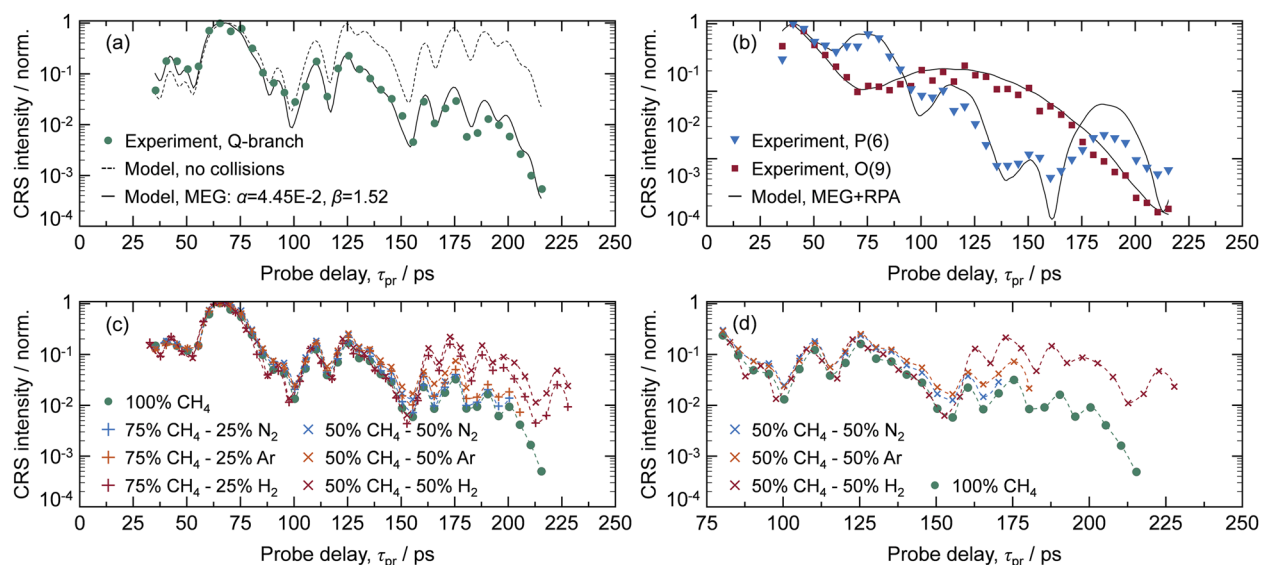


FIG. 6. Collisional dephasing of the CH_4 ν_2 CRS signal. (a) Comparison of the experimental dephasing of the spectrally integrated Q-branch signal and the MEG model in a room-temperature atmospheric CH_4 flow. (b) Experimental dephasing of isolated lines of the CH_4 ν_2 spectrum in a room-temperature atmospheric CH_4 flow: O(9) and P(6) are selected as an example. The experimental behavior is compared to the linewidths computed adopting the RPA. (c) Experimental dephasing of the spectrally integrated isotropic Q-branch signal in binary mixtures with N_2 , H_2 , and argon, compared to the dephasing in a pure CH_4 flow. (d) Detail of the experimental dephasing of the CH_4 signal self-perturbed and in 50%–50% binary mixtures with N_2 , H_2 , and argon for probe delays 100–225 ps.

MEG scaling law for the state-to-state relaxation rates. The effect of inter-branch coupling, on the other hand, is taken into account in the collisional dephasing model for the O-, P-, R-, and S-branch lines, by introducing the RPA. The corresponding dephasing rates are thus computed according to Eq. (13) from the dephasing of the Q-branch lines, previously calculated by the MEG scaling law. Figure 6(b) shows the experimental dephasing of two isolated lines in the O- and P- branches [namely O(9) and P(6)] and presents a comparison to the prediction of the RPA applied to the MEG-modeled Q-branch dephasing coefficients. The model predicts the collisional dephasing of these lines to a reasonable agreement with the experimental data, particularly for line O(6) although the predicted dephasing introduces an undue damping of the coherence beating of P(9) for probe delays larger than ~ 150 ps.

The effect of different collisional partners in binary mixtures with CH_4 was also investigated, as shown in Fig. 6(c). The experimental decay of the CH_4 ν_2 Q-branch signal, as measured in pure CH_4 or in the binary mixtures with N_2 , H_2 , and argon, shows negligible dependence on the actual mixture composition for a probe delay lesser than ~ 80 ps. This observation confirms for the ν_2 mode spectrum the same behavior reported for the CH_4 ν_1 Q-branch in binary mixtures with N_2 by Engel *et al.*⁴⁸ Chen *et al.* defined a critical probe delay for the collisional partner independence of the CH_4 ν_1 CRS signal and estimated this to be ~ 100 ps;⁵² they experimentally confirmed it by measuring the collisional dephasing of the CH_4 CRS signal in binary mixtures with argon, as well as N_2 . They employed a simple model based on gas kinetics to estimate the mean time between molecular collisions and thus define a collision sensitivity time. Below such a timescale, coherence beating dominates the temporal evolution of the CH_4 ν_2 CRS signal in Fig. 6(c). For probe delays larger than ~ 80 ps the dephasing of the signal shows sensitivity to the collisional environment, as shown more in detail in Fig. 6(d). N_2 and argon present similar behaviors as collisional partners to the CH_4 molecules: this confirms the findings by Chen *et al.* who estimated the collisional linewidths of the CH_4 ν_1 Q-branch at room temperature and 500 Torr to be 0.167 and 0.158 cm^{-1} for binary mixtures with 90% N_2 and argon, respectively. On the other hand, H_2 is a much weaker perturber to the CH_4 ν_2 CRS signal as demonstrated by the comparison with the self-perturbed behavior in Fig. 6(d): at 210 ps, the CH_4 CRS signal in a binary mixture with 50% H_2 is almost two orders of magnitude larger than the corresponding signal in a flow of pure CH_4 . This observation is explained by the largely different rotational energy manifolds of the CH_4 and H_2 molecules, such that any rotational state of the CH_4 scatterers has a low density of neighboring states of the H_2 perturbers. Thus, the collisional RET is significantly lessened in case of binary mixtures of CH_4 and H_2 . In the case of N_2 , on the other hand, the radiator and perturber have plenty of neighboring rotational energy states, promoting the RET in inelastic collisions.

C. CH_4 /air diffusion flame spectroscopy

The potential of CH_4 ν_2 CRS spectroscopy for *in situ* diagnostics in chemically reactive flows is demonstrated by performing spatially resolved measurements across a laminar CH_4 /air diffusion flame.

CRS measurements over the spectral range ~ 1100 – 2000 cm^{-1} were performed with a probe delay of 30.7 ps at each of the 25

locations, spaced by 0.5 mm across the flame front, moving from the fuel stream at the center of the burner ($y = 0$ mm) to the oxidizer stream ($y = 12$ mm). The probe delay was chosen to time-gate the NR background and to maximize the signal-to-noise ratio (SNR), balancing the detection of the different branches of the CH_4 spectrum as visible in Figs. 6(a) and 6(b). The single-shot detection of the CH_4 ν_2 CRS signal was achieved up to location 6.5 mm, at temperatures as high as ~ 800 K: Fig. 7 presents different examples of single-shot CH_4 ν_2 CRS spectra acquired at three different locations ($y = 2.5$, 5.5, and 6.5 mm) in the flame. Moving from the center of the burner toward the chemical reaction zone of the flame, the CH_4 ν_2 CRS signal is negatively impacted by the combination of increasing temperature and reducing CH_4 concentration so that, at $y = 6.5$ mm, the SNR is reduced to less than 10. In order to extend the detection limit of the CH_4 ν_2 CRS signal and validate the time-domain model at higher temperatures, ten-shot-averaged spectra were acquired at each flame location: the CH_4 could be detected up to 1020 K at $y = 8$ mm. Figure 8 shows the ultrabroadband CRS spectra acquired across the CH_4 /air flame front. Four different chemical species are detected in the vibrational fingerprint region of the Raman spectrum from 1100 to 2000 cm^{-1} : namely, the ro-vibrational ν_2 spectrum of CH_4 (from 1300 to 1950 cm^{-1}), the CO_2 Fermi dyad (with fundamental Q-branch at 1285 and 1388 cm^{-1}), the pure-rotational spectrum of H_2 (four O-branch lines at 1246, 1447, 1637, and 1815 cm^{-1}), and the ro-vibrational spectrum of O_2 (with fundamental Q-branch at 1556 cm^{-1}). Thus, spatially resolved ultrabroadband CRS provides a window to monitor the physical-chemical processes *in situ* by measuring the local temperature and detecting some of the major reactants and products. The spectrum in Fig. 8(a) is representative of the fuel stream from the center of the burner to $y \approx 5$ mm (as shown also in the comparison of the single-shot spectra at $y = 2.5$ and 5.5 mm in Fig. 7): the spectrum is dominated by the ro-vibrational ν_2 mode lines of CH_4 and by the characteristics Fermi dyad of the CO_2 spectrum, particularly the “blue” fundamental band of CO_2 at 1388 cm^{-1} is the single highest spectral feature. Such an abundance of CO_2 at the center of the fuel stream can be explained by the buoyancy of the lightweight CH_4 molecule (relative molecular mass: $m = 16$) against the heavier CO_2 ($m = 44$) produced in the

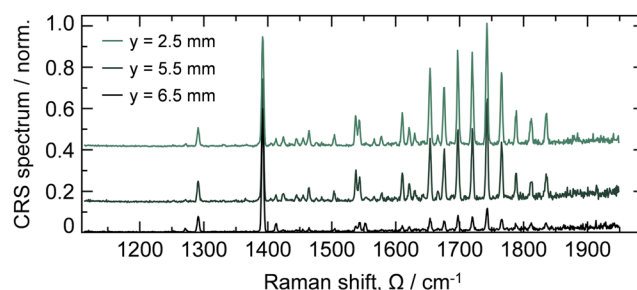


FIG. 7. Single-shot CRS spectra acquired at locations $y = 2.5$, 5.5, and 6.5 mm across the CH_4 /air diffusion flame front. The probe delay of 30.7 ps is for all spectra. The peaks at 1265 and 1285 cm^{-1} belong to the “red” dyad of the CO_2 spectrum, while the peaks at 1388 and 1409 cm^{-1} belong to the “blue” dyad. O_2 can also be detected by its ro-vibrational Q-branch spectrum at 1553 cm^{-1} , which presents a clear increment moving toward the oxidizer stream. The rest of the spectral lines are attributed to the ν_2 mode spectrum of CH_4 .

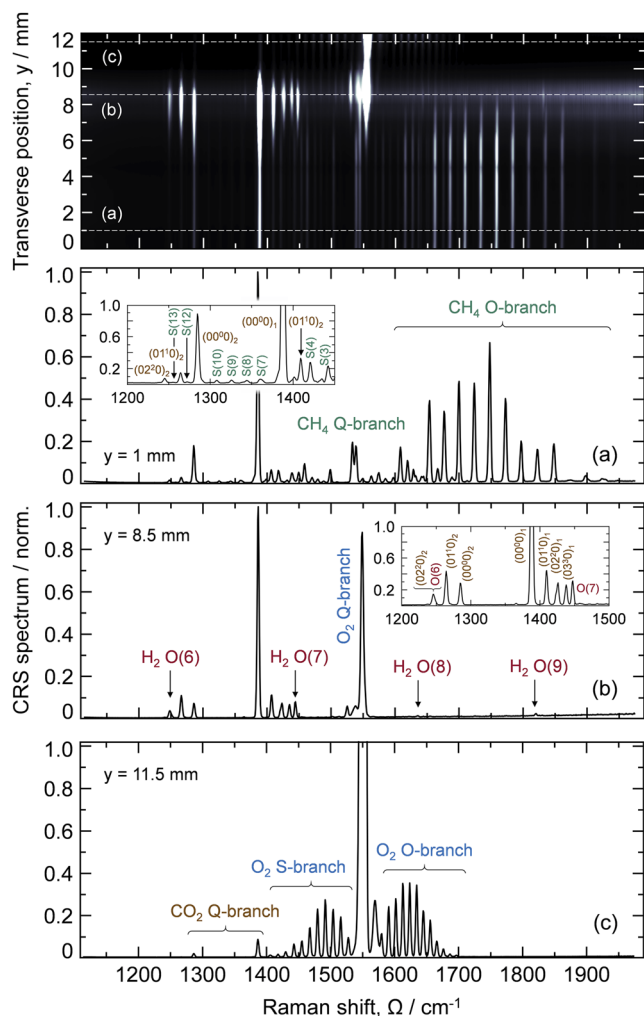


FIG. 8. Ultrabroadband CRS spectra in the molecular fingerprint region, measured across the laminar CH_4/air diffusion flame. The probe delay of 30.7 ps is for all spectra. Note that gamma compression is employed in the image post-processing. (a) Fuel steam ($y = 1$ mm): low-temperature ro-vibrational spectra of CO_2 and CH_4 . (b) Reaction layer ($y = 8.5$ mm): moving toward the reaction zone, CH_4 undergoes pyrolysis producing H_2 , which is then rapidly consumed in the chemical reaction. Four lines of the pure-rotational Raman spectrum of H_2 are detected in the window $\sim 1100\text{--}2000\text{ cm}^{-1}$. The increased temperature is evident in the multiple hot bands of the ro-vibrational CO_2 and O_2 spectra. (c) Oxidizer stream ($y = 11.5$ mm): the oxidizer is ambient air so that the Raman spectrum in the fingerprint region is dominated by the ro-vibrational spectrum of O_2 . A small amount of the CO_2 produced in the combustion reaction diffuses into the oxidizer stream.

whole reaction zone volume of the flame. This results in the internal recirculation of the high-temperature CO_2 and its local mixing with the room-temperature CH_4 flow at the burner inlet: upon thermalization, the local temperature is higher than 296 K as attested by the clear detection of the first hot band in the CO_2 spectrum⁵⁴ at ~ 1265 and $\sim 1409\text{ cm}^{-1}$, and even a second hot band in the red dyad at 1244 cm^{-1} as shown in the inset in Fig. 8(a). It is worth noting here that this local CO_2 recirculation is expected to impact the

combustion chemistry due to its large heat capacity: this point will be further discussed when presenting the results of CRS thermometry. A small amount of fuel mixing with air is also visible in the spectrum: the small peak at 1556 cm^{-1} is indeed the ro-vibrational Q-branch of O_2 . This becomes much more pronounced in the high-temperature spectrum in Fig. 8(b), acquired in the reaction zone at $y = 8.5$ mm.

In this region of the flame, the combustion reaction sustains itself by releasing heat that, in turns, provides the activation energy for the dissociation of the CH_4 molecules into radicals leading to the generation of H_2 , which is then rapidly oxidized. Indeed, the CH_4 ν_2 CRS spectrum is barely detectable at this location, while four lines of the pure-rotational H_2 O-branch spectrum are identified: namely, O(6) at 1246 cm^{-1} (overlapping with the second hot band of the CO_2 spectrum, as highlighted in the inset), O(7) at 1447 cm^{-1} , O(8) at 1637 cm^{-1} , and O(9) at 1815 cm^{-1} . The heat released in the oxidation of the fuel increases the local temperature as marked by the pronounced hot bands in the Q-branch spectra of CO_2 ; these are marked in the figure by the lowest vibrational state, denoted according to the notation in Herzberg⁸⁷ as $\nu_1\nu_2^l\nu_3$. As the oxidizer is ambient air, the temperature rapidly drops moving away from the reaction zone and the spectrum at $y = 11.5$ mm is dominated by the low-temperature O_2 CRS signal, with the ro-vibrational O- and S-branch spectra also visible in Fig. 8(c).

Figure 9 illustrates how the CH_4 ν_2 CRS spectrum is affected by temperature by comparing a single-shot spectrum acquired in at room temperature, and two 10-shot-averaged spectra acquired at two different locations across the flame and fitted to the time-domain model measuring the local temperature to 550 and 700 K, respectively. As the temperature increases, a higher rotational state becomes populated, as reflected by the spread of the spectral envelope; in addition, the coherence beating of the O- and P-branch lines, and possibly of the vibrational hot bands, determines a shift of the strongest line from O(6) at room temperature to O(10) in the flame, while the probe delay is kept at 30.7 ps for all the spectra. An analogous behavior is shown by the S-branch spectrum, with the dominating line at room temperature being S(0) and shifting to S(2) at 550 K. The small peak at 1447 cm^{-1} at 700 K is attributed to line O(7) in the pure-rotational H_2 CRS spectrum. A significant change is also observed in the unresolved Q-branch spectrum as the temperature increases, although the effect is somewhat obscured

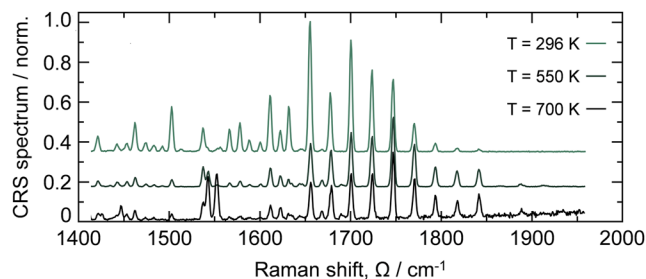


FIG. 9. Temperature dependence of the CH_4 ν_2 CRS spectrum. Single-shot spectrum acquired in an open room-temperature CH_4 flow (light green), and 10-shot-averaged spectra acquired at 550 K (dark green) and 700 K (black) at location $y = 3.5$ and 6 mm, respectively.

at $T = 700$ K because of the overlap with the ro-vibrational O_2 Q-branch spectrum whose intensity increases moving toward the oxidizer stream.

Ultrabroadband CRS in the fingerprint region has so far provided us with qualitative insights into the mixing and chemical processes taking place across the diffusion flame. In order to perform quantitative measurements, it is necessary to validate the time-domain model for the CH_4 ν_2 CRS spectrum at high temperatures. This is done by performing direct CH_4 thermometry by matching the experimental ten-shot-averaged spectra up to $y = 6.5$ mm against the synthetic library and comparing the results to CO_2 CRS thermometry, developed and validated in our previous study.⁵⁴ An example of the fitting of an experimental CH_4 ν_2 CRS spectrum (at $y = 0$ mm) to the time-domain CRS model is shown in Fig. 10(a). All the five ro-vibrational branches of the CH_4 ν_2 spectrum in the region ~ 1400 – 1950 cm^{-1} are fitted simultaneously to the synthetic spectra in the library employing a damped least-squares algorithm. As the temperature increases toward the reaction zone, the hot bands of the CO_2 spectrum cover the S- and R-branch lines of the CH_4 spectrum; similarly, the increasing oxygen concentration moving to the oxidizer stream results in the mixing of the Q-branch lines of the CH_4 and O_2 spectra. Hence, from location $y = 4$ mm up to 7 mm, the fit was limited to the P- and O-branch lines of the CH_4 ν_2 spectrum. All the five ro-vibrational branches of the CH_4 ν_2 spectrum in the region ~ 1400 – 1950 cm^{-1} are fitted simultaneously to the synthetic spectra in the library employing a damped least-squares algorithm. As the temperature increases toward the reaction zone, the hot bands of the CO_2 spectrum cover the S- and R-branch lines of the CH_4 spec-

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Figure 10(b) presents the validation of CH_4 ν_2 CRS thermometry (green markers) by means of comparison with ro-vibrational CO_2 CRS (black markers): the abundance of CO_2 in the diffusion flame tested and its detection at all the measurement locations allow us to use CO_2 to measure the temperature profile across the whole flame front. As briefly mentioned in the previous paragraph, the large concentration of CO_2 in the center of the burner, as detected in the ultrabroadband CRS spectra in Fig. 8, has a significant impact on the combustion chemistry, given the large heat capacity of the CO_2 molecules that thus act as heat sinks in the reaction zone on the flame. The result is a smoother temperature profile than expected for a laminar CH_4 /air diffusion flame:⁸⁸ the temperature at the center of the fuel stream ($y = 0$ mm) is larger than the room temperature of the inlet CH_4 flow, while the maximum temperature in the reaction zone (~ 1430 K) is significantly lower than expected for a laminar axisymmetric methane/air diffusion flame.⁸⁸ The thermal effect of the back-diffusing CO_2 on the measured flame temperature in our experiment is comparable to the reduction in the adiabatic flame temperature measured in laminar premixed CH_4 /air flames for a 20% CO_2 dilution.⁸⁹ Comparing the temperature measurements obtained by CH_4 ν_2 CRS and CO_2 CRS, we can validate our time-domain CH_4 CRS model at temperatures as high as ~ 800 K and quantify the accuracy of CH_4 ν_2 CRS thermometry. The two ther-

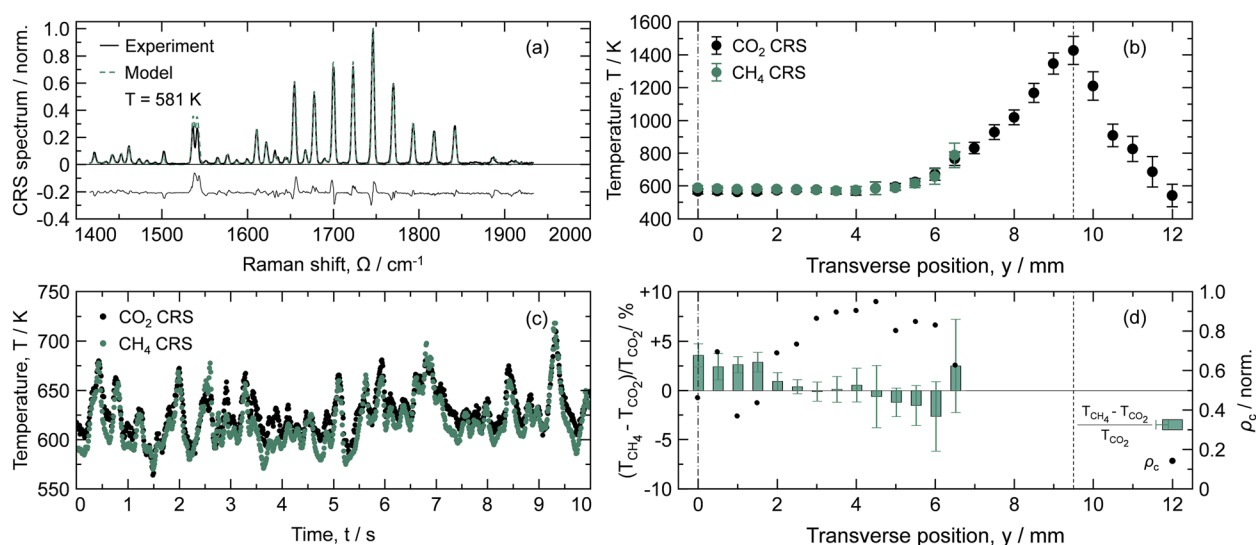


FIG. 10. Ro-vibrational CH_4 ν_2 CRS thermometry. (a) Experimental ten-shot-averaged CH_4 ν_2 CRS spectrum acquired at location $y = 0$ mm and comparison to the time-domain CRS model for thermometry. (b) Average temperature profile across the laminar CH_4 /air diffusion flame as measured by ultrabroadband fs/ps CRS thermometry: ro-vibrational CO_2 thermometry (black) provides validation to ro-vibrational CH_4 thermometry (green). (c) Comparison of CO_2 and CH_4 CRS thermometry over a sample of 1000 ultrabroadband spectra, acquired at $y = 5.5$ mm: the same dynamics is reproduced by the two independent methods, which proves the physical nature of the temperature oscillations. (d) Accuracy and precision of CH_4 ν_2 CRS thermometry as compared to CO_2 CRS thermometry: the thermometric accuracy (green bars) is better than 3% at all measurement locations, while the precision is fundamentally limited by the physical fluctuations in the flame. Concordance correlation (black dots) quantifies the agreement of the temperature dynamics measured by CO_2 and CH_4 CRS at each location.

metric techniques show a very satisfactory agreement not only in terms of the average temperature measured at each flame location but also within each temporal sequence of 1000 frames, as shown in Fig. 10(c) for $y = 0$ mm.

The agreement between CH_4 ν_2 and CO_2 CRS thermometry is quantified in Fig. 10(d) evaluating the systematic bias between the two methodologies (bar plot), and by means of Lin's concordance correlation coefficient, defined as⁹⁰

$$\rho_c = \frac{2\sigma_{\text{CO}_2, \text{CH}_4}}{\sigma_{\text{CO}_2}^2 + \sigma_{\text{CH}_4}^2 + (\mu_{\text{CO}_2} - \mu_{\text{CH}_4})^2}, \quad (14)$$

where μ_{CO_2} and μ_{CH_4} are the mean temperature measured by CO_2 and CH_4 CRS, respectively, with corresponding standard deviations σ_{CO_2} , σ_{CH_4} , and covariance $\sigma_{\text{CO}_2, \text{CH}_4}$.

The formula in Eq. (14) quantifies the correlation between the two temperature measurements independently of the possible systematic bias between them and of the temperature fluctuations in the sample. The temperature dynamics measured by CO_2 CRS thermometry can thus be adequately reproduced by CH_4 ν_2 CRS, with an accuracy better than 3% at all measurements location. The satisfactory frame-by-frame comparison between CO_2 and CH_4 ν_2 CRS thermometry, as quantified by the concordance correlation factor, also allows us to establish the origin of the fluctuations in the measured temperature due to oscillations in the flame.

Figure 11(a) shows the spectra of the temperature dynamics as measured by CO_2 and CH_4 ν_2 CRS thermometry at $y = 3$ mm and

represented by circles in Fig. 11(b). These spectra present common frequencies in the range ~ 0.1 –12.5 Hz (the upper limit being the value where the spectra of T_{CH_4} and T_{CO_2} diverge), corresponding to the frequencies associated with the flame oscillations. Applying a spectral filter in this window,⁹¹ it is possible to isolate the inherent noise in the measurements, represented by the solid lines in Fig. 11(b), even though a lower-frequency oscillation (<0.1 Hz), unresolved over the 10 s acquisition window, can still be made out in the smoothened profiles. Thus, the thermometric precision was improved from 1.9% for CH_4 ν_2 CRS and 2.4% for CO_2 CRS to 0.57% and 0.91%, respectively. The identification of the frequency at which the two spectra diverge depends on the metric used to define this divergence, but this has only a minor effect on the quantification of the measurement precision: filtering oscillations only up to 10 Hz resulted in a precision of 0.61% for CH_4 ν_2 CRS and 1.0% for CO_2 CRS, while extending the filter to 15 Hz changed it to 0.54% and 0.90%, respectively. The 12.5 Hz value was found to be the lowest applicable to all the datasets acquired in the flame: the resulting thermometric precision was better than 2% at all measurement locations for both CO_2 and CH_4 ν_2 CRS thermometry.

IV. CONCLUSIONS

We have reported the first investigation of the ν_2 mode Raman spectrum of CH_4 by means of coherent Raman spectroscopy: ultra-broadband two-beam fs/ps CRS was employed to perform time-resolved measurements of the ro-vibrational CH_4 ν_2 spectrum and to demonstrate its application as a combustion diagnostic tool. Our CRS instrument employs a single regenerative fs amplifier to generate both broadband fs and narrowband ps pulses; fs laser-induced filamentation is employed *in situ* to compress the fs pulse to <20 fs so as to excite the ro-vibrational Raman modes in the vibrational fingerprint region. This spectral region is of particular interest in the investigation of chemical reactions in gas-phase environments, as a number of Raman-active species have a spectral signature in the range ~ 1100 –2000 cm^{-1} .

We developed a time-domain CRS model for the CH_4 ν_2 spectrum employing the MeCaSDa calculated spectroscopic database with the position and cross-section of $\sim 10 \times 10^6$ Raman lines, as computed at the University of Burgundy.^{72,73} We employed this line list by rescaling the spontaneous Raman cross section by the differential Boltzmann population between the two ro-vibrational states involved in the CRS process, and we applied a spectral filter, considering only CRS lines stronger than a cut-off value of the strongest lines, to reduce the computational time of the temperature-dependent synthetic CRS libraries. We then validated the CRS model by performing CH_4 ν_2 thermometry in a laminar CH_4 /air diffusion flame: the temperature estimations by CH_4 ν_2 CRS were found to converge for cut-off values smaller than 1×10^{-6} for all the measured temperatures up to ~ 800 K.

The CRS model includes collisional Raman linewidths computed by a modified energy-gap scaling law for the Q-branch lines and adopting the random phase approximation for the derivation of the O-, P-, R-, and S-branch collisional dephasing rates. The collisional dephasing of the self-perturbed CRS signal in a room-temperature CH_4 flow is used to fit the α and β parameters of the MEG model. The use of the sum rule for the calculation of the

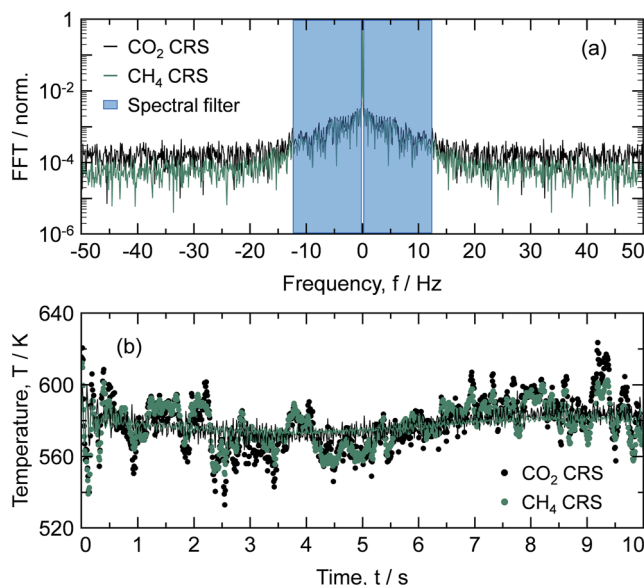


FIG. 11. Temperature dynamics measured by ro-vibrational CRS thermometry performed on CO_2 (black) and CH_4 (green). (a) Spectrum of the temperature dynamics: common frequencies, due to physical oscillations in the flame environment, are identified in the range of ~ 0.1 –12.5 Hz. (b) Original (circle) and filtered (solid line) temperature dynamics: removing the temperature fluctuations with common frequencies we estimate the inherent precision of CH_4 ν_2 CRS thermometry to be better than 1% for this sample.

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