

An inquiry into the limits of applicability  
of the refraction - diffraction equation .

Thesis in partial  
fulfillment of the  
requirements of the  
ingenieurs examination.

D.O. van Nieuwenhuijze

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Contents

page

1	<u>Introduction</u>	- 1 -
2	<u>Analytical Description</u>	
2.00	Introduction	- 4 -
2.01	Concepts of vector analysis	- 6 -
2.02	Potential flow	- 7 -
2.03	Boundary conditions of the field	- 9 -
2.04	Side View Model	- 9 -
2.05	Flat bottom solution	- 10 -
2.06	Top View Model	- 11 -
2.07	Comparison	- 12 -
2.08	A model for comparative calculation	- 15 -
2.09	The influx and efflux boundaries	- 17 -
2.10	Summary of the analytical description	- 20 -
3	<u>Numerical Description</u>	
3.00	Introduction	- 21 -
3.01	Finite Element Method	- 22 -
3.02	The Galerkin Method	- 24 -
3.03	Lagrange Polynomials	- 25 -
3.04	Algorithm for the translation	- 28 -
3.05	Describing the equations in numerical form	- 29 -
3.06	Computing element matrices and vectors	- 36 -
3.07	Summary of the numerical description	- 44 -
4	<u>Description of the computation</u>	
4.00	Introduction	- 45 -
4.01	Basic structure of the program	- 46 -
4.01.1	Preparing the computer for the job	- 47 -
4.01.2	Generate elements for the Finite Elm. Meth.	- 48 -
4.01.3	Prescription of the equations	- 49 -
4.01.4	Providing the input for the equations	- 49 -
4.01.5	Computing matrices and vectors per element	- 50 -
4.01.6	Assembling matrix " $\tilde{A}$ " and vector " $\tilde{g}$ "	- 50 -
4.01.7	Solving $\varphi$ from $\tilde{A}\varphi = \tilde{g}$	- 50 -
4.01.8	Presenting the solution	- 51 -
4.01.9	Signaling the end of the program	- 51 -



4.02	Using the computer program.	- 46 -
5	<u>Testing the program</u>	- 48 -
5.00	Introduction	- 48 -
5.01	Configuration for verification	- 48 -
5.02	Verifying the predictions	- 49 -
5.02.1	Wavelength	- 49 -
5.02.2	Magnitude of the potential:	- 52 -
5.03	Alternative verification	- 55 -
6	<u>Experimentation</u>	- 56 -
6.00	Introduction	- 56 -
6.01	Goal of the test	- 57 -
6.02	Parameters of the wavefield	- 57 -
6.03	Selected bottom configuration	- 58 -
6.04	Discussion of the parameters	- 61 -
6.05	Characteristic parameter of the wavefield	- 63 -
6.06	Three main types of variation	- 66 -
6.07	Selected parameter combinations	- 67 -
7	<u>Results</u>	- 71 -
7.00	Introduction	- 71 -
7.01	The generated mesh	- 72 -
7.02	Form of presentation of the results	- 74 -
7.03	Gathering data from the graphs	- 77 -
7.04	Results in tables	- 77 -
8	<u>Interpretation of the results</u>	- 80 -
8.00	Introduction	- 80 -
8.01	Presentation per model and per CASE	- 80 -
8.02	Evaluation of the predictions of $\ \varphi\ $	- 84 -
8.03	Evaluation of the values of $\frac{\ \varphi_t\ _r}{\ \varphi_t\ _s} - 1$	- 85 -
9	<u>Discussion of the results</u>	- 87 -
9.00	Introduction	- 87 -
9.01	Results of CASE 3	- 88 -
9.02	Influence of the parameters ScD, Ratio, ScB	- 90 -



10	<u>Conclusions</u>	- 97 -
11	<u>Recommendations for further studies</u>	- 99 -
12	<u>Acknowledgements</u>	- 101 -



## SAMENVATTING

In deze studie zijn twee modellen met elkaar vergeleken welke elk de waarden voor een golfpotentiaal voorspellen ( numeriek ).

Het eerste model, " Side View Model ", is gebaseerd op de vergelijking van Laplace en beschrijft in een verticaal vlak. Het tweede model, " Top View Model ", geeft een beschrijving in een horizontaal vlak ( het oppervlak ) en is gebaseerd op de Refractie-Diffractie vergelijking, zoals gepresenteerd door Berkhoff (1976).

Doel van het onderzoek was: het komen tot een inzicht ten aanzien van de gevoeligheid voor de verandering van de bodenhelling, in het Top View Model.

Voor de studie is een computerprogramma gemaakt dat, voor elk der twee modellen, de waarde van de golfpotentiaal aan het oppervlak berekent, voor een gekozen bodemconfiguratie. Het computerprogramma is gebaseerd op de eindige elementen methode en gebruikt het in Delft ontwikkelde pakket "AFEP" .

Als bodemconfiguratie is gekozen: een trapeziumvormige dam op een vlakke bodem; de hoogte van de dam, de helling en de lengte van het talud van de dam zijn gevarieerd en de voorspelde waarden zijn voor beide modellen vergeleken voor de identieke gevallen.

Het is gebleken dat er geen duidelijke directe correlatie is tussen de vergeleken waarden ( van de beide modellen ) en een verandering van de bodenhelling bij de gebruikte toetsingsparameter ( welke afgeleid is van de door Berkhoff gepresenteerde toetsingsparameter ). Uitgaande van de hier gepresenteerde resultaten zou verdergaand onderzoek tot meer inzicht kunnen leiden.



Wave phenomena have been studied with great interest throughout the ages, due to the importance of knowledge of this matter for various human activities ( such as commerce: shipping ; habitation: shore protection ).

A number of models have been designed, " equations ", which predict the changes in the characteristics of water waves as a result of the variations of the water-depth over which the wave travels.

One of these models, the " refraction - diffraction equation ", has been published by Berkhoff (1976).

The equation describes a wavefield with a free surface condition over an impermeable rigid bottom, which is flat or 'nearly' flat.

The requirement of the bottom to be at most 'nearly' flat follows from the way the equation has been derived:

Using a known solution for the case of a flat and horizontal bottom ( Lamb, 1975 ), Berkhoff has assumed that this solution might also be used as an approximative description for a 'nearly' flat bottom, with 'rather gentle' ( in quantitative sense ) slopes.

This study has been undertaken in order to find out what is to be understood by " rather gentle " slopes.

That is, it is an inquiry into the limits of applicability of the refraction - diffraction equation.



In this study, the refraction - diffraction equation, as well as the set of equations from which it has been derived, are used to predict the values of the wave potential at the water surface. The computation is done numerically, using A Finite Element Package ( " AFEP " ) which has been developed at Delft University ( Segal, 1975 ).

The predictions by the two models are compared for identical cases and a correlation between the slope of the bottom and the discrepancy in the predictions is sought.

Only simple linear waves are considered in this study.

This thesis is made up as follows:

Chapter 2 contains the presentation of the analytical core of the study: the set of equations which describe the wave field over a rigid impermeable bottom with a free surface condition, the refraction - diffraction equation which is derived from it and the relation between the two.

Two models are created based on these descriptions, respectively: " Side view model " and " Top view model " of which the predictions are to be compared for identical cases.

Chapter 3 presents the numerical form of these two models as they are to be used in AFEP.

Chapter 4 gives a description of the computer program.



The correctness of the predictions of the computer model are verified, for the case of a flat and horizontal bottom, in chapter 5.

The experimentation, described in chapter 6, involves the selection of the configuration of the bottom and a discussion of the characteristic parameters of the problem.

This leads to the selected combinations of parameters for the experimentation.

Chapter 7 presents the output of the computer program: the generated mesh; the form of presentation of the predicted values of the wave potential and the results.

The results are interpreted in chapter 8 and discussed in chapter 9, where the influence of all the parameters on the difference in the predictions of the two models is investigated.

These investigations show that no direct correspondence can be found between the variation of the characteristic parameters and the difference in the predictions of the two models. These conclusions are described in chapter 10.

Chapter 11 contains a number of recommendations for further studies which follow from this study.

Chapter 12 contains acknowledgements.



2.00            Introduction

This chapter starts with a brief summary of potential formulation, using vector notation.

The concepts are then used to describe a potential flow and boundary conditions of the fluid field: a free surface condition and impermeable rigid bottom.

The equations are presented which describe the same field in linearised form, in a vertical plane, for a simple harmonic wave. ( " Side view model " ).

The analytical solution of these equations is shown as it has been derived for the case of a flat bottom.

It is shown, briefly, how this solution is used by Berkhoff (1976) to reduce the initial set of equations to the refraction - diffraction equation, describing the wave characteristic ( the potential ) at the water surface. ( " Top view model " ).

The Side View Model and the Top View Model are compared and a simple setup is presented which may be used to indicate the difference in predictions by the two models.



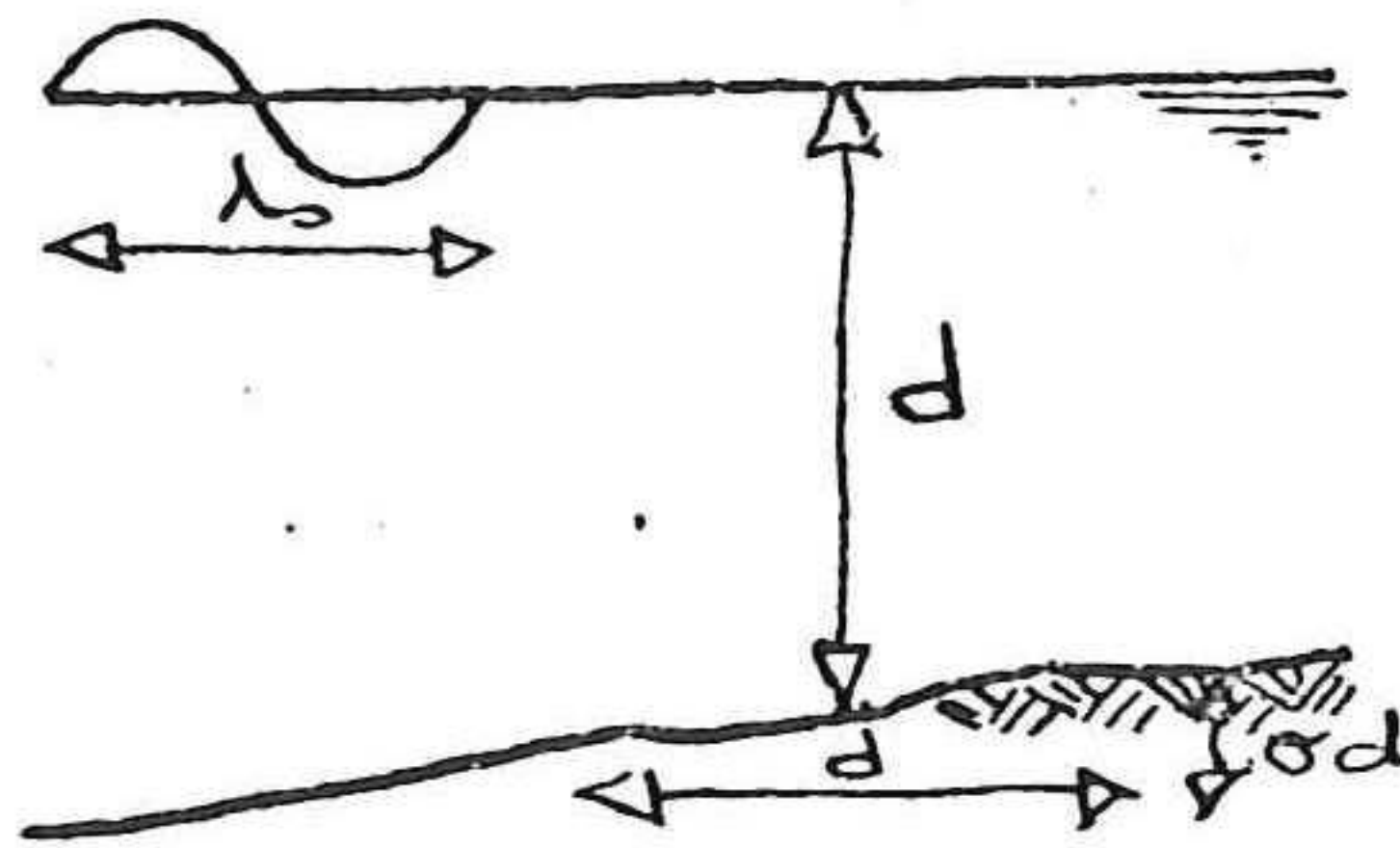
This study deals only with simple harmonic linear water waves. The fluid field is assumed to be free of currents; to be characterised by the Laplace equation.

Ignored are the effects of viscosity, compressibility, surface tension.

The wave potential  $\varphi$  is described in the complex plane:

$$\varphi = \varphi_1 + i \varphi_2 \quad ( i : \text{imaginary unit} )$$

Dimensionless parameters are used for the description.



$d$  : mean water depth

$T$  : wave period

$$\lambda_0 = g \div \omega^2$$

$\sigma$  : mean bottom slope



2.01 Concepts of vector analysis

The description of the waves will be presented using the wave potential as the characterising parameter.

For this reason the necessary concepts from vector analysis are presented, leading to the description of the potential.

- a vectorfield may be defined by  $r_3 : (x, y, z)$
- a gradient is the vector, at location  $(x, y, z)$  which shows the direction of largest increment of a scalar-field " $\Phi$ "; the length of the vector indicates the size of this largest increment.

$$\text{grad } \Phi = \nabla \Phi = \left\{ \frac{\partial \Phi}{\partial x}, \frac{\partial \Phi}{\partial y}, \frac{\partial \Phi}{\partial z} \right\} \quad (2.01.1)$$

- the divergence of a vectorfield is a scalar which, at a location  $(x, y, z)$ , indicates the strength of the source (or: sink) at that point of the vectorfield.

$$\begin{aligned} \text{div } \vec{\Phi} &= \nabla \cdot \vec{\Phi} = \left( \frac{\partial v_1}{\partial x} + \frac{\partial v_2}{\partial y} + \frac{\partial v_3}{\partial z} \right) = \\ &= \lim_{\text{vol} \rightarrow 0} \frac{1}{\text{vol}} \oint \vec{\Phi} \cdot d\vec{n} \end{aligned} \quad (2.01.2)$$

- the rotation of a vectorfield is a vector which, at a location  $(x, y, z)$ , indicates the intensity and direction of the change of orientation of the vectorfield.

$$\begin{aligned} \text{rot } \vec{\Phi} &= \nabla \times \vec{\Phi} = \lim_{\text{area} \rightarrow 0} \frac{1}{\text{area}} \oint \vec{\Phi} \cdot d\vec{s} = \\ &= \left\{ \frac{\partial v_3}{\partial y} - \frac{\partial v_2}{\partial z}; \frac{\partial v_1}{\partial z} - \frac{\partial v_3}{\partial x}; \frac{\partial v_2}{\partial x} - \frac{\partial v_1}{\partial y} \right\} \end{aligned} \quad (2.01.3)$$



- for a field without sources or sinks:

$$\nabla \cdot \vec{v} = 0 \quad ( 2.01.4 )$$

- for a field without rotation:

$$\nabla \times \vec{v} = 0 \quad ( 2.01.5 )$$

- " nabla " is the vector operator:

$$\nabla \doteq \left\{ \frac{\partial}{\partial x} , \frac{\partial}{\partial y} , \frac{\partial}{\partial z} \right\}$$

## 2.02 Potential flow

A fluid field may be characterised by its velocity vectorfield. More convenient than the description as a vectorfield is a description using a scalarfield, " Potential ", using equation ( 2.01.1 ):

$$\vec{v} = \nabla \phi$$

here  $\vec{v}$  is: the velocity vectorfield.

For a velocity field which is expressed as the gradient of the potential the condition of a rotation free field, ( 2.01.5 ), will always be fulfilled:

$$\begin{aligned} \nabla \times \vec{v} &= \nabla \times (\nabla \phi) \doteq \left\{ \frac{\partial}{\partial y} \frac{\partial \phi}{\partial z} - \frac{\partial}{\partial z} \frac{\partial \phi}{\partial y} ; \frac{\partial}{\partial x} \frac{\partial \phi}{\partial z} - \frac{\partial}{\partial z} \frac{\partial \phi}{\partial x} ; \frac{\partial}{\partial x} \frac{\partial \phi}{\partial y} - \frac{\partial}{\partial y} \frac{\partial \phi}{\partial x} \right\} = \\ &\equiv \vec{0} \end{aligned}$$



For a velocity field which is expressed as the gradient of a potential and which is free of sources and sinks the condition will be:

$$\nabla \cdot \nabla \phi = 0$$

That is: a fluid field without sources or sinks is characterised by:

$$\nabla \cdot \nabla \phi = \nabla^2 \phi = 0$$

the " Laplace equation ".

This description is used for the fluid field of this study:

The velocity vectorfield is described using a " potential ":

$$\vec{v} = \nabla \phi \quad ( 2.02.1 )$$

The fluid field is assumed to be free of sources and sinks :

$$\nabla^2 \phi = 0 \quad ( 2.02.2 )$$

This is the field condition.

The next section presents the boundary conditions.



2.03

Boundary conditions of the field

Any field is determined by field condition and the boundary conditions.

Disregarding ( untill section 2.09 ) the 'side' boundaries, the boundary conditions are:

the - linearised - free surface condition ( see appendix 1 ):

$$\frac{\partial \phi}{\partial z} + \frac{1}{g} \frac{\partial^2 \phi}{\partial t^2} = 0 \quad \text{at } z = 0 \quad ( 2.03.1 )$$

and the rigid impermeable bottom condition:

$$\frac{\partial \phi}{\partial n} = 0 \quad \text{at } z = d \quad ( 2.03.2 )$$

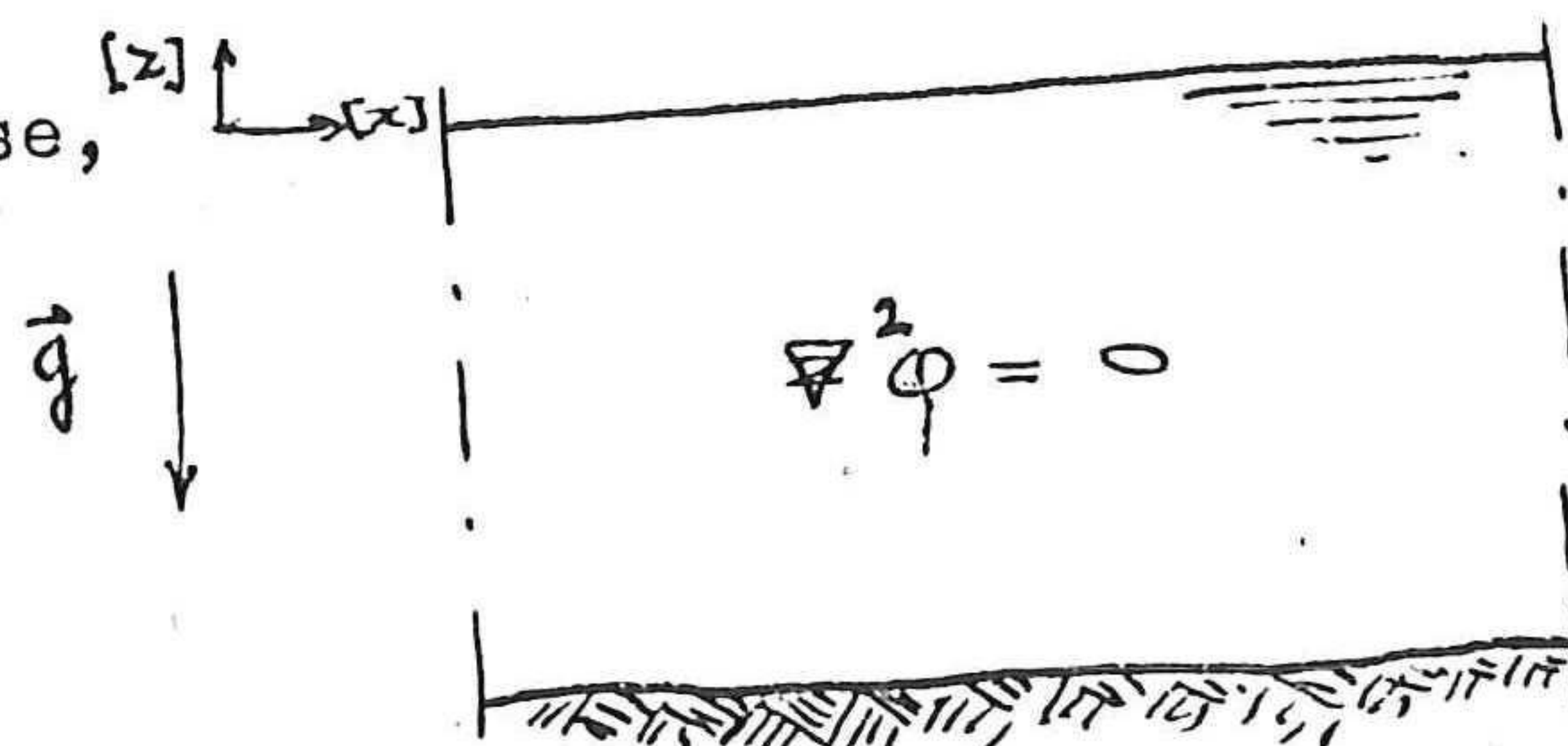
d = waterdepth

2.04

Side View Model

Combining the field condition ( section 2.02 ) with the boundary conditions of the field ( section 2.03 ) and disregarding the variation in the " y-direction " a set of equations is obtained which will be called the " Side View Model " :

visualise,



$$\frac{\partial \phi}{\partial z} + \frac{1}{g} \frac{\partial^2 \phi}{\partial t^2} = 0$$

$$\frac{\partial \phi}{\partial n} = 0$$

( 2.04.1 )

where:  $\nabla = \left\{ \frac{\partial}{\partial x}, \frac{\partial}{\partial z} \right\}$  ;  
 $\phi = \phi ( x, z, t )$



2.05

Flat bottom solution.

An analytical solution is known for the Side View Model for the case of a flat horizontal bottom. ( $\frac{\partial \varphi}{\partial z} = 0$ ). In deriving this solution it has been assumed that the potential is time-space periodical.

Lamb (1975) shows that this solution is:

$$\varphi = \frac{g a}{\omega} \frac{\cosh (k (h+z))}{\cosh (k h)} \cdot e^{i(\omega t - kx)} \quad (2.05.1)$$

- where:
- g : gravitational gradient intensity
  - a : amplitude
  - $\omega$  : wave time frequency: " wave frequency "
  - k : wave space frequency: " wave number "
  - h : waterdepth
  - z : depth below the water surface
  - t : time
  - i : unit of the imaginary number set

The displacement at the surface is:

$$\eta = i a \cdot e^{i(\omega t - kx)} \quad (2.05.2)$$

There is a relation between the wave time frequency and the wave space frequency determined by the "elasticity" of the surface:

$$\omega^2 = g k \cdot \tanh (k h) \quad (2.05.3)$$



The potential "  $\varphi$  " has three main components:

- 1)  $\frac{g a}{\omega}$  an amplitude factor depending on  $\omega$  ( 2.05.4 )
- 2)  $\frac{\cosh(k(h+z))}{\cosh(kh)}$  a distribution factor ( over the vertical ) ( 2.05.5 )
- 3)  $e^{i(\omega t - kx)}$  a phase factor ( 2.05.6 )

This knowledge will be used in the formulation of the Top View Model.

## 2.06 Top View Model

In case the bottom variation is 'very small' ( as assumed by Berkhoff ) the solution for the case of a flat bottom ( 2.05.1 ) will be approximately valid:

$$\varphi \langle x, y, z, t \rangle = \varphi_0 \langle x, y \rangle \cdot \frac{\cosh(k(h+z))}{\cosh(kh)} \cdot e^{i(\omega t - kx)} \quad ( 2.06.1 )$$

where:  $\varphi_0 \langle x, y \rangle$  is a first order approximation of  $\varphi \langle x, y, z \rangle$  in case of a nearly flat bottom.

Berkhoff (1976) shows that the set of equations ( 2.04.1 ) then reduces to: ( see appendix 3 for a short derivation ):

$$\nabla \cdot (c c_g \nabla \varphi_0) + k^2 c c_g \varphi_0 = 0 \quad ( 2.06.2 )$$

where:  $c$  : wave celerity  
 $c_g$  : wave group celerity  
 $\varphi_0$  : wave potential at the surface (  $z=0$  )  
 $k$  : wave number

$$\nabla = \left\{ \frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right\} \quad \text{and,}$$



$$c = \frac{g}{k} \quad ( 2.06.3 )$$

$$c_g = \frac{1}{2} \left( 1 + 2 \frac{kh}{\sinh(2kh)} \right) * c \quad ( 2.06.4 )$$

$$\omega^2 = gk \cdot \tanh(kh) \quad ( 2.06.5 )$$

Field equation ( 2.06.2 ), together with the boundary conditions of the field, will be called the " Top View Model " :

visualise,

$$\nabla \cdot (cc_g \nabla \phi_0) + k^2 cc_g \phi_0 = 0 \quad ( 2.06.6 )$$

where:  $\nabla = \left\{ \frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right\}$

### 2.07 Comparison

Two methods are now available for the prediction of the values of the potential throughout the wave field:

- ( 2.04.1 ) : the Side View Model, based on the Laplace equation.
- ( 2.06.6 ) : the Top View Model, based on the refraction - diffraction equation.

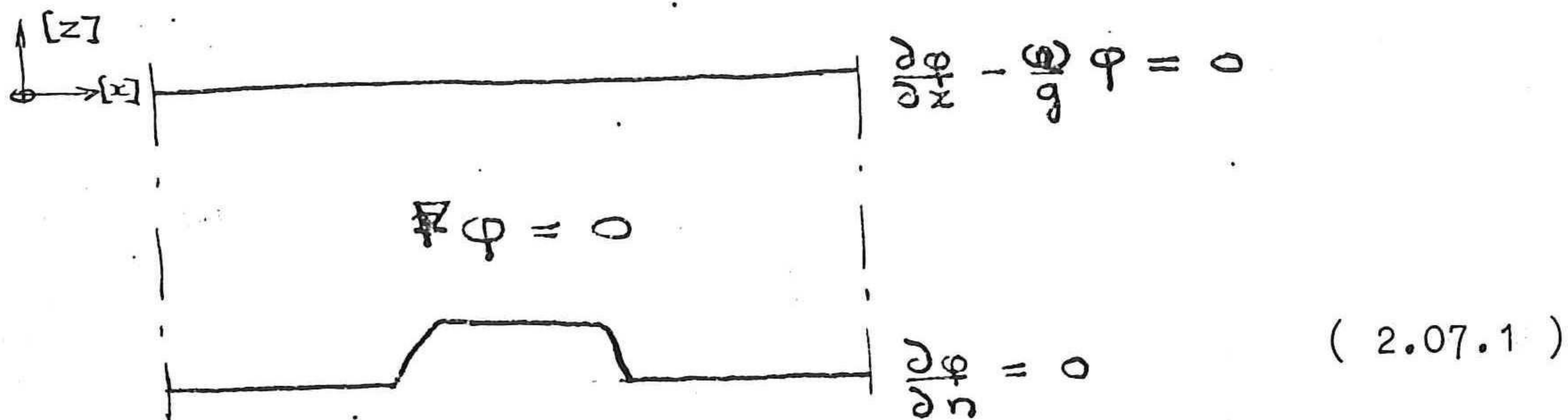
Since the time displacement factor,  $e^{i(\omega t)}$ , is known and the same for both models, this factor may be isolated and filtered out of the description:



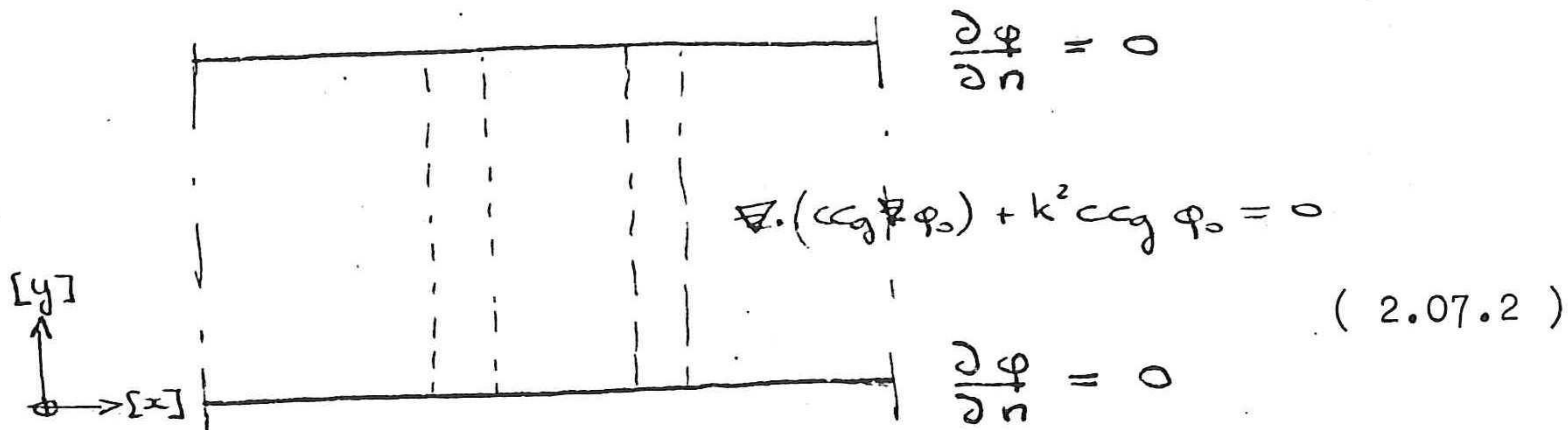
The factor  $e^{i(\omega t)}$  is further disregarded.

The two models which will be compared are:

Side View Model:



Top View Model:



These two models should both be able to predict the values of the wave potential in the field.

The Side View Model and the Top View Model differ in the following aspects:

- The Side View Model provides the values of the wave potential in a vertical plane whereas the Top View Model provides them for a horizontal plane: the water surface.

This has the practical consequence that, for the comparison, only the common surface values have to be compared.



- The derivation of the refraction - diffraction equation ( see appendix 3. ) is based on the assumption that the variation of the bottom will be 'very small' . The Top View Model is based on this equation and thus may be sensitive to the bottom slope in its predictions; the Side View Model works for any bottom slope. This difference forms the basis of this study which is meant to provide insight in this particular characteristic of the Top View Model.

- The Top View Model is capable of dealing with variations in the plane of the horizontal which the Side View Model can not. ( A 3-D model could, )

This has the practical consequence that, for the comparison, variations to the side of the common axis are disregarded. ( The Top View Model is thus used essentially one-dimensional This is the consequence of the aim to use the model for investigations of variations in the horizontal plane after succesful completion of the study. Unfortunately this has not been possible and, in retrospect, a one-dimensional model would have sufficed.)

- The Top View Model is based upon a distribution of the potential over the vertical as described by ( 2.06.1 ); The distribution of the potential over the vertical has to be prescribed for the Side View Model.

This has the practical consequence that, for the comparison, at the influx boundary the potential of the Top View Model is also prescribed for the Side View Model.



Practical aspects of the comparison:

- The bottom configurations must be the same for both models per tested prediction of the potential.
- The potential distribution over the vertical is made the same at the influx boundary of both models.
- The influx boundaries should be sufficiently far removed from the bottom section with the variation of the bottom slope.
- Predictions of the potential at the surface will be compared.

2.08            A model for comparative calculation

The important difference between the Side View Model and the Top View Model is the implied dependency on the bottom slope of the Top View Model.

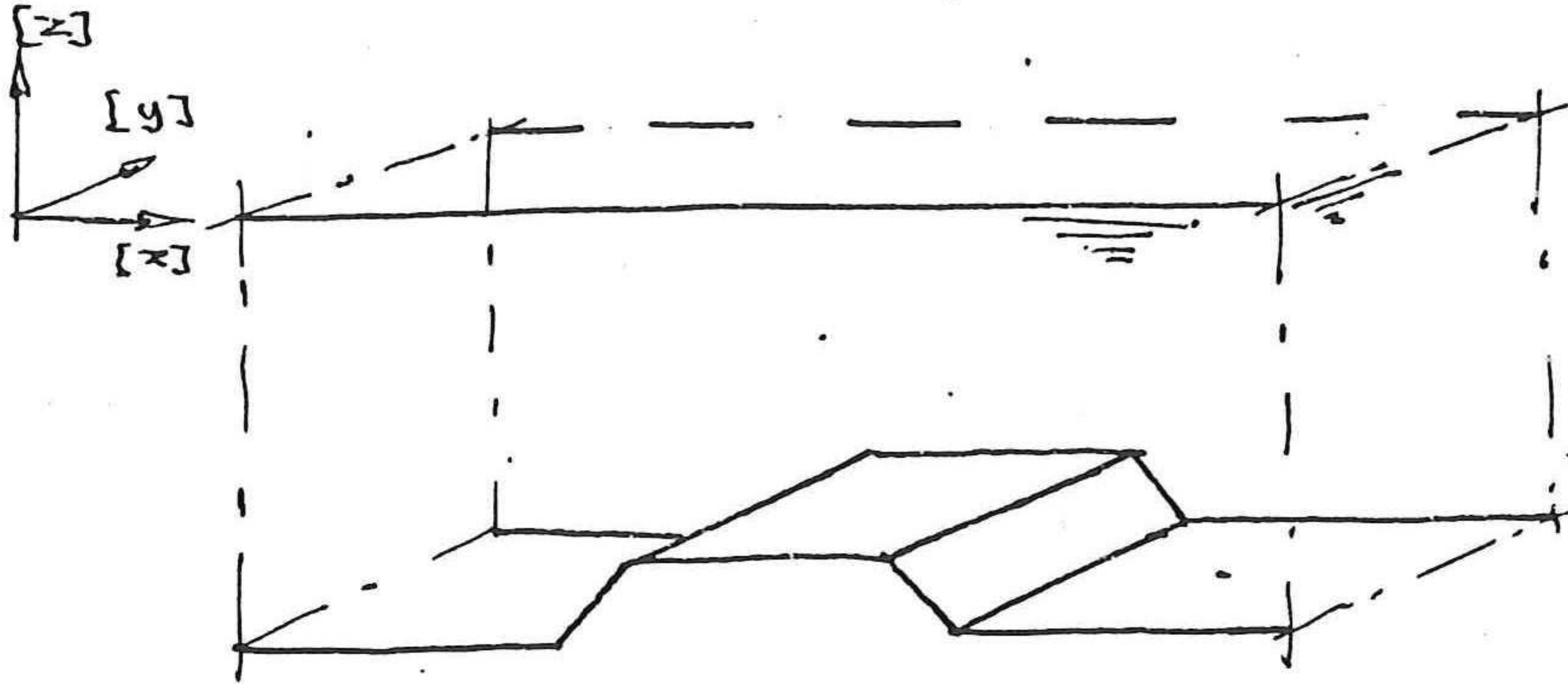
The bottom contour which is to be used for the comparison of both models must therefore have a section characterised by the bottom slope which may be varied.

A very simple bottom contour has been selected which fits this description:

A trapezoidal submerged barrier on a flat, horizontal bottom.

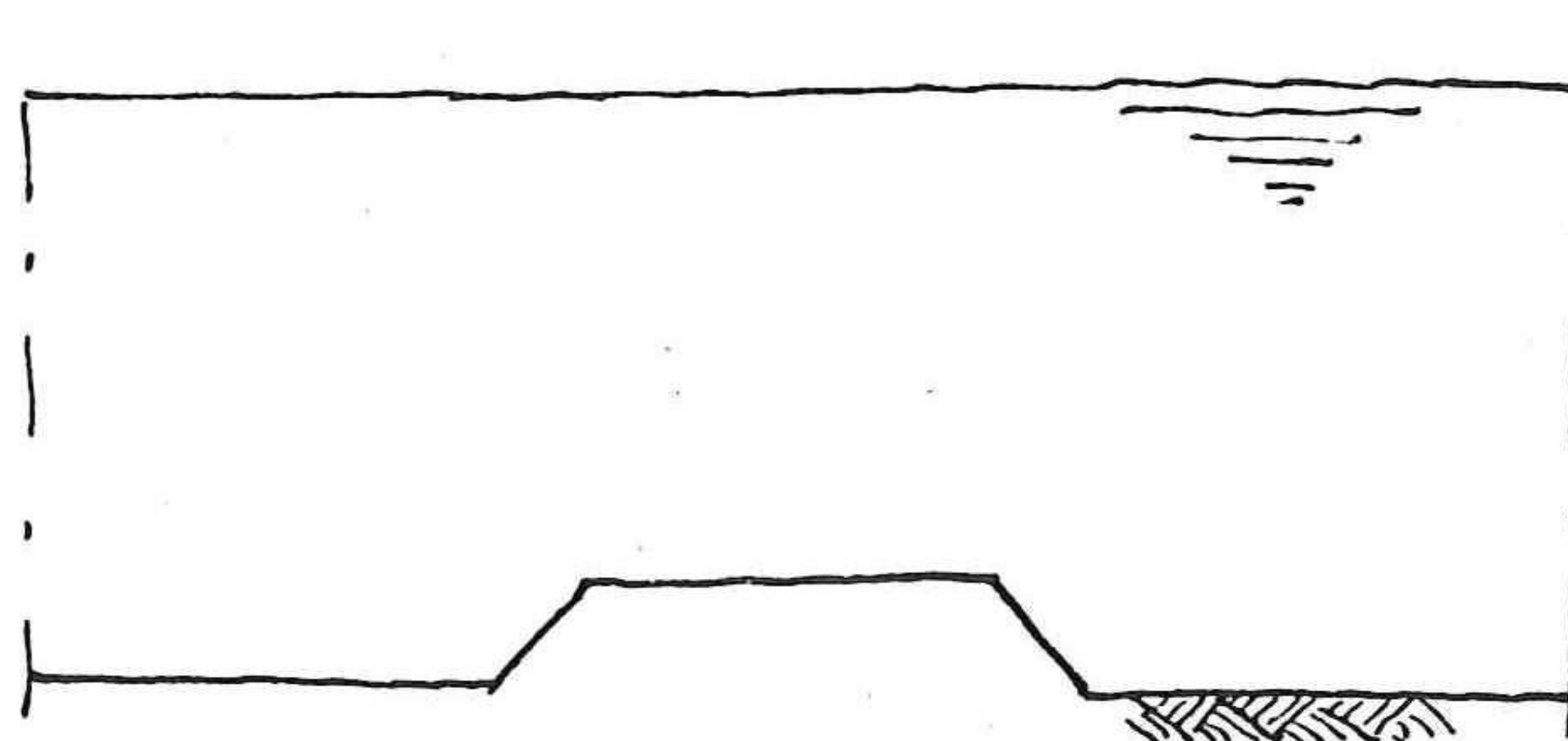


Visualise the trapezoidal submerged barrier:



( 2.08.1 )

The Side View Model uses:



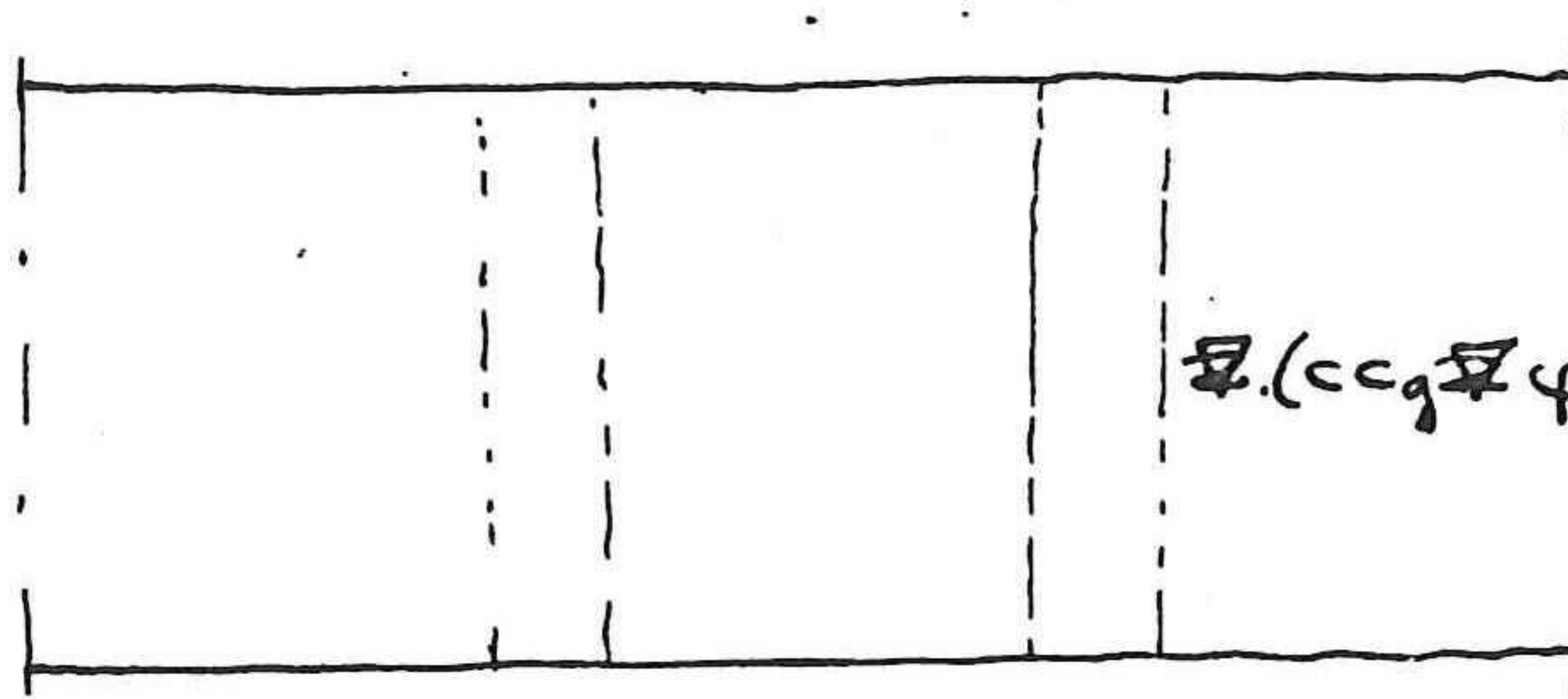
$$\frac{\partial \phi}{\partial z} - \frac{g}{g} \phi = 0$$

$$\nabla^2 \phi = 0$$

$$\frac{\partial \phi}{\partial n} = 0$$

( 2.08.2 )

The Top View Model uses:



$$\frac{\partial \phi}{\partial n} = 0$$

$$\nabla \cdot (c c_g \nabla \phi) + k^2 c c_g \phi = 0$$

$$\frac{\partial \phi}{\partial n} = 0$$

( 2.08.3 )

where the values of  $c$ ,  $c_g$ , and  $k$  have to be provided throughout the field.

For the side boundaries we have  $\frac{\partial \phi}{\partial n} = 0$ , because variation in the  $y$ -direction is not considered.

The influx and efflux boundaries still have to be considered. This will be done next.



2.09 The influx and efflux boundaries

Up to this point two boundaries have not yet been considered: the influx and efflux boundaries, where the incoming and outgoing waves are to be described.

Both the Side View Model and the Top View Model require this information.

In this study only simple harmonic waves are considered.

These are described by:

$$\varphi = \hat{\varphi} \cdot e^{i(\omega t - kx)} \quad ( 2.09.1 )$$

( Recall that the time variation is disregarded. )

outgoing wave

Any wave coming from within the field must be able to pass freely through the boundary.

At any place within the field:

$$\frac{\partial \varphi}{\partial x} = -i k \hat{\varphi} \cdot e^{i(\omega t - kx)} \quad ( 2.09.2 )$$

that is:

$$\frac{\partial \varphi}{\partial x} = -i k \varphi \quad ( 2.09.3 )$$

This should also be the case for the outgoing waves at the boundary.

By using  $\frac{\partial \varphi}{\partial n}$  instead of  $\frac{\partial \varphi}{\partial x}$  the orientation of the efflux is in all cases correctly accounted for in the equations:

$$\frac{\partial \varphi}{\partial n} = -i k \varphi \quad ( 2.09.4 )$$



This is called the " radiation condition ":  $\frac{\partial \varphi}{\partial n} = -i k \varphi$ .

With the radiation condition imposed at the boundary waves can freely leave the field.

The radiation condition holds for both the Side View Model and the Top View Model.

incoming wave

It is assumed that a wave enters the field at the left-hand boundary. ( See ( 2.08.2 ) and ( 2.08.3 ) ).

At that boundary, at any time, two waves may exist together: an incoming wave and an outgoing wave:

$$\varphi = \varphi_i + \varphi_r \quad ( 2.09.5 )$$

where:  $\varphi_i$  : incoming wave  
 $\varphi_r$  : outgoing wave

The outgoing wave,  $\varphi_r$ , can in the computation be determined only from:

$$\varphi_r = \varphi - \varphi_i \quad ( 2.09.6 )$$

This outgoing wave has to fulfill the radiation condition ( 2.09.4 ):

$$\begin{aligned} \frac{\partial \varphi_r}{\partial n} &= -i k \varphi_r \\ \Rightarrow \frac{\partial \varphi}{\partial n} - \frac{\partial \varphi_i}{\partial n} &= -i k \varphi + i k \varphi_i \end{aligned}$$

$\frac{\partial \varphi_i}{\partial n}$  may be determined from the radiation condition for the incoming wave ( for which the direction of travel is opposite to the normal at the boundary ) :

$$\begin{aligned} \frac{\partial \varphi_i}{\partial n} &= +i k \varphi_i \\ \Rightarrow \frac{\partial \varphi}{\partial n} &= -i k \varphi + 2 i k \varphi_i \quad ( 2.09.7 ) \end{aligned}$$



Taking  $\hat{\varphi}_i = 1$  the influx condition becomes,

for the Top View Model:

$$\frac{\partial \varphi}{\partial n} = -ik\varphi + 2ik \quad ( 2.09.8 )$$

for the Side View Model:

$$\frac{\partial \varphi}{\partial n} = -ik\varphi + 2ik \frac{\cosh(k(h+z))}{\cosh(kh)} \quad ( 2.09.9 )$$

( As has been stated in section 2.07 the Side View Model has to be given the same distribution of the potential over the vertical as is implied for the Top View Model, at the influx boundary. At this boundary the bottom must be horizontal. )

The choice of  $\hat{\varphi}_i = 1$  does not affect the generality of the description due to the linearity of the description.

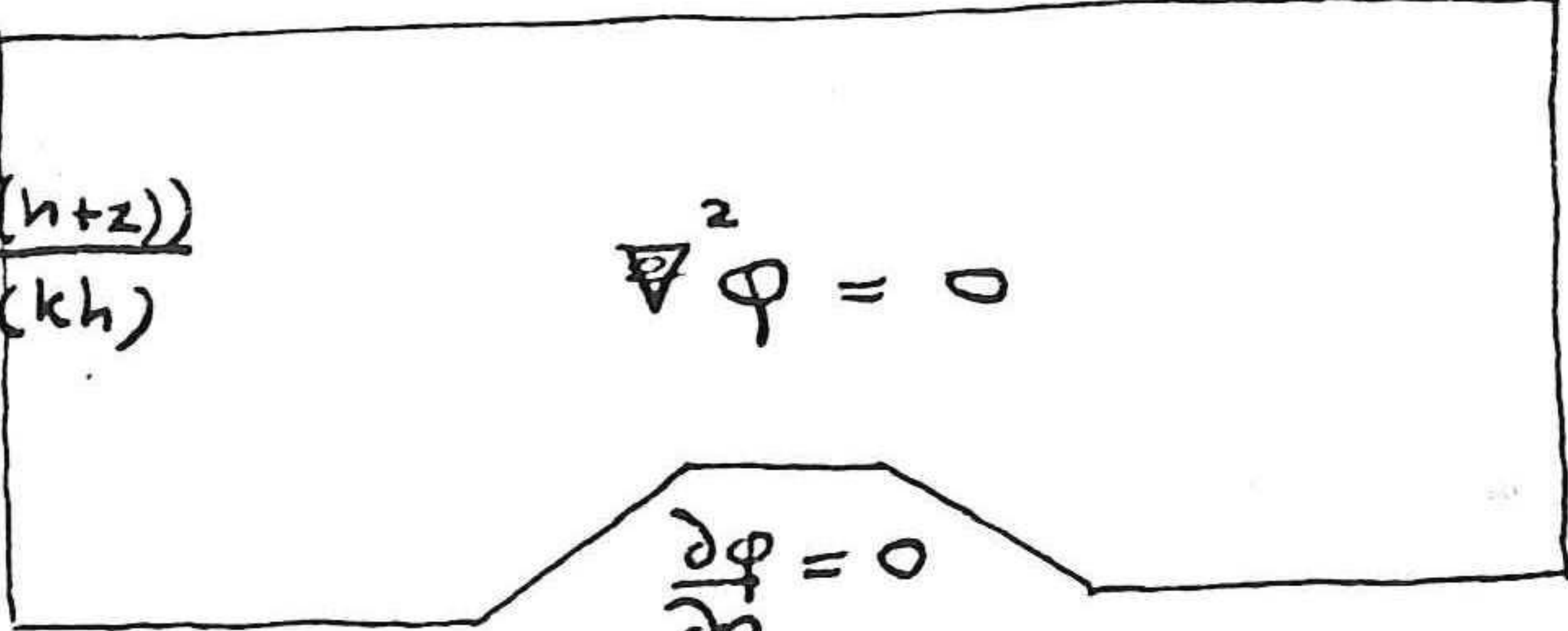




2.10 Summary of the analytical description

These are the two models which will be used for the comparative calculation:

Side View Model :

$$\frac{\partial \phi}{\partial n} = \frac{\omega^2}{g} \phi$$


$$\frac{\partial \phi}{\partial n} = -ik\phi + 2ik\hat{\phi} \frac{\cosh(k(h+z))}{\cosh(kh)}$$

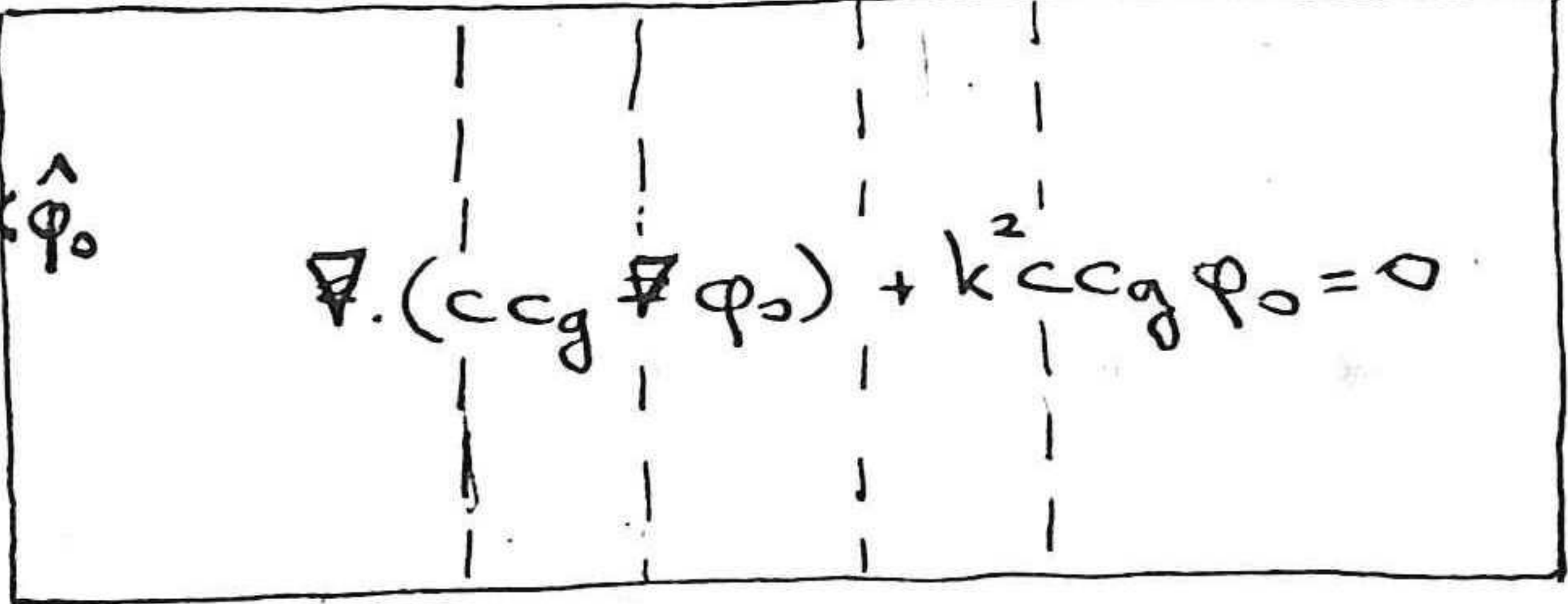
$$\nabla^2 \phi = 0$$

$$\frac{\partial \phi}{\partial n} = -ik\phi$$

$$\frac{\partial \phi}{\partial n} = 0$$

( 2.10.1 )

Top View Model :

$$\frac{\partial \phi_0}{\partial n} = 0$$


$$\frac{\partial \phi_0}{\partial n} = -ik\phi_0 + 2ik\hat{\phi}_0$$

$$\nabla \cdot (cc_g \nabla \phi_0) + k^2 cc_g \phi_0 = 0$$

$$\frac{\partial \phi_0}{\partial n} = -ik\phi_0$$

$$\frac{\partial \phi_0}{\partial n} = 0$$

( 2.10.2 )

( To ensure that the Top View Model uses the same bottom configuration as the Side View Model the values of  $c$ ,  $c_g$  and  $k$ , as required for the Top View Model, will be determined using the values of the local depth in the Side View Model. )



3.00      Introduction

The sets of equations ( 2.10.1 ), the Side View Model, and ( 2.10.2 ), the Top View Model, need to be written in numerical form, in order to prepare them for computation in the computer.

The translation from analytical to numerical form is done in this chapter.

The resulting sets of equations are: ( 3.07.1 ) and ( 3.07.2 ) .

The numerical form of the description is based on the Finite Element Package ( " AFEP " ) which has been developed at the T.H. Delft.

In order to provide some background information for the Finite Element Method, the chapter starts off with a very brief general description of the concepts of the Finite Element Method, followed by a description of Lagrange Polynomials and the Galerkin Method, used in AFEP.

↔



3.01 Finite Element Method

This section presents a very brief description of the concepts of the Finite Element Method.

For a more complete covering consult any of the many books on the topic, such as Zienkiewicz (1971), van Kan & Segal (1979), Strang & Fix (1973).

The Finite Element Method is based on "elements", which are small subregions of the field in which the computation is to be done.

The elements are generated by 'covering' the field with a mesh.

The unknowns are determined in the nodes of the mesh. By interpolating between the nodes, using an assumed interpolation function, an approximation of the continuous description can be achieved.

Depending on the total number of nodes and their distribution it is possible, within certain limits, to determine a balance between accuracy and laboriousness.

Basically the Finite Element Method works as follows ( Segal, 1975 ):



- a finite continuous region is divided into subregions called " elements ".
- nodal points are selected within each element.
- it is indicated which unknowns are associated with each nodal point. It is not necessary that these be the same at each nodal point.
- an approximate solution generating algorithm using the selected unknowns is provided for each element.
- the approximate solution, expressed in the selected unknowns, is determined for each element.
- the contributions from each element to the system of equations are stored in an element matrix resp. element vector.
- the system of equations is constructed by 'adding' the contributions per element together.
- the system of equations is solved for the unknowns, thereby providing the solution.

The essential part of this procedure is the description of the algorithm for the approximate solution per element. In doing this use can be made of a method of weighted residuals ( Finlayson, 1972 ) in combination with interpolation functions.

In AFEP Lagrange polynomials of the first kind are used for the interpolation but also as weighting functions. This kind of method of weighted residuals is called the " Galerkin Method ". Both the Galerkin Method and the Lagrange polynomials will now be described.



3.02      The Galerkin Method

The Galerkin Method is a method of weighted residuals which uses the same functions for the weighting as for interpolation.

The method of weighted residuals aims at reducing the residual which results when an operation is performed in discrete space instead of in continuous space.

Let the operation in continuous space be :

$$L \cdot \varphi = 0 \quad ( 3.02.1 )$$

where: L : linear operator.

In discrete space this will generally be:

$$L \tilde{\varphi} = R \quad ( 3.02.2 )$$

where:  $\tilde{\varphi}$  : the approximation of  $\varphi$

R : residual

In order to minimalise the residual ( that is: to obtain a better approximation of the operation in the discrete form ) a weighting function is introduced in the operation.

A requirement for this weighting function is that it nullifies the residual, at least averaged over an interval " F ":

$$\int_F R \cdot W \cdot dF = 0 \quad ( 3.02.3 )$$

where: W : weighting function.



The operation then becomes:

$$\int_F L. \bar{\varphi} . W . dF \quad ( 3.02.4 )$$

In discrete form, continuous functions are described using interpolation functions between the nodes.

The Galerkin Method uses the same function  $\Phi_j$  as weighting function (  $W =: \Phi_j$  ) and as interpolation function (  $\bar{\varphi} =: \varphi \cdot \Phi_i$  ).

Using the Galerkin Method the operation will be:

$$\int_F L. \bar{\varphi} . \Phi_j . dF \quad ( 3.02.5 )$$

where:  $F$  : interval of the weighting operation  
 $\Phi$  : interpolation polynomial  
 $j$  : nodal index  
 $L$  : linear operator

In this study the Galerkin method is used in the translation of the analytical description of the models into numerical form.

The interpolation is done using Lagrange Polynomials.

### 3.03 Lagrange Polynomials

A continuous function can be approximated by a discrete function using nodal points ( where the values of the unknowns are determined ) and interpolation polynomials.



A general interpolation formula is the Lagrange Interpolation Formula ( Abramowitz, 1970)

$$f(x) = \sum_0^n l_i(x) \cdot f(x_i) + R_n(x)$$

where:

$$l_i = \frac{\prod_n(x)}{(x-x_i) \prod_n(x_i)}$$

$$= \frac{(x-x_0) \dots (x-x_{i-1})(x-x_{i+1}) \dots (x-x_n)}{(x_i-x_0) \dots (x_i-x_{i-1})(x_i-x_{i+1}) \dots (x_i-x_n)}$$

$$R_n = \prod_n(x) \cdot \frac{f^{n+1}(\xi)}{(n+1)!} \quad : \text{ the remainders}$$

In this study only the linear Lagrange Polynomial will be used:

$$l_i(x) = \frac{x-x_{i-1}}{x_i-x_{i-1}}$$

: the Lagrange Polynomial at point  $x_i$ .

Likewise for point  $x_{i-1}$  :

$$l_{i-1}(x) = \frac{(x_i-x)}{(x_i-x_{i-1})}$$

Visualise the linear Lagrange Polynomials:



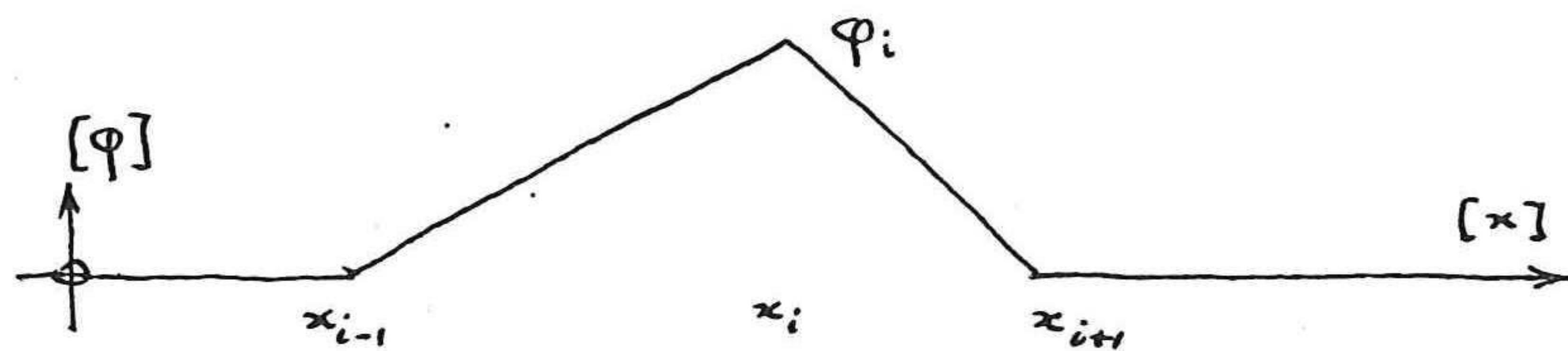
The interpolation function,  $x$ , is composed of linear Lagrange Polynomials.

For brevity  $\Phi_i(x)$  will be written as  $\Phi_i$  , etc.



Describing a continuous function using the linear Lagrange Polynomials for the approximation is done as follows:

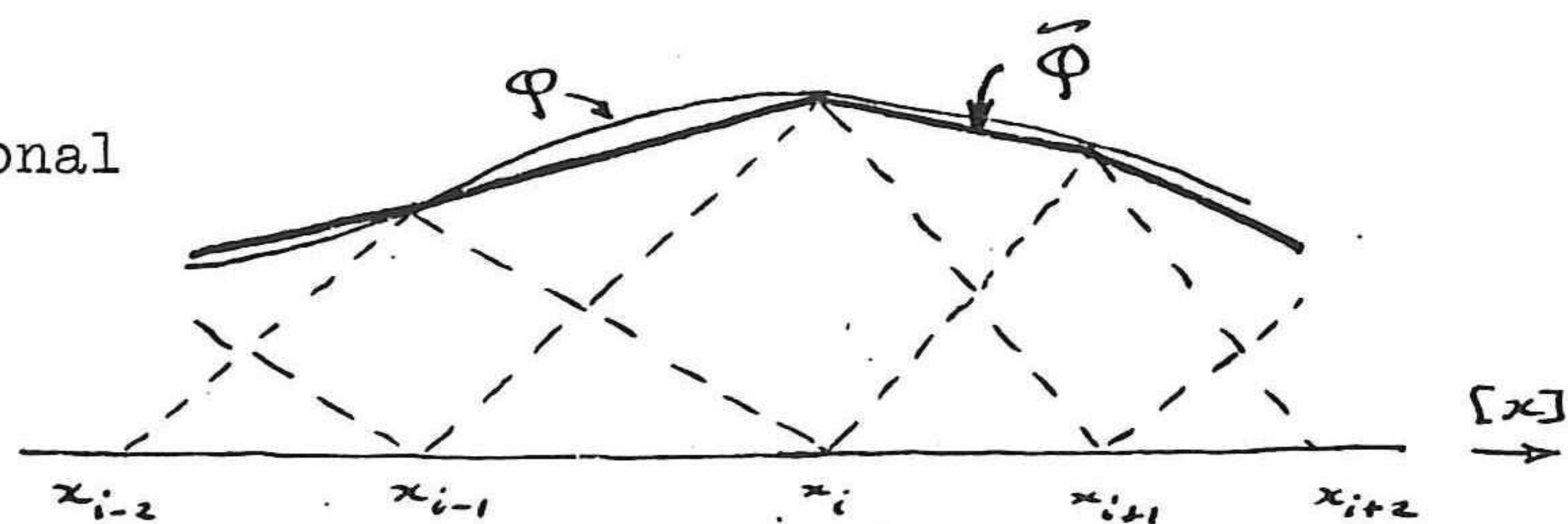
- at the nodes the value of the function is used.
- between the nodes interpolation is used, using the linear Lagrange Polynomials.
- the description over the interval will be composed of the local approximated descriptions using superposition.
- the local description will have the following form:



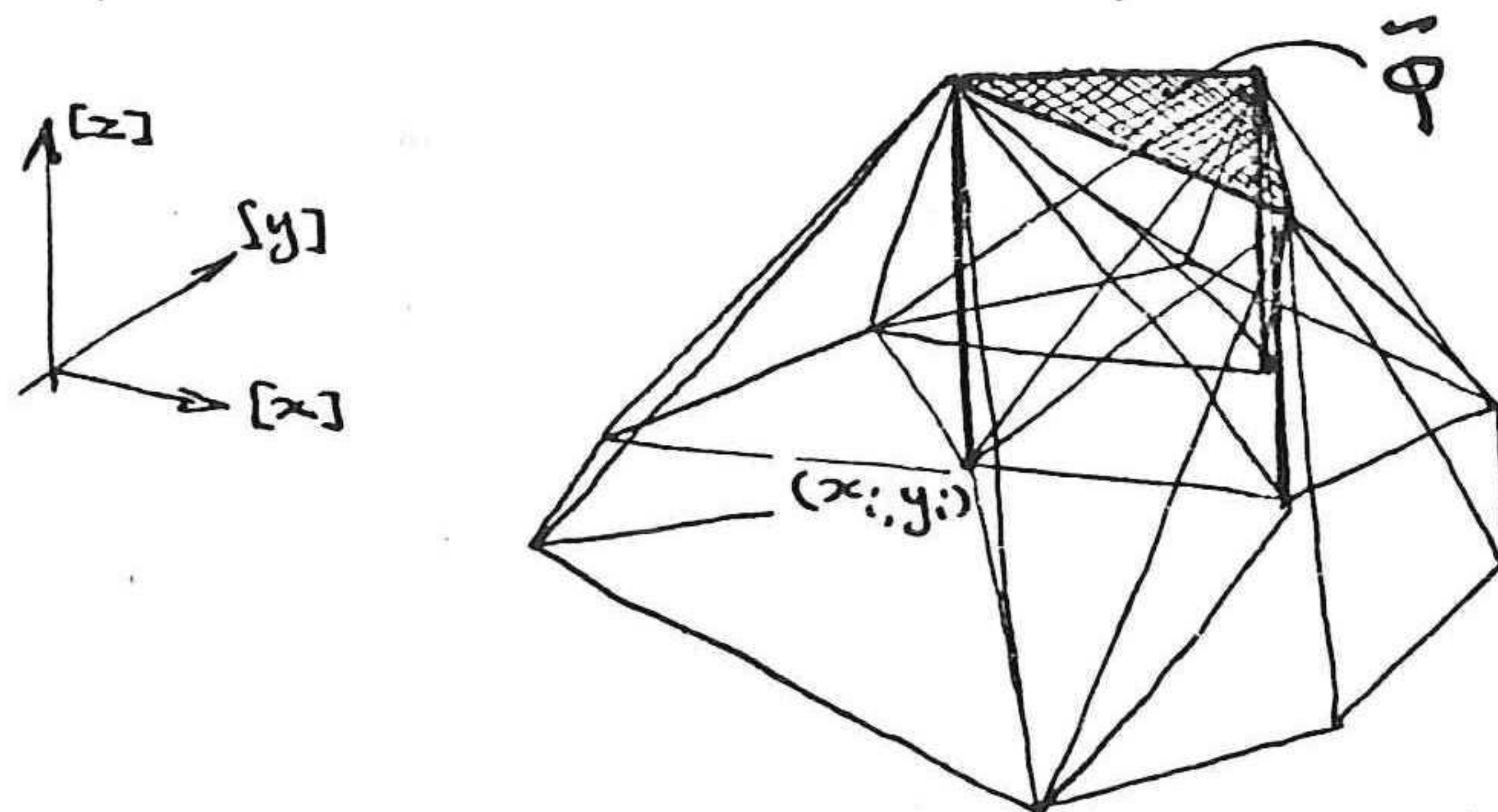
where:  $\varphi_i$  : the value of  $\varphi$  at point  $x_i$  (  $x_i = \varphi(x_i)$  )

Visualise the result of an approximation done this way:

one dimensional



two dimensional ( the local description uses  $\varphi_i(x, y)$  )



In dealing with approximations done this way the generalised differential of the Sobolev space is to be used, due to non-differentiability of the function  $\tilde{\varphi}$  at the



nodes and boundaries of the element. See Segal (1974) for a more detailed description.

3.04 Algorithm for the translation

The following is a brief outline of the method for rewriting the analytical equations ( ( 2.10.1 ), ( 2.10.2 ) ) into numerical form ( ( 3.07.1 ), ( 3.07.2 ) ):

- 1)- Rewrite the field equations using the Galerkin Method ( 3.02.5 ) using linear Lagrange Polynomials for the approximative description in discrete form.

$$L\phi = 0 \quad \rightarrow \quad \int_F L \bar{\phi} \phi_i dF$$

- 2)- Apply Gauss' theorem:

$$\text{e.g.:} \quad \int_F \nabla^2 \bar{\phi} \phi_i dF \quad \Leftrightarrow \quad - \int_F \nabla \bar{\phi} \nabla \phi_i dF + \int_F \nabla(\bar{\phi} \phi_i) dF$$

- 3)- Apply Greens' theorem:

$$\int_F \nabla(\bar{\phi} \phi_i) dF = \int_{\square} \frac{\partial \bar{\phi}}{\partial n} \phi_i d\alpha$$

- 4)- Substitute the boundary conditions. ( In this case these are of the general form:  $\frac{\partial \phi}{\partial n} = -a\phi + q$  ).
- 5)- Provide the computer with an algorithm for the evaluation of the integral.
- 6)- Store, separately, the contributions from each element to the system of equations ( as described in section 3.01 ) to the element matrix resp. the element vector.



Due to the difference between field elements and boundary elements the contributions from the field elements and the boundary elements are treated differently.

- 7)- Assemble all the contributions to the matrix and the vector and solve the system  $\vec{A} \vec{\varphi} = \vec{g}$   
Standard methods may be used for this purpose.

These steps will be applied to the equations ( 2.10.1 )  
and ( 2.10.2 ).

### 3.05 Describing the equations in numerical form

In the following the Side View Model and the Top View Model will be considered together so as to show the similarities between the operations performed on the two models.

The field equations are:

Side View Model:  $-\nabla^2 \varphi = 0$

Top View Model :  $-\nabla \cdot (c c_g \nabla \varphi_0) - k^2 c c_g \varphi_0 = 0$

( The negative sign is chosen in order to adhere to AFEP usage. )



1) - Apply Galerkins Method:

$$\text{S.V.M.} \quad - \varphi_i \int_F \nabla^2 \phi_i \phi_j dF = 0$$

$$\text{T.V.M.} \quad - \varphi_i \int_F \{ \nabla \cdot (c c_g \nabla \phi_i) - k^2 c c_g \phi_i \} \phi_j dF = 0$$

In this description use has been made of the property of interpolation by linear Lagrange Polynomials:

$$\tilde{\varphi} = \varphi(x_i) \cdot \phi_i =: \varphi_i \phi_i \quad (\text{summation being implied})$$

From now on  $\nabla$  is written as  $\nabla$ , the two dimensionality already being implied in the integration sign.

2) - Apply the Gauss' theorem and Green's theorem.

3) -

$$\text{S.V.M.} \quad + \varphi_i \int_F (\nabla \phi_i) (\nabla \phi_j) dF - \oint_{\square} \frac{\partial}{\partial n} (\varphi_i \phi_j) \phi_j d\alpha = 0$$

$$\text{T.V.M.} \quad + \varphi_i \int_F (c c_g \nabla \phi_i \nabla \phi_j - k^2 c c_g \phi_i \phi_j) dF - \oint_{\square} c c_g \frac{\partial}{\partial n} (\varphi_i \phi_j) \phi_j d\alpha = 0$$

where:  $\square$  : boundary

4) - Substitute the general boundary conditions

$$\frac{\partial \varphi}{\partial n} = -a\varphi + g$$

$$\text{S.V.M.} \quad + \varphi_i \int_F (\nabla \phi_i) (\nabla \phi_j) dF - \int_{\square_k} (-a\varphi \phi_i + g) \phi_j d\alpha = 0$$

$$\text{T.V.M.} \quad + \varphi_i \int_F \{ c c_g (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) \} dF - \int_{\square_k} c c_g (-a\varphi \phi_i + g) \phi_j d\alpha = 0$$

( The boundary may be considered in parts and described in the summation in the equations. )



Separating the Real and Imaginary parts of the complex numbers:

( The Real part is given index " 1 ", the Imaginary part has index " 2 " . )

$$\begin{array}{l}
 \text{S.V.M.} \\
 \left. \begin{array}{l}
 + \varphi_{1i} \int_F (\nabla \phi_i)(\nabla \phi_j) dF - k \sum_{\square_k} (-a_{1i} \varphi_{1i} \phi_i + a_{2i} \varphi_{2i} \phi_i + g_{1i}) \phi_j d\alpha \\
 + \varphi_{2i} \int_F (\nabla \phi_i)(\nabla \phi_j) dF - \sum_{\square_k} (-a_{2i} \varphi_{1i} \phi_i - a_{1i} \varphi_{2i} \phi_i + g_{2i}) \phi_j d\alpha
 \end{array} \right\}
 \end{array}$$

T.V.M.

$$\left. \begin{array}{l}
 + \varphi_{1i} \int_F c_{ij} (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) dF - \sum_{\square_k} c_{ij} (-a_{1i} \varphi_{1i} \phi_i + a_{2i} \varphi_{2i} \phi_i + g_{1i}) \phi_j d\alpha \\
 + \varphi_{2i} \int_F c_{ij} (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) dF - \sum_{\square_k} c_{ij} (-a_{2i} \varphi_{1i} \phi_i - a_{1i} \varphi_{2i} \phi_i + g_{2i}) \phi_j d\alpha
 \end{array} \right\}$$

Rewriting these sets of equation using vector description

and the algorithm:

$$\begin{bmatrix} a & 0 \\ 0 & b \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} ax \\ by \end{bmatrix}$$

S.V.M.

$$\begin{bmatrix} \int_F \nabla \phi_i \nabla \phi_j dF + \sum_{\square_k} a_{1i} \phi_i \phi_j d\alpha & - \sum_{\square_k} a_{2i} \phi_i \phi_j d\alpha \\ \sum_{\square_k} a_{2i} \phi_i \phi_j d\alpha & \int_F \nabla \phi_i \nabla \phi_j dF + \sum_{\square_k} a_{1i} \phi_i \phi_j d\alpha \end{bmatrix} \begin{bmatrix} \varphi_{1i} \\ \varphi_{2i} \end{bmatrix} = \begin{bmatrix} \sum_{\square_k} g_{1i} \phi_j d\alpha \\ \sum_{\square_k} g_{2i} \phi_j d\alpha \end{bmatrix}$$



T.V.M.

$$\begin{pmatrix} \int_F c c_g (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) dF + \sum_{\alpha_k} c c_g (a_{ii} \phi_i \phi_j) d\alpha & - \sum_{\alpha_k} c c_g a_{2i} \phi_i \phi_j d\alpha \\ \sum_{\alpha_k} c c_g a_{2i} \phi_i \phi_j d\alpha & \int_F c c_g (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) dF + \sum_{\alpha_k} c c_g a_{ii} \phi_i \phi_j d\alpha \end{pmatrix} \begin{pmatrix} \phi_{ii} \\ \phi_{2i} \end{pmatrix} = \begin{pmatrix} \sum_{\alpha_k} c c_g g_{ii} \phi_j d\alpha \\ \sum_{\alpha_k} c c_g g_{2i} \phi_j d\alpha \end{pmatrix}$$

It is convenient to use the property of superposition to separate the matrices into parts.

This also has the advantage of allowing a separation into matrices describing the different aspects: the field condition versus the boundary conditions.

S.V.M. 
$$\begin{pmatrix} \int_F \nabla \phi_i \nabla \phi_j dF & 0 \\ 0 & \int_F \nabla \phi_i \nabla \phi_j dF \end{pmatrix} \begin{pmatrix} \phi_{ii} \\ \phi_{2i} \end{pmatrix} + \quad (3.05.1)$$

$$+ \begin{pmatrix} \sum_{\alpha_k} a_{ii} \phi_i \phi_j d\alpha & - \sum_{\alpha_k} a_{2i} \phi_i \phi_j d\alpha \\ \sum_{\alpha_k} a_{2i} \phi_i \phi_j d\alpha & \sum_{\alpha_k} a_{ii} \phi_i \phi_j d\alpha \end{pmatrix} \begin{pmatrix} \phi_{ii} \\ \phi_{2i} \end{pmatrix} = \begin{pmatrix} \sum_{\alpha_k} g_{ii} \phi_j d\alpha \\ \sum_{\alpha_k} g_{2i} \phi_j d\alpha \end{pmatrix}$$

T.V.M.

(3.05.2)

$$\begin{pmatrix} \int_F c c_g (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) dF & 0 \\ 0 & \int_F c c_g (\nabla \phi_i \nabla \phi_j - k^2 \phi_i \phi_j) dF \end{pmatrix} \begin{pmatrix} \phi_{ii} \\ \phi_{2i} \end{pmatrix} +$$

$$+ \begin{pmatrix} \sum_{\alpha_k} c c_g a_{ii} \phi_i \phi_j d\alpha & - \sum_{\alpha_k} c c_g a_{2i} \phi_i \phi_j d\alpha \\ \sum_{\alpha_k} c c_g a_{2i} \phi_i \phi_j d\alpha & \sum_{\alpha_k} c c_g a_{ii} \phi_i \phi_j d\alpha \end{pmatrix} \begin{pmatrix} \phi_{ii} \\ \phi_{2i} \end{pmatrix} = \begin{pmatrix} \sum_{\alpha_k} c c_g g_{ii} \phi_j d\alpha \\ \sum_{\alpha_k} c c_g g_{2i} \phi_j d\alpha \end{pmatrix}$$



It can be seen that for both models the first section is a ( symmetrical ) matrix describing the field condition in the element.

The second part is a ( non symmetrical ) matrix for the boundary elements.

The last part, the right hand side of the equation, is the vector of the boundary elements.

The three parts are called, respectively:

- the element matrix of the field,
- the element matrix of the boundary,
- the element vector of the boundary.

Note the similarity between the two models.

5)- Evaluate the integrals.

Equations ( 3.05.1 ) and ( 3.05.2 ) contain integrals.

That is an operation which has to be defined for the computer.

This can be done either by providing an algorithm which computes the value of the solution, or, in the case of a simple integration, by providing the solution directly. Because of the simplicity of the integration involved in this study the latter method has been used.

It is demonstrated how this is done using the boundary element. For the field element the integration is done analogously; this is not demonstrated.



Lagrange Polynomials are used for the description of the approximation of the unknowns in the element.

For a two-point boundary element the Lagrange Polynomial will have either of two forms:



The nodes have been labeled ① resp ②.

In an integral a multiple of combinations of  $\Phi_i$  may exist,

$$\text{e.g.: } \int_{\Omega} c c_g a_{1i} \Phi_i d\Omega = \int_{\Omega} c c_g a_1 (\Phi_k \Phi_i \Phi_j) d\Omega$$

(  $a_{1i}$  has been rewritten using Lagrange Polynomials. )

In the integration  $i, j, k$  may each take either value ① or ②. Thus a number of possible combinations of  $\Phi_1, \Phi_2$  may exist in the integration, depending upon the nodal point from which the computation is done (① or ②) and which combination is made to the other nodes within the element (① or ②).

As an example it is demonstrated how this is done for the evaluation of the integral  $\int_{\Omega} c c_g a_1 (\Phi_k \Phi_i \Phi_j) d\Omega$ .

A)-

For  $\Phi_k =: \Phi_1 ; \Phi_i =: \Phi_1 ; \Phi_j =: \Phi_1$

$$\begin{aligned} \int_{\Omega} c c_g a_1 (\Phi_1)^3 d\Omega &= \int_0^H c c_g a_1 \left(\frac{x}{H}\right)^3 dx = c c_g a_1 \left. \frac{x^4}{4H^3} \right|_0^H \\ &= c c_g a_1 \cdot \frac{H}{4} \end{aligned}$$

B)-

For  $\Phi_k =: \Phi_1 ; \Phi_i =: \Phi_2 ; \Phi_j =: \Phi_2$

$$\int_0^H c c_g a_1 \left(\frac{x}{H}\right) \left(1 - \frac{x}{H}\right)^2 dx = c c_g a_1 \left( \frac{x^2}{2H} - \frac{2x^3}{3H^2} + \frac{x^4}{4H^3} \right) \Big|_0^H =$$



$$= c c_g a_1 \cdot \frac{H}{12}$$

Following this method:

$$\int_{\Omega} \phi_i d\Omega = \frac{H}{2}$$

$$\int_{\Omega} \phi_i \phi_j d\Omega = \frac{H}{3} \quad (i=j)$$

$$= \frac{H}{6} \quad (i \neq j)$$

Per element the matrix description of  $\varphi_{1i}$  and  $\varphi_{2i}$  may be expanded.

$i = 1, 2, 3$  ( field element ) produces a 6 x 6 matrix,

$i = 1, 2$  ( boundary element ) produces a 4 x 4 matrix.

( The right hand part of the equation for the boundary element will be a 1 x 4 matrix: " the element vector " . )

evaluating the integrals in the manner described above while expanding the matrix over the number of nodes per element the equation will take the form indicated below:

The general form of the matrix of the boundary element ( compare with the second term in equations ( 3.05.1 ) and ( 3.05.2 ) . ) :

$$\frac{H}{12} \begin{bmatrix} 3d_{11}+d_{12} & -(3d_{21}+d_{22}) & d_{11}+d_{12} & -(d_{21}+d_{22}) \\ 3d_{21}+d_{22} & 3d_{11}+d_{12} & d_{21}+d_{22} & d_{11}+d_{12} \\ d_{11}+d_{12} & -(d_{21}+d_{22}) & 3d_{11}+d_{12} & -(3d_{21}+d_{22}) \\ d_{21}+d_{22} & d_{11}+d_{12} & 3d_{21}+d_{22} & 3d_{11}+d_{12} \end{bmatrix}$$



The general form of the vector of the boundary element  
( compare with the third term in equations ( 3.05.1 )  
and ( 3.05.2 ). ) :

$$\frac{H}{6} \begin{pmatrix} 2e_{11} + e_{21} \\ 2e_{21} + e_{22} \\ e_{11} + 2e_{12} \\ e_{21} + 2e_{22} \end{pmatrix}$$

where: H : the length of the interval

$d_{ij}$  : the coefficient of  $\phi_j$  in node "j".

$e_{ij}$  : the coefficient of  $\phi_i$  in node "j".

$d_{ij} = a_{ij}$  for the Side View Model

$= a_{ij} \cdot cc_g$  for the Top View Model

$e_{ij} = g_{ij}$  for the Side View Model

$= g_{ij} \cdot cc_g$  for the Top View Model

( see equations ( 3.05.1 ) and ( 3.05.2 ). )

For the matrices of the field element the description  
is obtained in a similar way.

### 3.06 Computing element matrices and vectors

The subroutines ( of the computer program ) which com-  
pute the contributions of each element to the system  
of equations ( following the method described in sec-  
tion 3.05 ) are called:

#### Side View Model:

ELM148 field element matrix

ELM149 boundary element matrix

ELR149 boundary element vector



Top View Model:

ELM301 field element matrix  
ELM302 boundary element matrix  
ELR302 boundary element vector

That is:

ELM148 describes the Laplace equation

ELM301 describes the refraction - diffraction equation

The boundary condition:

$$\frac{\partial \varphi}{\partial n} = -\partial \varphi + q$$

is described by:

ELM 149 and ELR149 ( for the Side View Model )

ELM 302 and ELR302 ( for the Top View Model )

A listing of these subroutines is presented in Appendix 7.

Note that the boundary elements all require input:  $\hat{a}$ ,  $\hat{g}$ ,  
and all the elements of the Top View Model require  
the input:  $cc_g$  ( for the field elements:  $kcc_g$  ).

These data have to be provided in the program for computation.

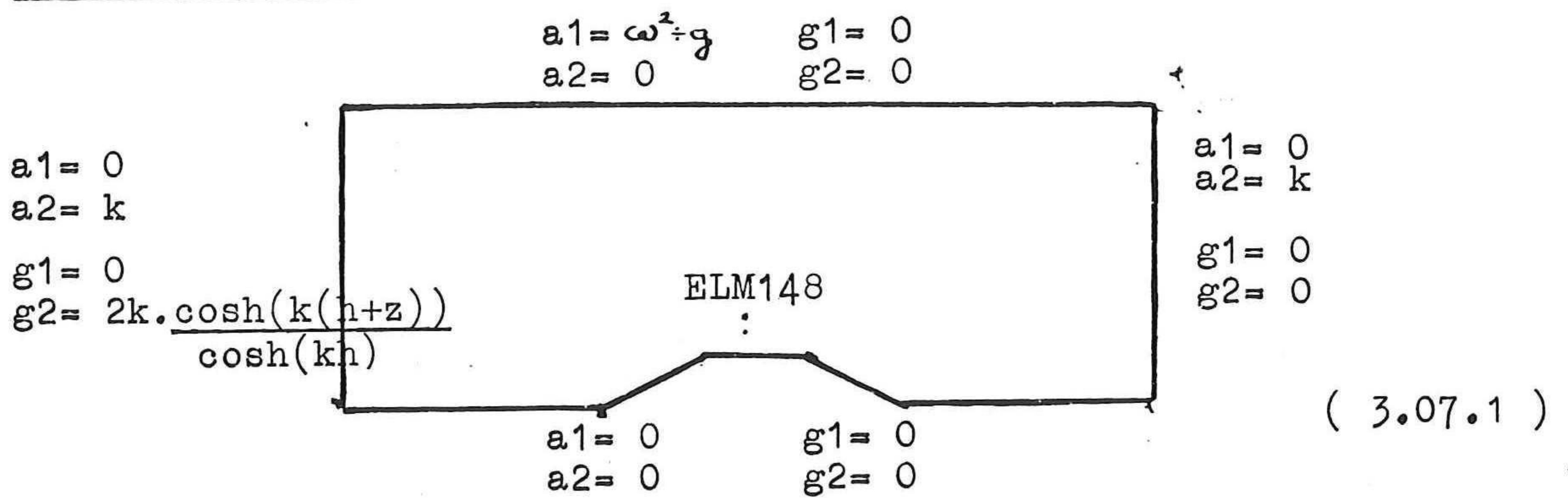


3.07

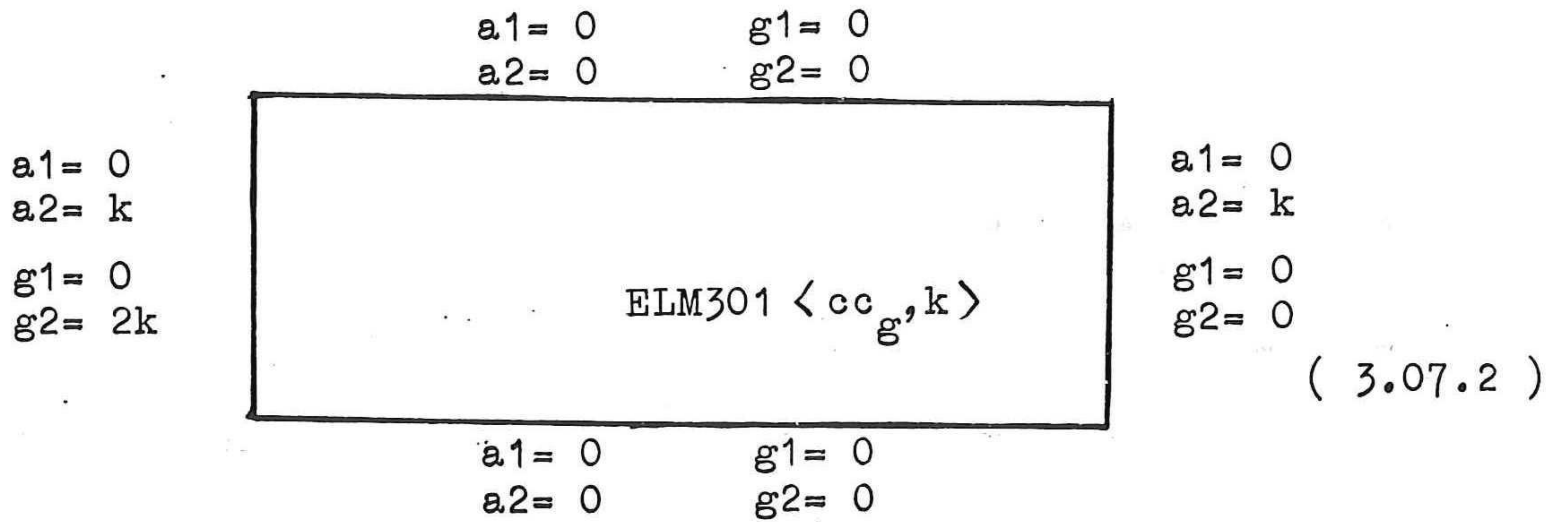
Summary of the numerical description

The numerical form of the two models which will be used for the comparative calculation is:

Side View Model :



Top View Model :



For the description of the boundary elements,  
the Side View Model uses: ELM149 and ELR149  
the Top View Model uses: ELM302 and ELR302.  
The required input data are indicated above.









4.01 Basic structure of the program

The program contains the complete description of all the actions to be performed by the computer for the achieving of the desired goal: obtaining predictions of the potentials of the wave field for the two models: the Side View Model and the Top View Model. (These predictions are then compared. )

For both models the line of action is similar:

- 1) - Preparing the computer for the job
- 2) - Generate elements for the finite element method
- 3) - Prescribe the equations for the approximate solution per element ( ELM148,ELM149,... )
- 4) - Provide the required input for " 3) "
- 5) - Compute the contribution from each element to the system of equations.
- 6) - Assemble the large matrix "  $\bar{A}$  " and the vector "  $\bar{g}$  " from the local contributions.
- 7) - Solve "  $\bar{\phi}$  " from the system of equations:  $\bar{A}\bar{\phi} = \bar{g}$  .
- 8) - Present the solution.
- 9) - Signal the computer that the job is done.

The steps 1 through 9 are described in the following paragraphs. The program can be found in Appendix 4.



4.01.1 Preparing the computer for the job

A number of things have to be taken care of before the job can be executed by the computer. ( the " job " is the set of statements which command the actions to be performed for the computation. )

The administrative affairs of how the computer is to treat and execute the program is described in the Job Control Language ( "JCL" ) section of the program:

- Assign a memory space for the job and a maximal time interval to be used.
- Indicate the account where the costs of the job are to be charged.
- Indicate which computer language and compiler are to be used.
- Specify all input and output datasets.
- Specify all libraries to be used ( AFEP, ... ).
- Provide the program of the job.
- Provide datacards as required by the job.
- Mark the end of the program.

AFEP forms an essential part of the structure of the program. Information concerning all AFEP subroutines is to be found in the AFEP User Manual.



The job program consists of a main program and a set of subroutines which are called by the main program.

In case the subroutines are not already available in one of the libraries linked to the program they have to be added after the main program.

Fortran IV is used as the computer language for the program of this study.

This requires that in every main program the following information is presented:

- the amount of memory space to be reserved for every array which is used.
- the code which is used for variables which are to be treated as Real numbers.
- the variables which are to be used in "common blocks".

For more detailed information on Fortran IV consult any of the available books, e.g. Chirlian (1973).

#### 4.01.2 Generate elements for the finite element method

The field is divided into subregions called "elements" using an AFEP subroutine: "MESB2D". ( See the AFEP Manual for details. )

This meshgenerator creates a mesh within vertices on the boundary of which the user provides the coordinates.

The user also has to specify how many nodal points are



to be made on each section of the boundary and how the spacing is to be done. ( Segal, 1975 )

The " meshgenerator " then uses this information to generate the mesh in the field.

Subroutine " INTAKE " provides the information the mesh-generator requires. ( See Appendix 4. )

The form of the field ( laid down in INTAKE ) is varied in this study. The variation of INTAKE is done by the subroutine " VARY ".

#### 4.01.3 Prescription of the equations

The equations to be used for the computations, for both field and boundaries, are indicated by their number in a datacard input section at the end of the program following AFEP convention.

( See ( 3.07.1 ), ( 3.07.2 ) and the AFEP User Manual ).

#### 4.01.4 Providing the input for the equations

As is described on page - 44 - the equations ( 3.07.1 ) and ( 3.07.2 ) require input.

This information is made available using arrays " USER " and " IUSER " which, following AFEP convention provide the input for the subroutines describing the equations ( ELM148, ELM149, ELR149; ELM 301, ELM302, ELR302 ).

( See AFEP User Manual, Appendix 4 and Appendix 5 ).



Before the values of the required parameters of the input (  $a, c, c_g, k$  ) can be stored they have to be calculated.

This is done in subroutine " DEEP " ( Side View Model ) resp. subroutine " LOCAL " ( Top View Model ).

Subroutine " AUSER " stores the calculated values in the arrays " USER " and " IUSER ". ( Appendix 5 ).

#### 4.01.5 Computing matrices and vectors per element

AFEP subroutine " COMMAT " does the job.

#### 4.01.6 Assembling matrix " $\hat{A}$ " and vector " $\hat{g}$ "

Done by AFEP subroutines " MATRIX " resp. " RHSIDE ".

#### 4.01.7 Solving $\hat{\varphi}$ from $\hat{A} \hat{\varphi} = \hat{g}$

AFEP subroutine " SOLVE " does the job. SOLVE is based on a direct method. Matrix  $\hat{A}$  is subjected to LU decomposition and used in the call of SOLVE in that form. SOLVE consists of 10 subroutines. A description can be found in the AFEP Programme's Guide.



4.01.8 Presenting the solution

The solution can be presented in several ways:

Subroutine PRINRV prints the values of the potential  $(\varphi_1, \varphi_2)$  at each nodal point.

PLOTCH makes a contourplot using the results of PRINRV. Input is required for PLOTCH; this is provided by subroutine INPLOT.

GRAPH prints a graph of the values along a specified boundary, drawing on the results of PRINRV.

In this study the last form of presentation has been used ( using the surface boundary of the Side View Model and the corresponding side boundary of the Top View Model ).

4.01.9 Signaling the end of the program

This is done by the statement " end " at the end of the program.

—||—



4.02            Using the computer program

The computer program has been constructed in such a way that equations ( 3.07.1 ) and ( 3.07.2 ) are each described completely in a program.

A link has been created between these two programs such that the parameters of the Top View Model (  $cc_g$  and  $k$  ) are calculated using the values of the depth of the configuration of the Side View Model.

The position of the nodal points is ensured to be the same for both models ( along the axis where the values of the predictions will be compared ) by using the coordinates of the Side View Model ( which have been generated by the mesh generator ) for the Top View Model as well.

In this manner it is certain that the two models both compute the potentials for identical configurations.

A practical consequence for the use of the program is that the computation of the Side View Model always has to be done prior to the computation of the Top View Model for the same configuration.



The computer program, as it has been designed, readily permits variation of:

( See Appendix 4. )

- the bottom configuration ( " ITEST ", " RATIO " )
- the waterdepth ( " SCD " )
- the wavefrequency ( " OMEGA " )

The next chapter describes how the program is tested.

Chapter 6 describes how the program is used.





5.00 Introduction

The computer program which has been described in chapter 4 ( the program can be found in Appendix 4 ) has to be tested. This is done by using the program to predict the values of the potential  $\phi$  for a configuration for which there exists another method of determining the values of the potential.

This chapter shows which configuration has been used for this purpose and how the predictions are verified.

5.01 Configuration for the verification

A flat and horizontal bottom has been selected as the configuration for the verification.

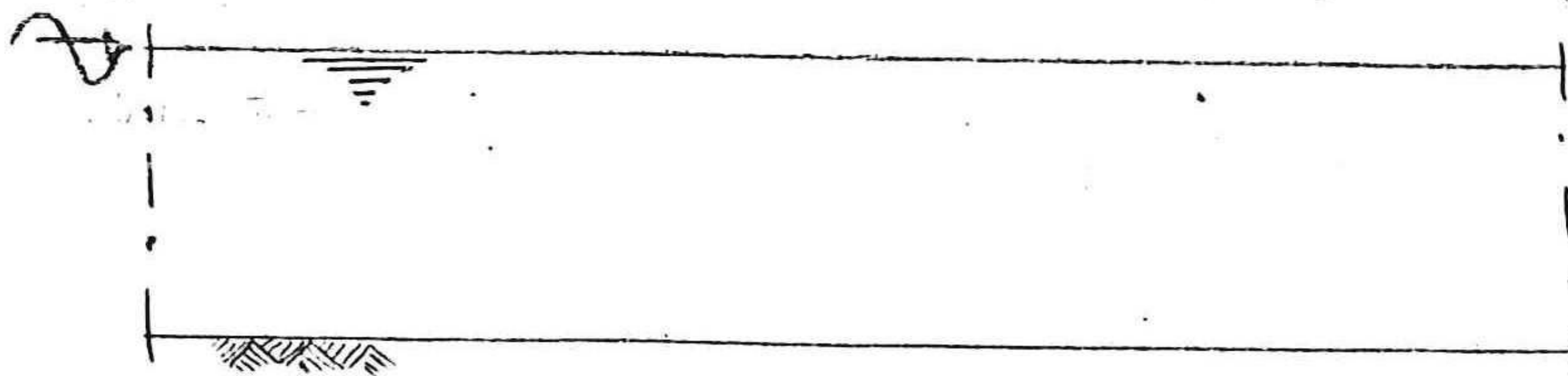


Figure 5.01.1 Sketch of the selected bottom configuration for the testing of the program.

The wave potential can be determined analytically for this case .



## 5.02 Verifying the predictions

There exist two ways for verifying the computed values of the wave potential  $\phi$  :

- 1) The wavelength must be correctly predicted,
- 2) The magnitude of the potential,  $\|\phi\|$ , must be invariant.

( A description of the relationship between the wave potential and its magnitude is provided in section 6.05 )

### 5.02.1 Wavelength

The wavelength has been determined for two values of the waterdepth. These two values are the same as have been used in the computations for the configurations described in the next chapter.

To determine the wavelength for the flat and horizontal bottom case using analytical methods is a simple matter, described in any book on wave mechanics (e.g. Lamb, 1975 ).

To determine the wavelength for the same case from the computer results is also rather straightforward but does require some further description of the interpretation of the results as presented by the computer ( this is done in section 7.02 ) and the way the accuracy of the thus obtained values ( which will be done next ).



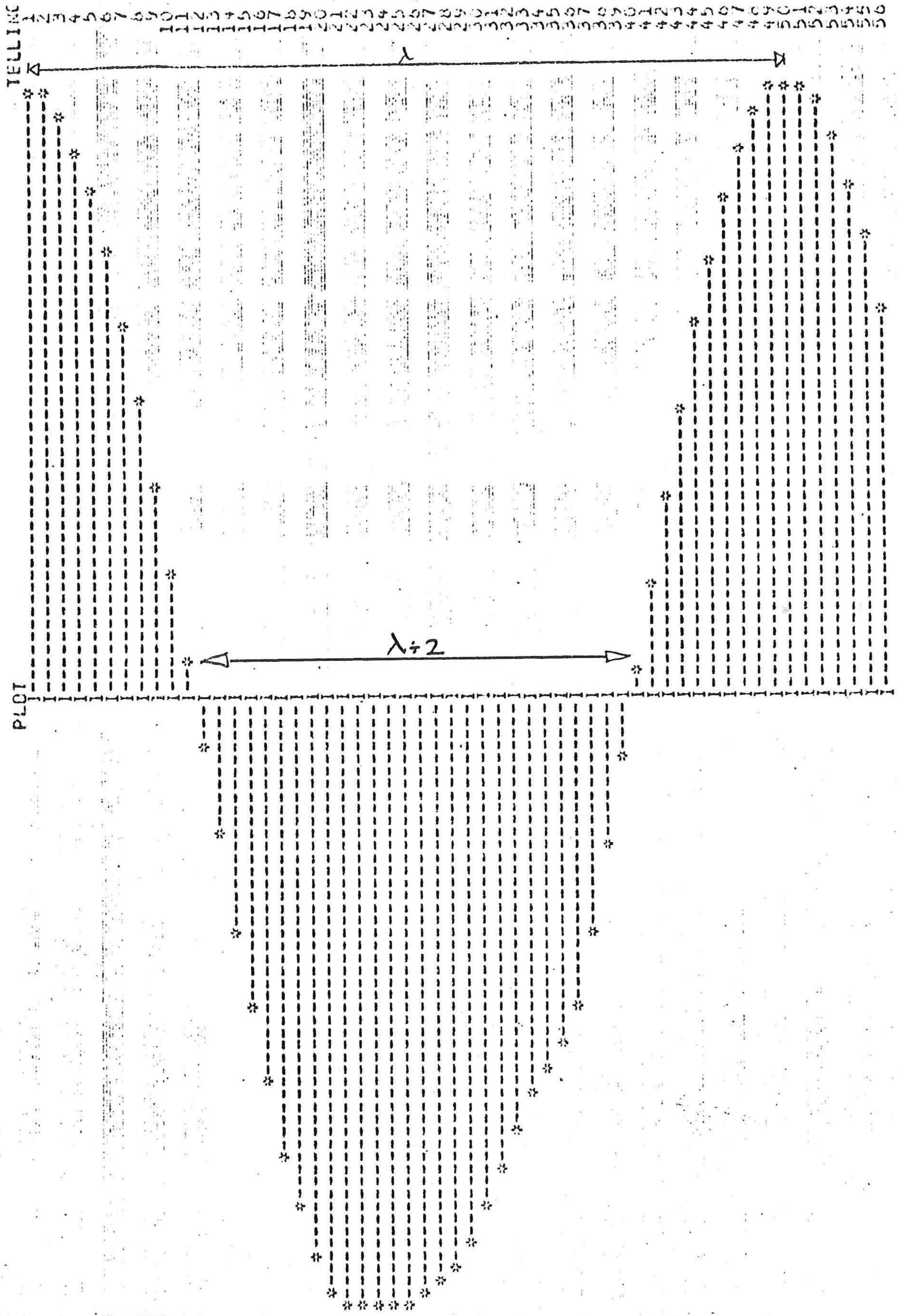


Figure 5.02.1.1

Example of the graph of the computed value of  $\phi_1$  from which the wavelength may be determined.



Figure ( 5.02.1.1 ) shows the computed values of the IReal part of the wave potential  $\varphi$ .

The figure clearly shows the wave pattern.

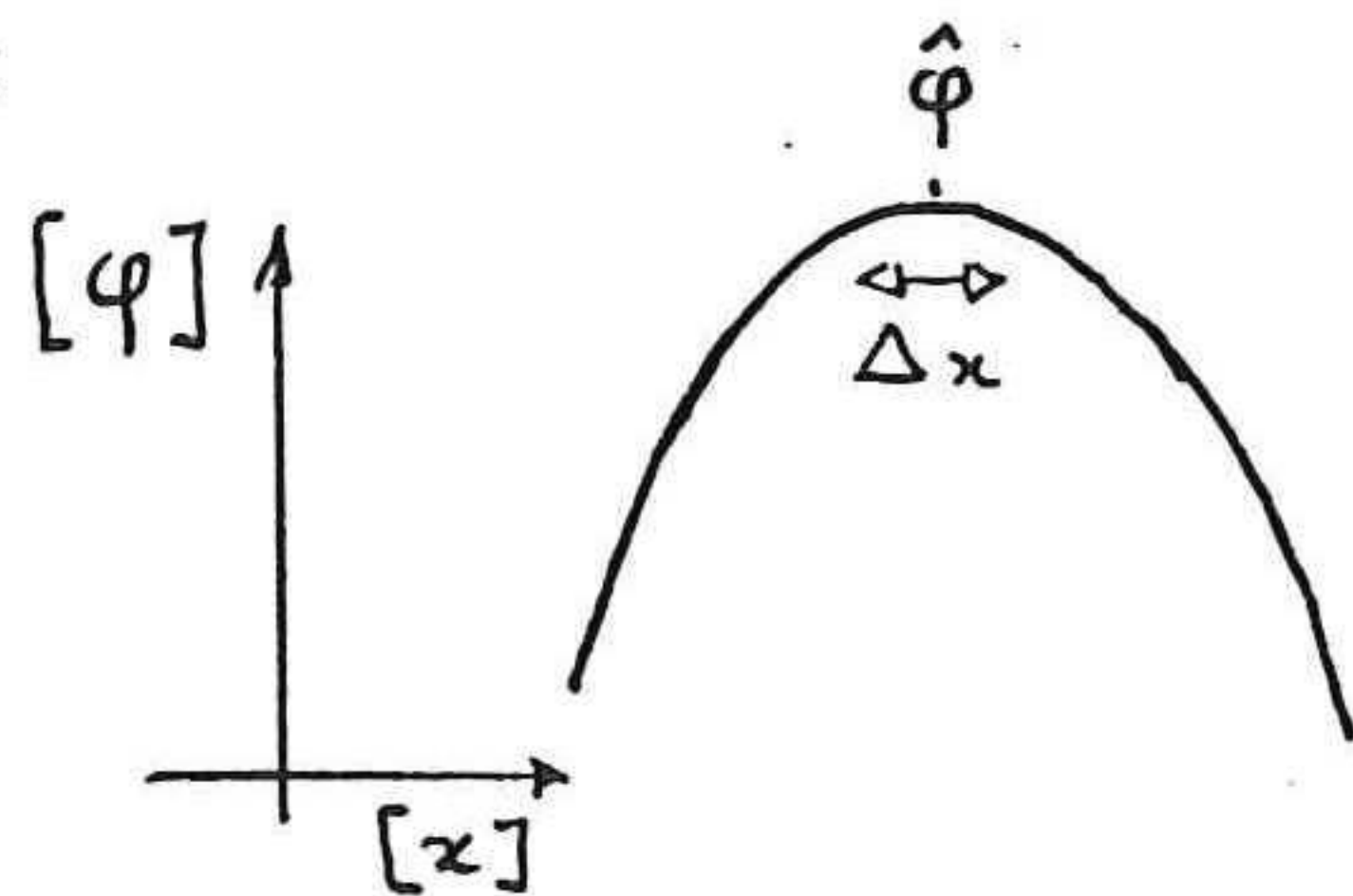
Observe that the depicted wave pattern is not exactly sinusoidal. This is caused by the fact that the steplength is not the same throughout the field. ( The stepsize has been made smaller in the area where the greatest change in the value of the potential would be expected: over the dam. See figures ( 7.01.1 ) and ( 7.01.2 ).)

( A discussion on the numerical aspects of the chosen step-length can be found in Appendix 6. )

The exact location of each node can always be determined from the input data of the program.

Since the location of a node evidently does not have to coincide with the location of a maximum or minimum, an error will be made in determining the wavelength.

If the wavelength is determined from the location of the maxima ( See Figure ( 5.01.1.1 ), the uncertainty of the value will be as great as the steplength, as follows from the following sketch:



$$x \langle \hat{\varphi} \rangle = x \langle \tilde{\varphi} \rangle \pm \frac{\Delta x}{2}$$

$\hat{\varphi}$  : max of  $\varphi$ , continuous function.

$\tilde{\varphi}$  : max of  $\tilde{\varphi}$ , approximate function.

$\Delta x$  : steplength.

( 5.01.1.2 )

( In the studied cases:  $\Delta x = 0.11$  )

A more accurate value of the wavelength may be obtained by interpolating for the location of the zero-crossings of the sinus function. In that case the uncertainty will be twice



as large as the round-off error of the value of the x-coordinate. ( Which, for the studied cases is:  $\pm 0.01$  round-off .)

The results of the calculated and graphically determined values of the wavelength are compared:

Table 5.01.1.3 Comparison of the graphically determined wavelength versus the calculated wavelength for a horizontal flat bottom. ( SVM= Side View Model; TVM= Top View Model )			
Depth	wavelength SVM	wavelength TVM	calculated wavelength
0.4	( 3.71 $\pm$ .02 )	( 3.71 $\pm$ .02 )	( 3.71 $\pm$ .01 )
0.6	( 4.39 $\pm$ .02 )	( 4.39 $\pm$ .02 )	( 4.38 $\pm$ .01 )

These results show that the results of the computation do indeed correspond with the calculated values of the wavelength.

#### 5.02.2 Magnitude of the potential: $\|\varphi\|$

Discussion of the relationship between the wave potential and its magnitude, as well as their interpretation, will be done in chapter 6.

At this moment it suffices to observe the behaviour of  $\|\varphi\|$  for the case of a flat and horizontal bottom. In this case the value of  $\|\varphi\|$  should be constant.

The values of  $\|\varphi\|$  for both the Side View Model and the Top View Model are displayed in figures ( 5.02.2.2 ) and ( 5.02.2.3 ).

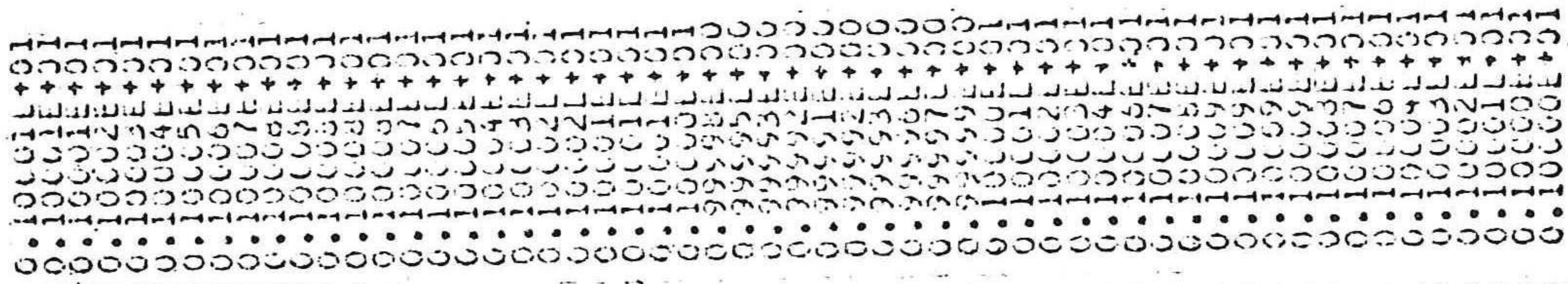
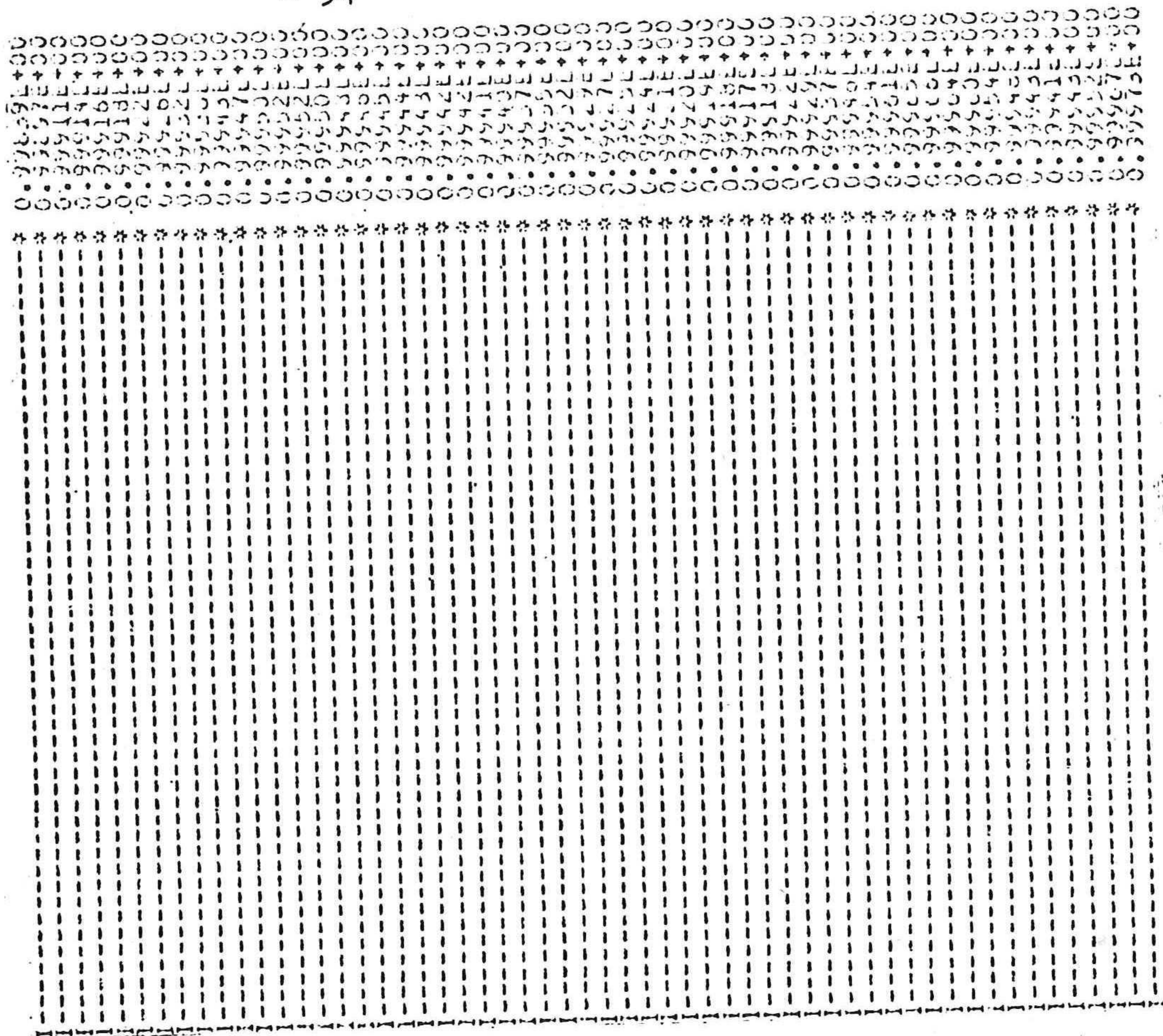
The two cases are the same as which have been used in section ( 5.02.1 ) for the verification of the computed wavelength.







$\|\varphi_t\|_s$



$\|\varphi_t\|_\tau$

Figure 5.02.2.2 Predicted values of  $\|\varphi_t\|_s$  and  $\|\varphi_t\|_\tau$  for the case of a flat horizontal bottom, for CASE 2.



Inspection of the results shows that the predicted values of  $\|\varphi\|$  are indeed constant for these cases:

Table 5.02.2.3 Predicted values of $\ \varphi\ $ for a flat horizontal bottom. ( See figures (5.02.2.1) and (5.02.2.2)		
Depth	Side View Model $\ \varphi\ $	Top View Model $\ \varphi\ $
0.4	( 0.9996 $\pm$ 0.0004 )	( 1.0000 $\pm$ 0.0013 )
0.6	( 0.9994 $\pm$ 0.0003 )	( 1.0000 $\pm$ 0.0009 )

### 5.03 Alternative verification

It has been possible to check the computations by the computer program for the Top View Model by comparing the results versus results obtained for identical cases using a one dimensional model bases on the Refraction-Diffraction equation. This model, made by N. Booy, produced identical results ( for the compared cases ) as the model which has been used in this study.

Appendix 6 shows the results of the model by N. Booy.

~



6.00      Introduction

The preceding chapters described how the Side View Model and the Top View Model have been derived and how a computer program has been made which, for both models, can compute the potentials of the wavefield for a selected bottom configuration.

This chapter shows how this program has been used for this study and which bottom configurations have been selected for the comparison.

Beginning with a summary of the parametric description for the general case of wave propagation and the specific case which is selected for this study, the chapter then continues with a description of the characteristic parameter which has been selected for the comparison. Finally a description of three main types of variations of the selected bottom configuration is presented.



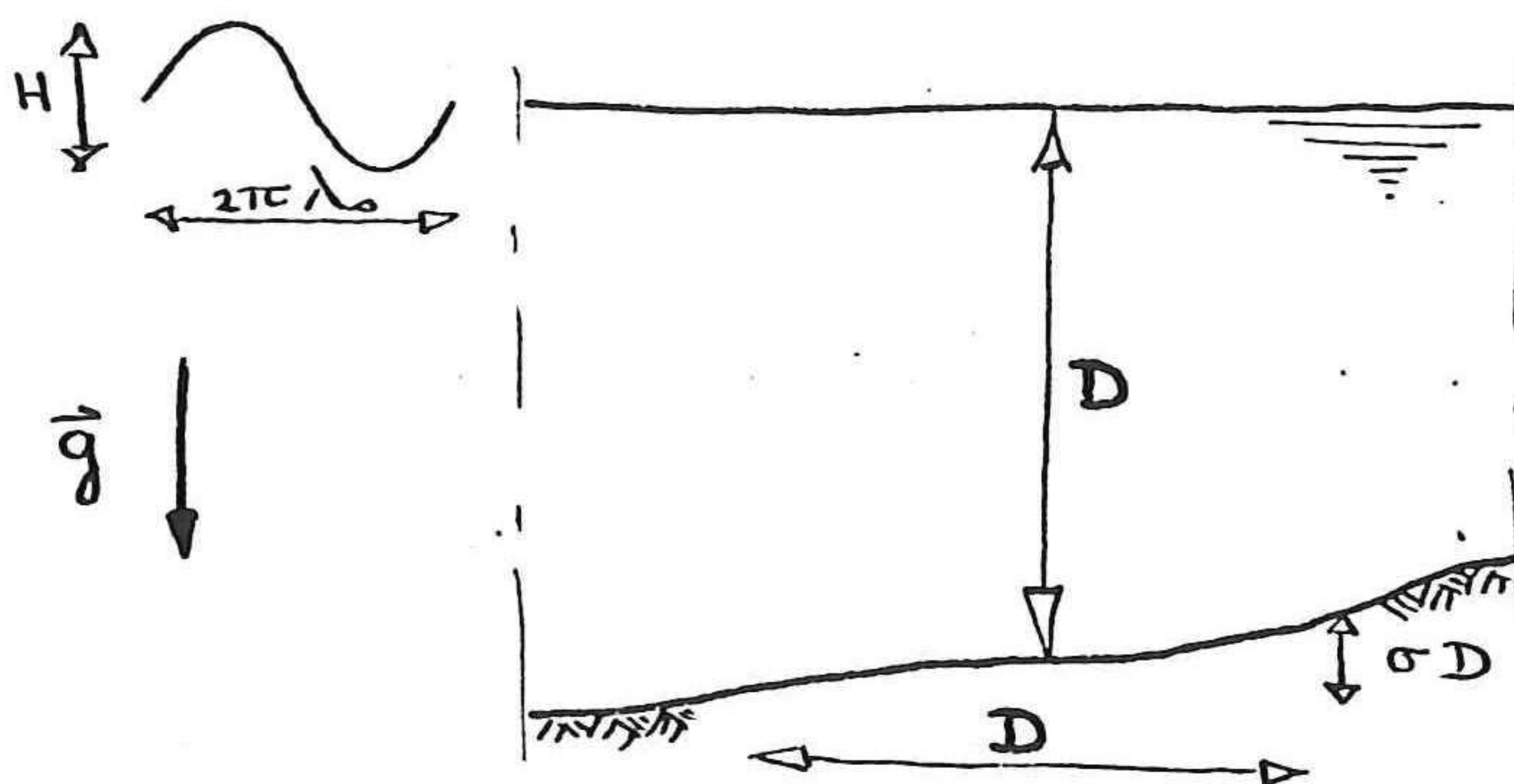
6.01            Goal of the test

The goal of the test is: to determine the difference in the prediction of the Side View Model and the Top View Model as a function of a difference in the bottom configuration. ( c<sub>q</sub> to determine if such a difference does exist. )

6.02            Parameters of the wave field

A number of restrictions and assumptions underlie the mathematical models ( 2.07.1 ) and ( 2.07.2 ).

Figure 2.00.1, repeated here for convenience, shows the characteristic quantities and parameters which are important in wave propagation problems. (Berkhoff, 1976 ).



( 6.02.1 )

where: D : mean water depth

H : wave height

$$\lambda_0 = g + \omega^2$$

$\sigma$  : mean slope of the bottom over a distance D.



For a linear model it is assumed that the wave steepness

$\gamma = H/\lambda_0$  is small.

*wave height?*

In the case of a small waterdepth the ratio  $\alpha = H/D$  must also be small.

In the refraction - diffraction equation it is assumed that the variation of the bottom is moderately small and noticeable for the surface waves ( $\frac{\sigma\lambda_0}{D} \ll 1$ ).

The curvature of the wave amplitude is taken into account.

A two dimensional equation in the horizontal is obtained by integration over the waterdepth.

( See Appendix 2. )

These assumptions do not have to be made for the Laplace model.

*but this is  
linear  
model*

### 6.03

### Selected bottom configuration

A trapezoidal dam has been selected because of the simplicity of its contour and the clearly defineable bottom slope.

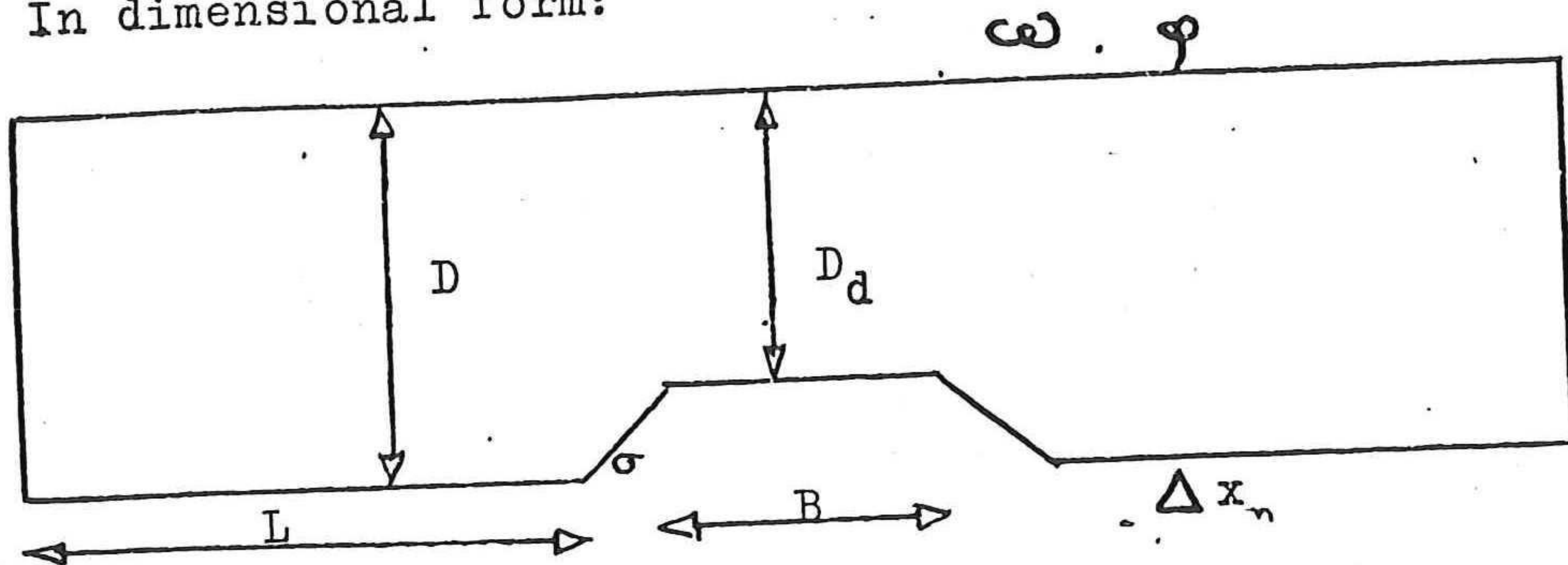
The selected setup is depicted in figure ( 6.03.1 ).

This figure shows the characteristic quantities and parameters which are important for the comparison of the Side View Model versus the Top View Model.

( Compare with ( 6.02.1 ) )



In dimensional form:



( 6.03.1 )

- where:
- $\omega$  : wave time frequency
  - $\lambda_0$  : wave length at deep water
  - H : wave height
  - D : waterdepth
  - $D_d$  : waterdepth over the dam
  - $\sigma$  : slope of the incline of the dam
  - B : width of the crest of the dam
  - L : length of the zone left of the dam
  - $\varphi$  : wave potential at the surface:  $\varphi = \varphi_1 + i \varphi_2$
  - g : gravitational potential
  - $\Delta x_n$  : numerical steplength of the computation

It is more convenient to use dimensionless parameters. The coordinates are made dimensionless using the free surface characteristic length  $\lambda_0 = g + \omega^2$

*(deep water wave length)*

As Berkhoff (1976) has done, the gradient of the waterdepth can be made of the order one by taking new horizontal coordinates (  $x'$ ,  $y'$  ) in the ( dimensionless ) waterdepth function  $\frac{D}{\lambda_0} \langle x, y \rangle$  :

$$\varepsilon \langle x', y' \rangle = \frac{\sigma}{D/\lambda_0} \langle x, y \rangle$$

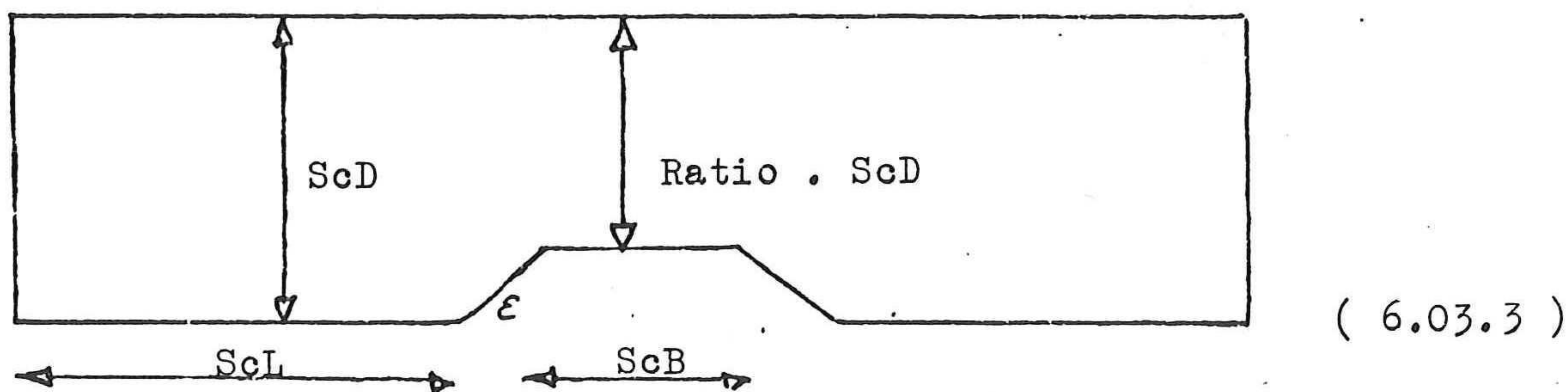
As a consequence of this the bottom slope is characterised by the parameter ( see figure ( 6.03.3 ) ) :

$$\varepsilon = \frac{\sigma \lambda_0}{D}$$

( 6.03.2 )



In dimensionless form:



where:

$$\begin{aligned} ScD &= D + \lambda_0 \\ \text{ratio} &= D_d + D \\ ScB &= B + \lambda_0 \\ ScL &= L + \lambda_0 \\ Sc\Delta x_n &= \Delta x_n + \lambda_0 \\ \epsilon &= \sigma + (D + \lambda_0) \end{aligned}$$

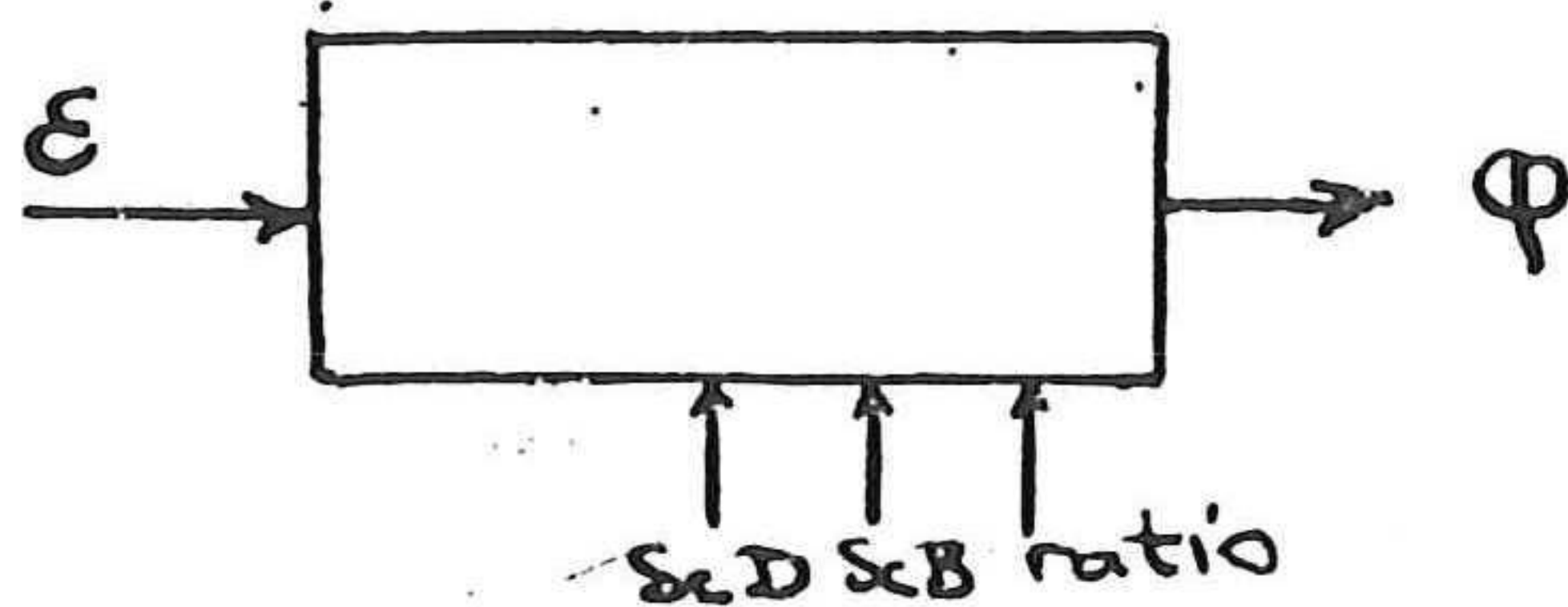
*Under linear model  $\Rightarrow H$  not relevant*

From the derivation of the refraction - diffraction equation ( Berkhoff, 1976 ) it may be expected that this parameter will be important for differences in predictions, by the two models, of the values of the wave potential.

Because of the importance of the assumption ( in the derivation by Berkhoff ) of the bottom variation being moderately small ( see ( 6.03.3 ) and Appendix 3. ), the computer program and the experimentation have been made to focus on the variation of  $\epsilon$ .

For that reason  $\epsilon$  has been varied in a number of cases where, per case, the other parameters remained constant, as depicted in the following system sketch:





( 6.03.3 )

System sketch. The computer program as a system for predicting values of  $\phi$  for varying values of  $\epsilon$  with a chosen setting of values of ratio,  $ScD$ ,  $ScB$ .

#### 6.04 Discussion of the parameters

In this section the parameters  $ScD$ , Ratio and  $ScB$  will be discussed.

Parameter  $\epsilon$  has already been described in ( 6.03.2 ), and  $\phi$  will be described in the next section.

The influence of parameters  $\Delta L$  ( length of the zone upwave of the dam ) and of  $\Delta x_n$  ( steplength ) is described in Appendix 6 . These parameters do not have to be varied in the computational model once they have been suitably determined.

ScD : Dimensionless measure of the depth of the wavefield.

Values of  $ScD$  greater than 1.0 will not be useful:

the effect of the bottom will not be noticeable.

For this reason four values have been selected:

$ScD = 0.6 ; 0.5 ; 0.4 ; 0.25 .$



Ratio : Dimensionless measure of the waterheight over the dam.

The value of Ratio is to be greater than zero, in order to be able to use the mesh generator. ( See section 7.01 ).

In this study no troughs are considered, that is:

Ratio is not greater than 1.0 .

In this study values for Ratio have been selected in the interval ( 0.1  $\rightarrow$  1.0 ).

ScB : Dimensionless measure of the width of the crest of the dam.

The value of ScB may in principle vary between zero and infinite. However, because of the requirement that the influx and efflux boundaries be sufficiently far removed from the barrier, there exists a maximum value of ScB which is determined by, amongst others, the length of the basin.

( The length of the basin is restricted by the maximum number of nodes of the computation, i.e., the costs of the computation. )

This restriction has restricted the values of ScB in this study to a maximum of  $ScB \approx 0.6$  .

By letting the crest of the dam go to the downwave side ( efflux ) boundary, an additional value of  $ScB = \infty$  is obtained.



6.05 Characteristic parameter of the wavefield

The computational model determines the value of the wave potential,  $\varphi$ , at the surface for the various combinations of the parameters  $\xi$ , ScD, Ratio and ScB which determine the bottom profile.

The wave potential,  $\varphi$ , consists of three components:

- $\varphi_i$  : the incident wave
- $\varphi_r$  : the reflected wave
- $\varphi_t$  : the transmitted wave

The comparison of the predicted wave potentials of the Side View Model and the Top View Model will be done using the magnitudes of the potential: ( $\|\varphi\| = (\varphi_1^2 + \varphi_2^2)^{\frac{1}{2}}$ ).

- $\|\varphi_i\|$  : magnitude of the incident wave
- $\|\varphi_r\|$  : magnitude of the reflected wave
- $\|\varphi_t\|$  : magnitude of the transmitted wave

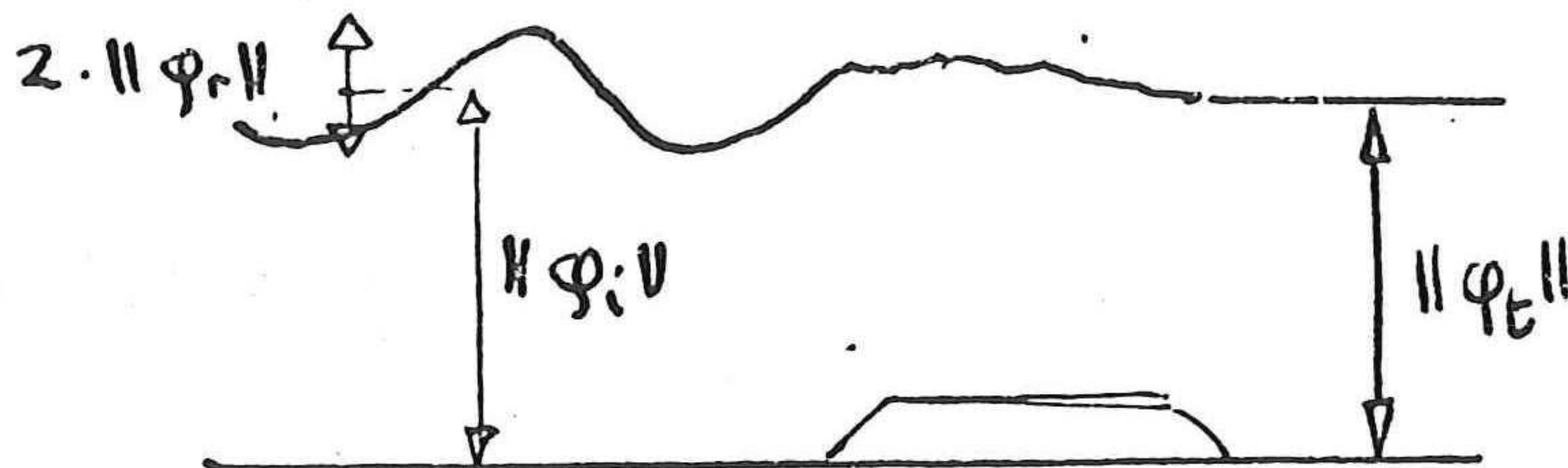
There is a fixed relationship between these three components ( Lamb, 1975 ):

$$\|\varphi_i\|^2 = \|\varphi_r\|^2 + \|\varphi_t\|^2 \quad ( 6.05.1 )$$

With the aid of figure ( 6.05.2 ) it will be shown how, from a graph depicting the values of  $\|\varphi\|$ , the three



characteristic components may be determined.



( 6.05.2 )

$$\text{where: } \begin{aligned} \|\varphi_i\| &= \frac{1}{2}(\|\varphi\|_{\max} + \|\varphi\|_{\min}) \\ \|\varphi_r\| &= \frac{1}{2}(\|\varphi\|_{\max} - \|\varphi\|_{\min}) \end{aligned}$$

In this study the value of the imposed  $\|\varphi_i\|$  is equal to one.

The oscillations of  $\|\varphi\|$  left of the barrier are caused by the interference of  $\varphi_i$  and  $\varphi_r$ ; the amplitude of the oscillations is the amplitude of the reflected component.

The average value of the oscillation is the value of the potential of the incoming wave.

The height of the potential downwave of the barrier is the magnitude of the potential of the transmitted wave.

Due to the chosen magnitude of the imposed wave,  $\|\varphi_i\| = 1$ , and because of the existence of the relationship between  $\varphi_i$ ,  $\varphi_r$  and  $\varphi_t$  as described by equation ( 6.05.1 ), the characterisation of the wavefield may be done using either  $\|\varphi_r\|$  or  $\|\varphi_t\|$  as the characteristic parameter of the wave field.

Because of the greater ease of determining  $\|\varphi_t\|$  ( see figure 6.05.2 ),  $\|\varphi_t\|$  has been selected as the characteristic parameter of the wavefield in the comparison of the two models.



Thus: the comparison of the Side View Model versus the Top View Model will be done by comparing the predicted values of  $\|\varphi_t\|$  of the two models for various bottom configurations.

The predictions will be compared in relative form:

$$\frac{\|\varphi_t\|_{T.V.M.} - \|\varphi_t\|_{S.V.M.}}{\|\varphi_t\|_{S.V.M.}} = \frac{\|\varphi_t\|_T}{\|\varphi_t\|_S} - 1 \quad (6.05.3)$$

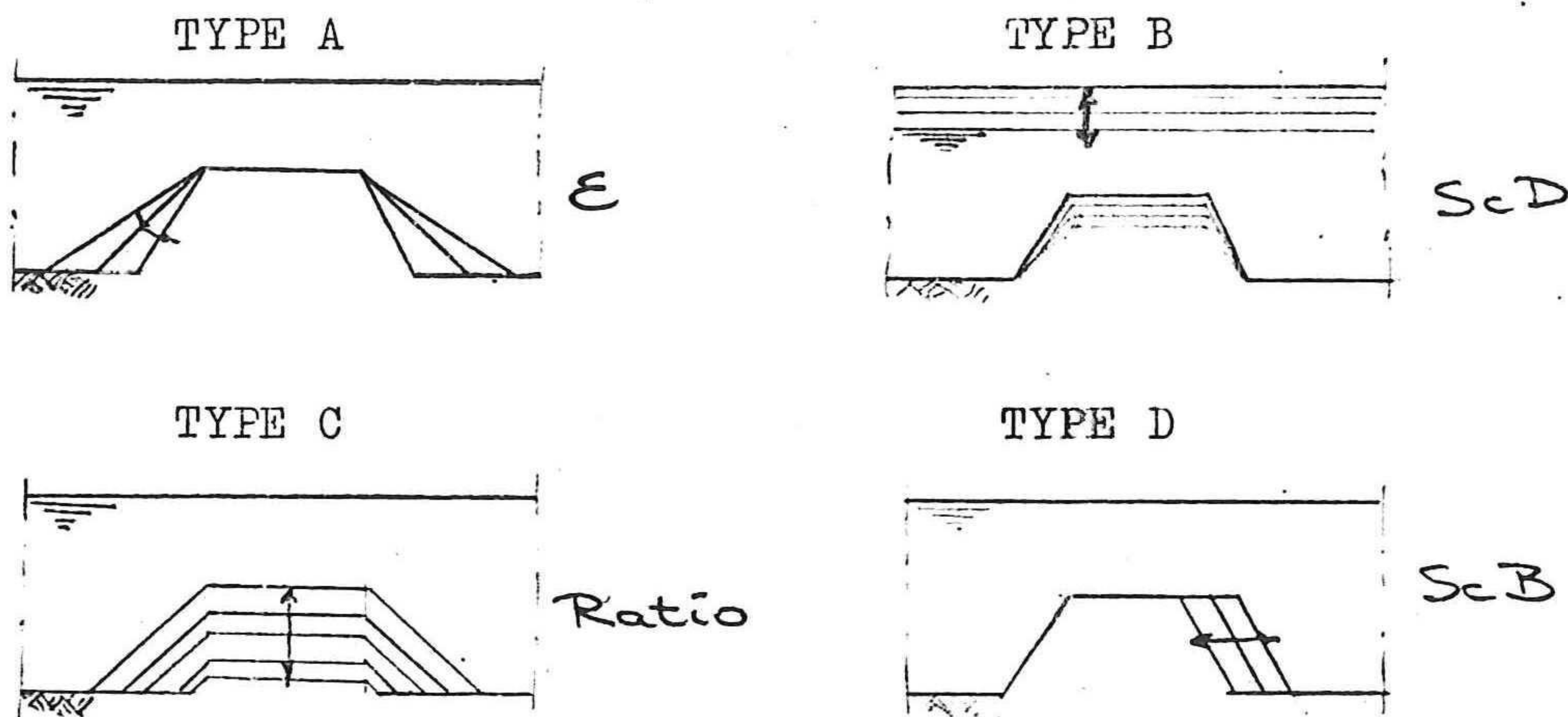
: The results of the Top View Model ( refraction - diffraction equation ) are compared versus those of the Side View Model ('laplace equation') which is taken as reference.



6.06

Three main types of variations

From the descriptions of sections ( 6.03 ) and ( 6.04 ) it follows that there are four basic types of variations, associated with the parameters  $\epsilon$ ;  $ScD$ , Ratio and  $ScB$  :



For the computations these four basic types have been combined into one type which allows for the variations of the four parameters: ITYPE 2 .

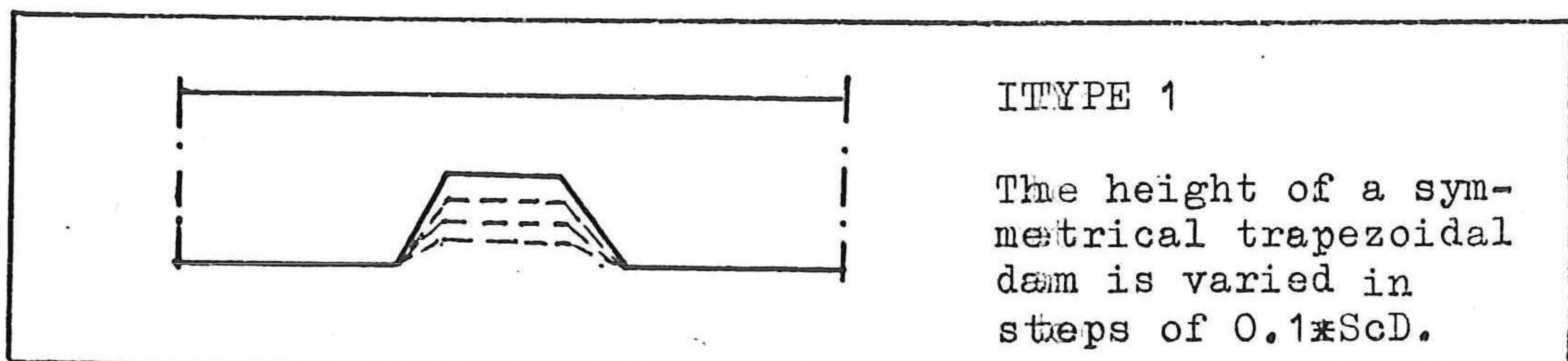
*allien type A*

In order to allow for the case of  $ScB = \infty$  ( as described in section ( 6.04 ) ) ITYPE 3 has been designed.

ITYPE 1 has been designed to provide direct insight into the effects of the variation of the dam height.

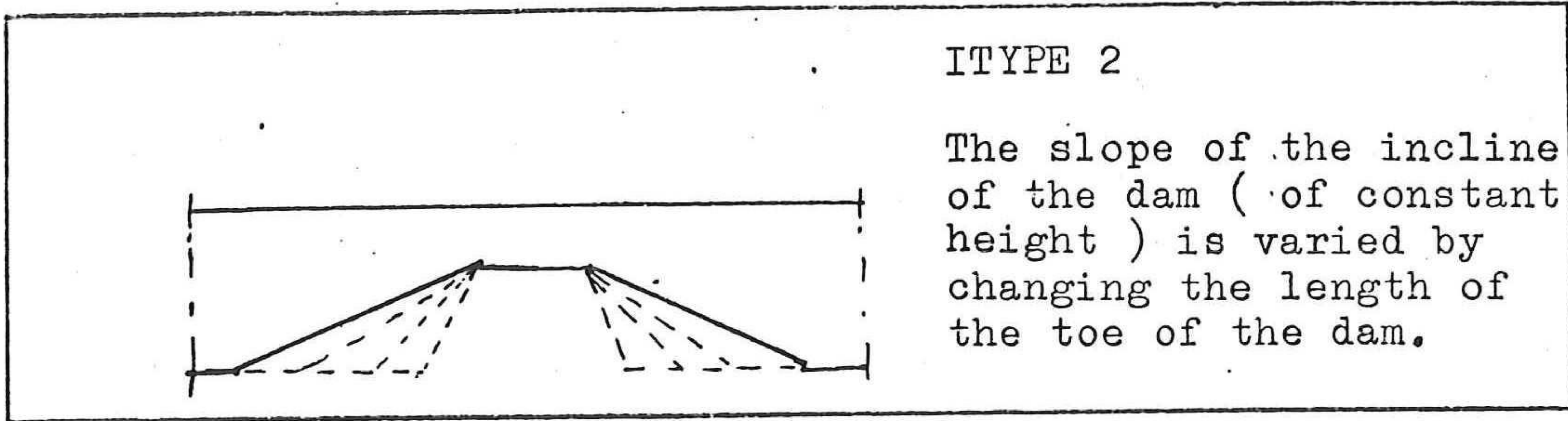
In each ITYPE the value of  $\epsilon$  is varied for a chosen setting of the values of  $ScD$ , Ratio,  $ScB$ .

The three ITYPEs are:

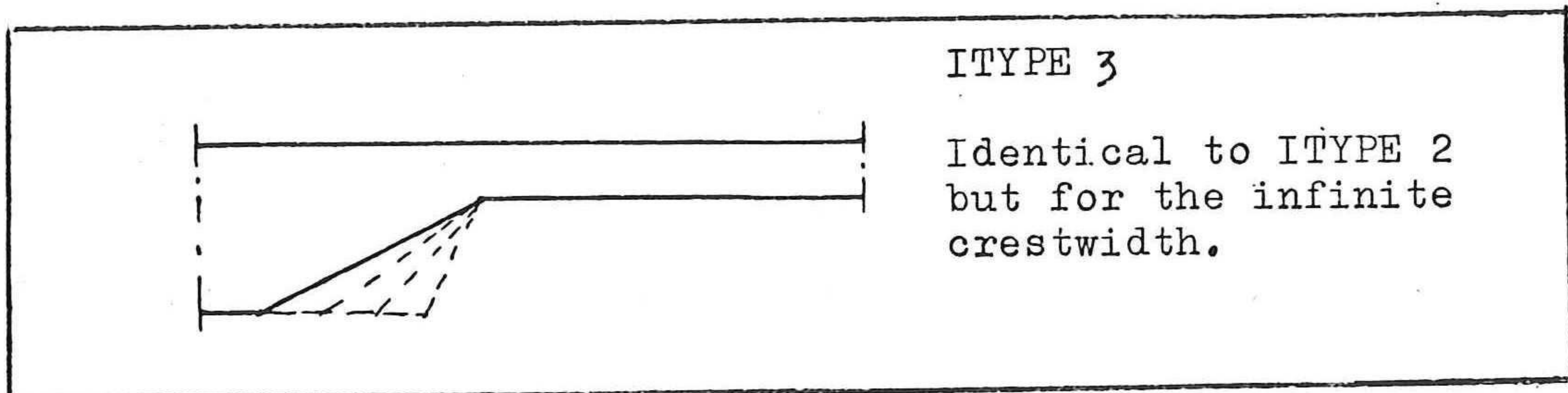


( 6.06.1





( 6.06.2 )



( 6.06.3 )

6.07

Selected parameter combinations

Diagram ( 6.07.1 ) presents the values of the parameters which have been used for the computations, for the three different ITYPES.

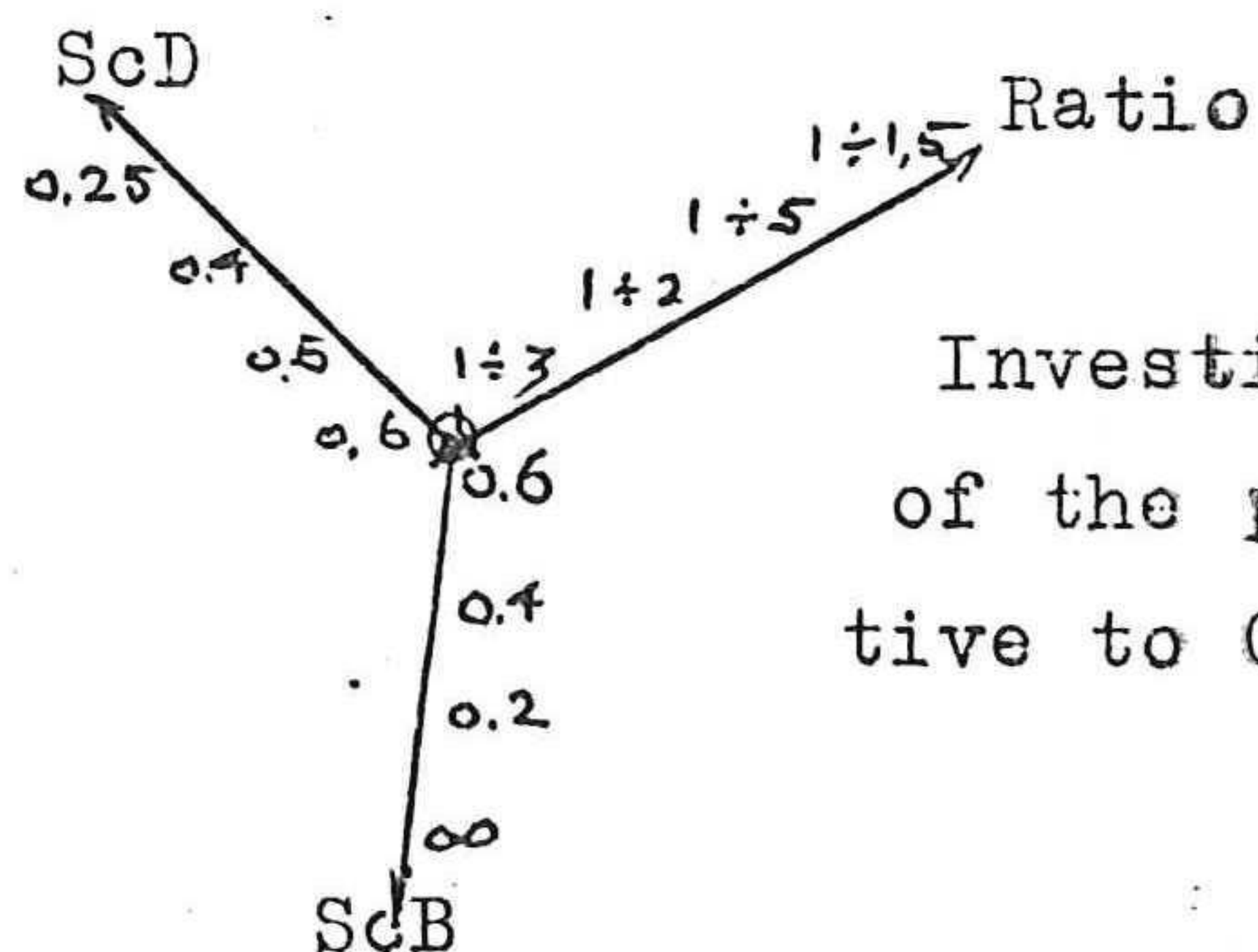
( The description of the ITYPES can be found in Appendix 4, and section 6.06. )

The results of the computations for these three ITYPES are presented in the next chapter.



Diagram 6.07.1 Tested configurations				
ITYPE 1	CASE nr.	ScD	Ratio	ScB
	1	0.4	varies	0.6
	2	0.6	varies	0.6
ITYPE 2	3	0.6	1+3	0.6
	4	0.5	1+3	0.6
	5	0.4	1+3	0.6
	6	0.25	1+3	0.6
	7	0.6	1+5	0.6
	8	0.6	1+1.5	0.6
	9	0.6	1+2	0.6
	10	0.6	1+3	0.4
	11	0.6	1+3	0.2
ITYPE 3	12	0.6	1+3	$\infty$
two cases are also computed for verification				
extended upwave length " L "				
	A	0.6	1+3	0.6
double stepsize				
	B	0.6	1+3	0.6

It can be seen from the above diagram that with respect to CASE 3 the influence of parameter " ScD " is investigated in cases 4 through 6; the influence of " Ratio " in cases 7 through 9 and the influence of " ScB " in cases 10 through 12:



Investigated variations of the parameters relative to CASE 3.



The variations which are already implied in the ITYPEs are the following: ( see section 6.06 ) :

ITYPE 1 varies the height of the dam in steps of a tenth of the waterdepth.

The only required extra information is the length of the toe of the dam: " toe ".

( In this study Toe is kept constant: Toe = 0.2 in ITYPE 1. )

ITYPE 2 varies the length of the toe of the dam for a fixed damheight. ( That is: the slope  $\sigma$  is made to vary by changing the run, keeping the rise constant. )

ITYPE 3 is identical to ITYPE 2 except for the crestwidth. ( See page - 67 - . )

Two cases are computed for the purpose of verification;

CASE A

This case is identical to CASE 3 except for the extension of the length "SGL" upwave of the dam ( see figure ( 6.03.3 ) ).

This case acts as a test for the location of the influx boundary, which should be sufficiently far removed from the barrier.

CASE B

In order to test the stepsize the most critical case of CASE 3 is also calculated with double stepsize.



( Doubling stepsize means less nodes, thus a cheaper computation. If the results are acceptable this does provide sufficient information. See the discussion on this topic in Appendix 6. )

←→



7.00      Introduction

In this chapter the results are presented of the comparative computation.

( In section ( 6.07 ) it has been shown which cases have been selected for the three IYPEs. )

Typical examples are given of a mesh generated by the mesh generator and of the results as presented by the computer program.

A description is given of the manner in which the computer presented data are interpreted to provide the data required for the comparison.

The predicted values of the wave potentials are presented.

←



7.01      The generated Mesh

One of the results which the computer program produces is the mesh.

This mesh plays an essential role in the computation because by the mesh the location of the nodes where the computation is done is determined.

Thus the mesh also shows the steplength with which the computation is done.

The mesh is generated by AFEP's meshgenerator MESB2D, based on user provided characteristic nodes ( vertices of the field, etc. ).

In this study the mesh generator is made to generate a mesh which has a refinement                      in the region where it is expected that the potential changes most: over the barrier.

Figures 7.01.1 and 7.01.2 show, at different scales, typical examples of generated meshes for the Top View Model resp. the Side View Model.

It is possible to let the computer display the details of the generated mesh such as the sections of which the total mesh is composed and the nodal numbering.

An example of the sections of the mesh is shown in figure 7.01.3 .



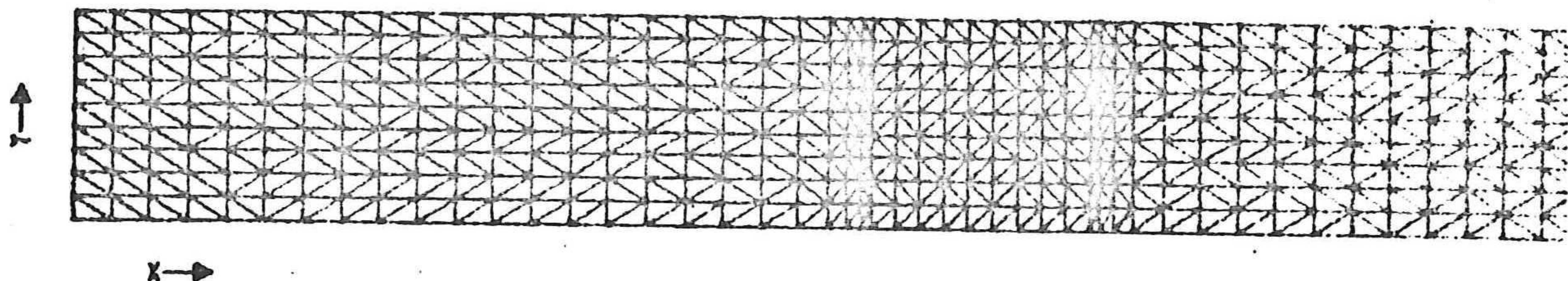


Figure 7.01.1

Example of a mesh generated in the Top View Model.

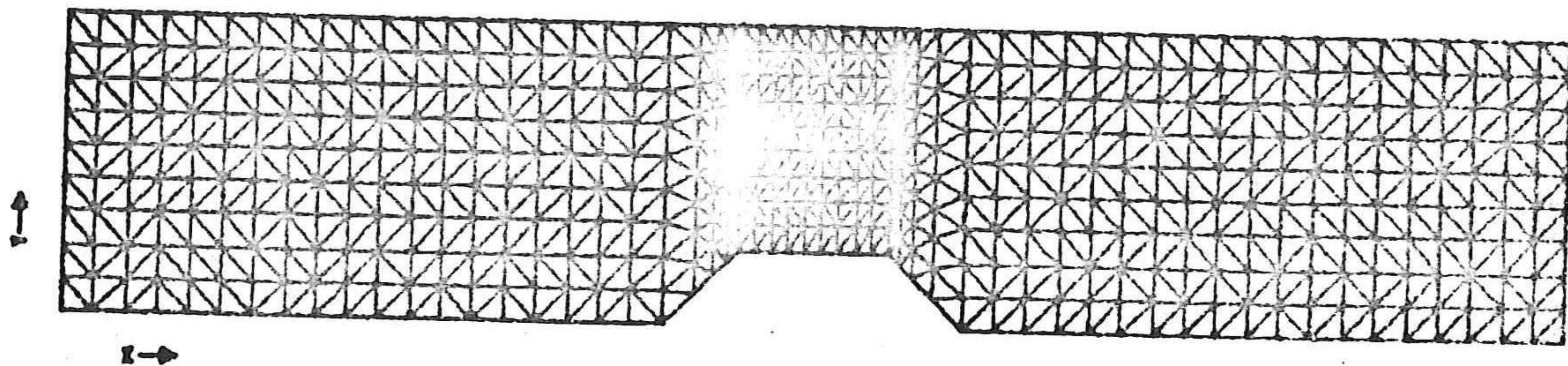


Figure 7.01.2

Example of a mesh generated in the Side View Model.

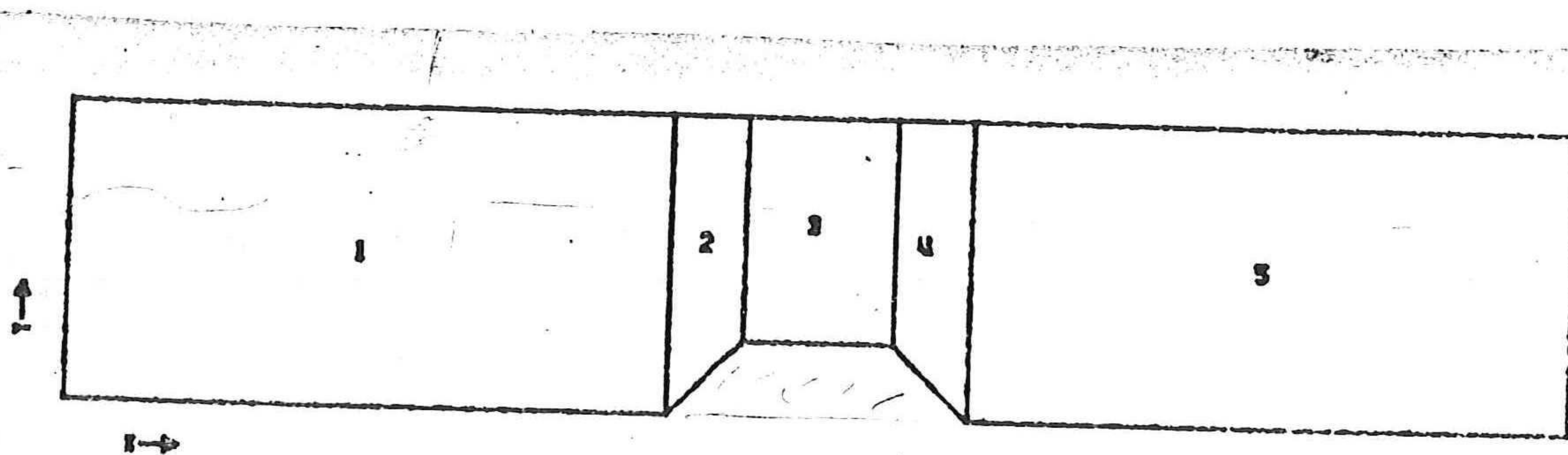


Figure 7.01.3

Example of the division of the field into subfields as done by the mesh generator of the Side View Model.



7.02

Form of presentation of the results

AFEP permits a number of forms of presentation, such as the drawing of contourlines of the potential ( see for an example figure ( 7.02.2 ) ) or a graph of the values of the potential along a selected interval. Both are made on the plotter.

In this study an alternative method has been used: the values of the potential along the surface boundary have been printed ( conform the remarks made in section 2.07 ) and presented in the form of a graph of  $\|\varphi\|$ . In this way the results became more readily available and at smaller costs ( as compared to the rather expensive plotter ) .

An example of the presentation of the results can be found in figure ( 7.02.3 ).

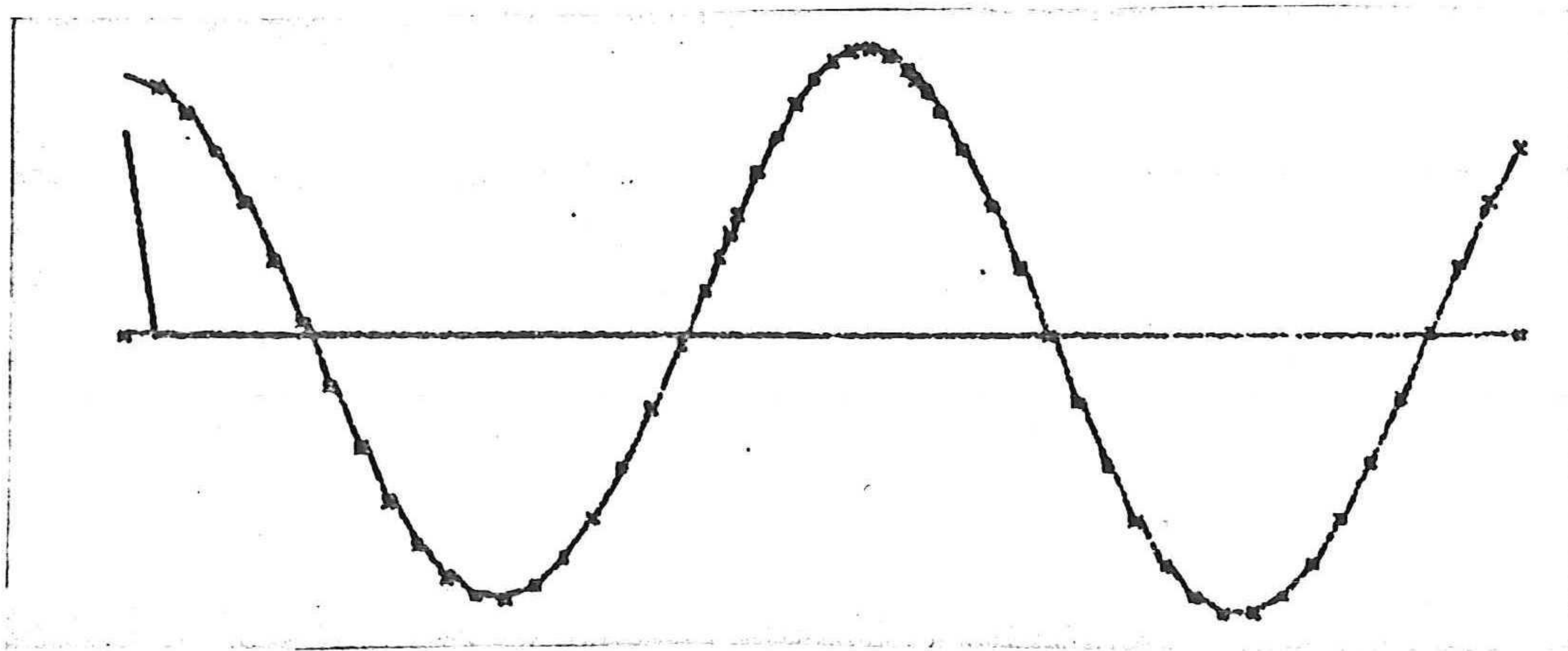


Figure 7.02.1

Example of the plotted values of the potential at the water surface ( in this case: for a flat horizontal bottom. )



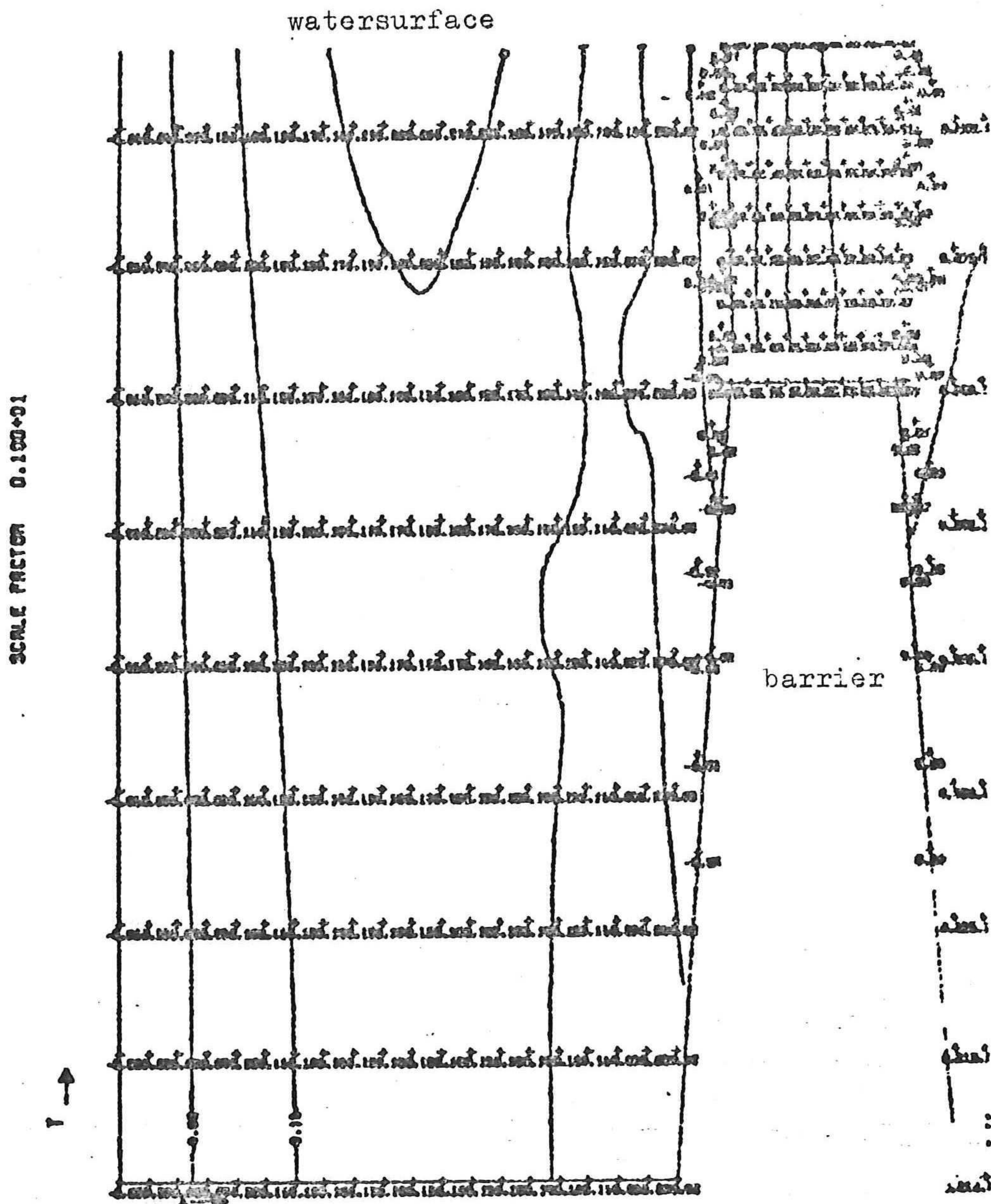


Figure 7.022

An example of a contourplot of the values of the potential ( in this case: of a section of the field of the Side View Model with a barrier height which is 0.6 times the waterdepth. )

Presentation in compressed scale.



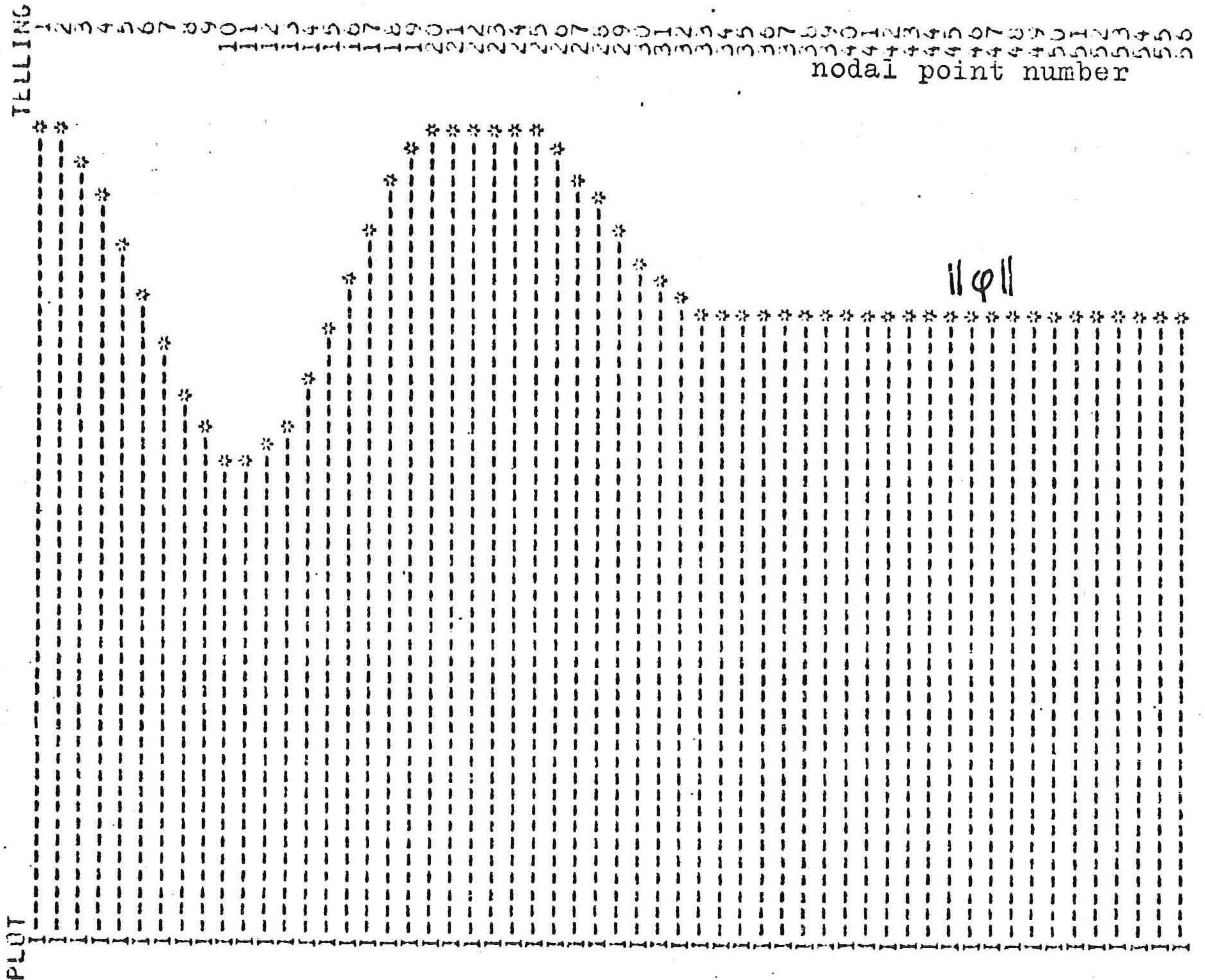
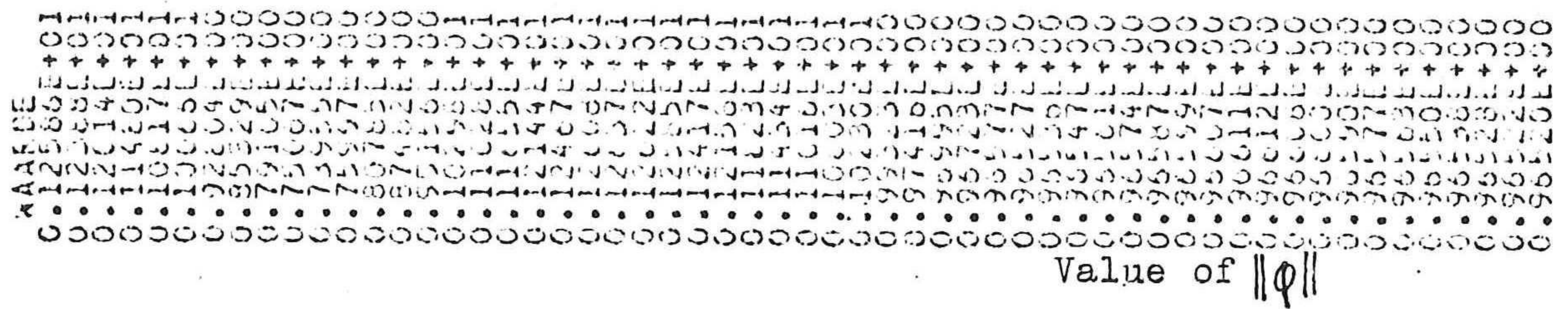


Figure 7.02.3

Example of the form of presentation used for this study.

Presentation of the values of  $\|\varphi\|$  at the nodes along the surface and a representation of these values in a graph.

For a description of the meaning of the presented characteristic parameter in this graph consult section 6.05 .





7.03            Gathering data from the graphs  
.....

The results of the computations of the Side View Model and the Top View Model all are presented in graphs of the form presented in figure ( 7.02.3 ).

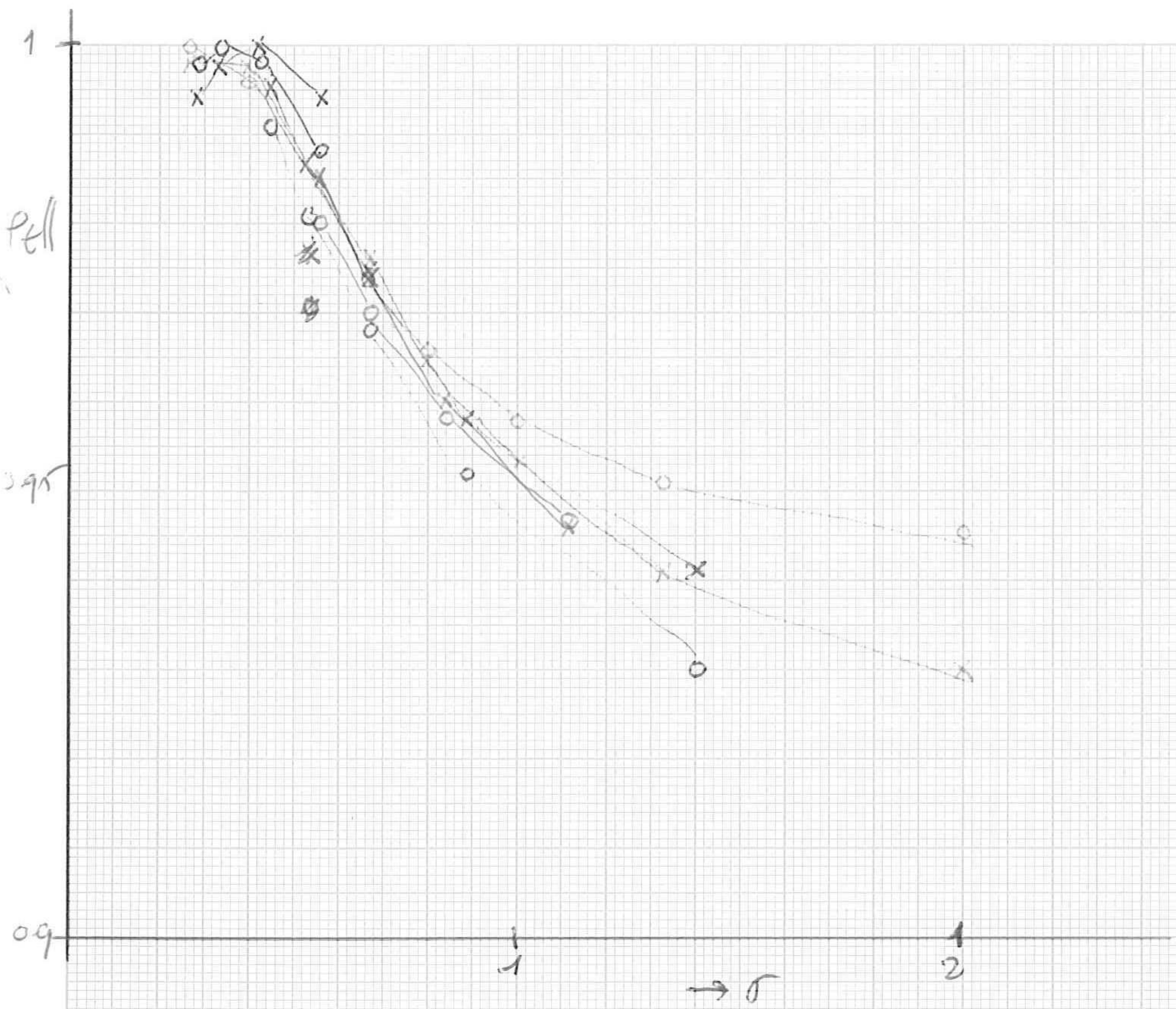
These graphs are interpreted as described in section 6.05 .

7.04            Results in tables  
.....

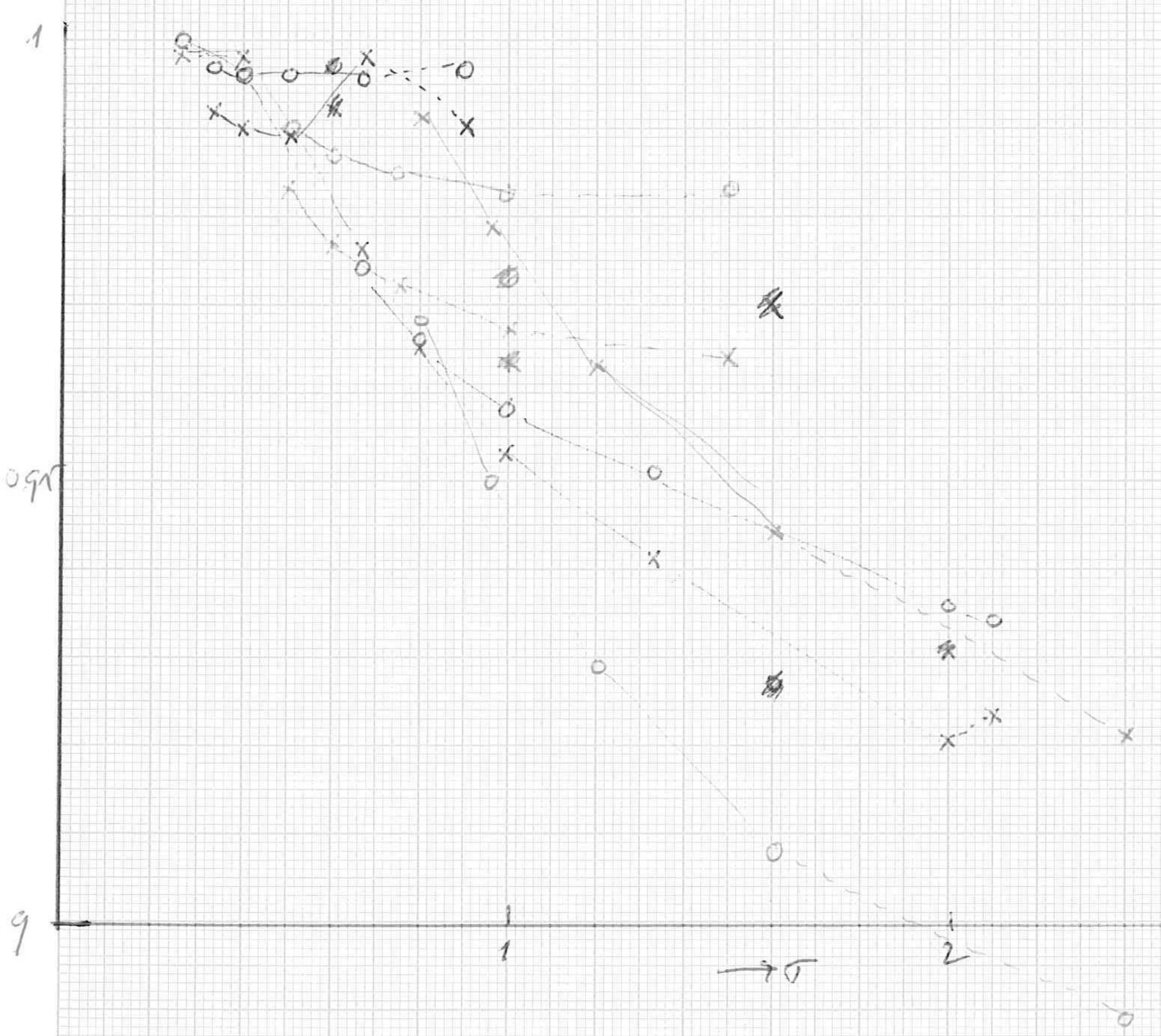
The following are the results of the computations for the cases described in section 6.07, interpreting the data as described in section 6.05 .

The numbering of the cases is the same as in section 6.07 .



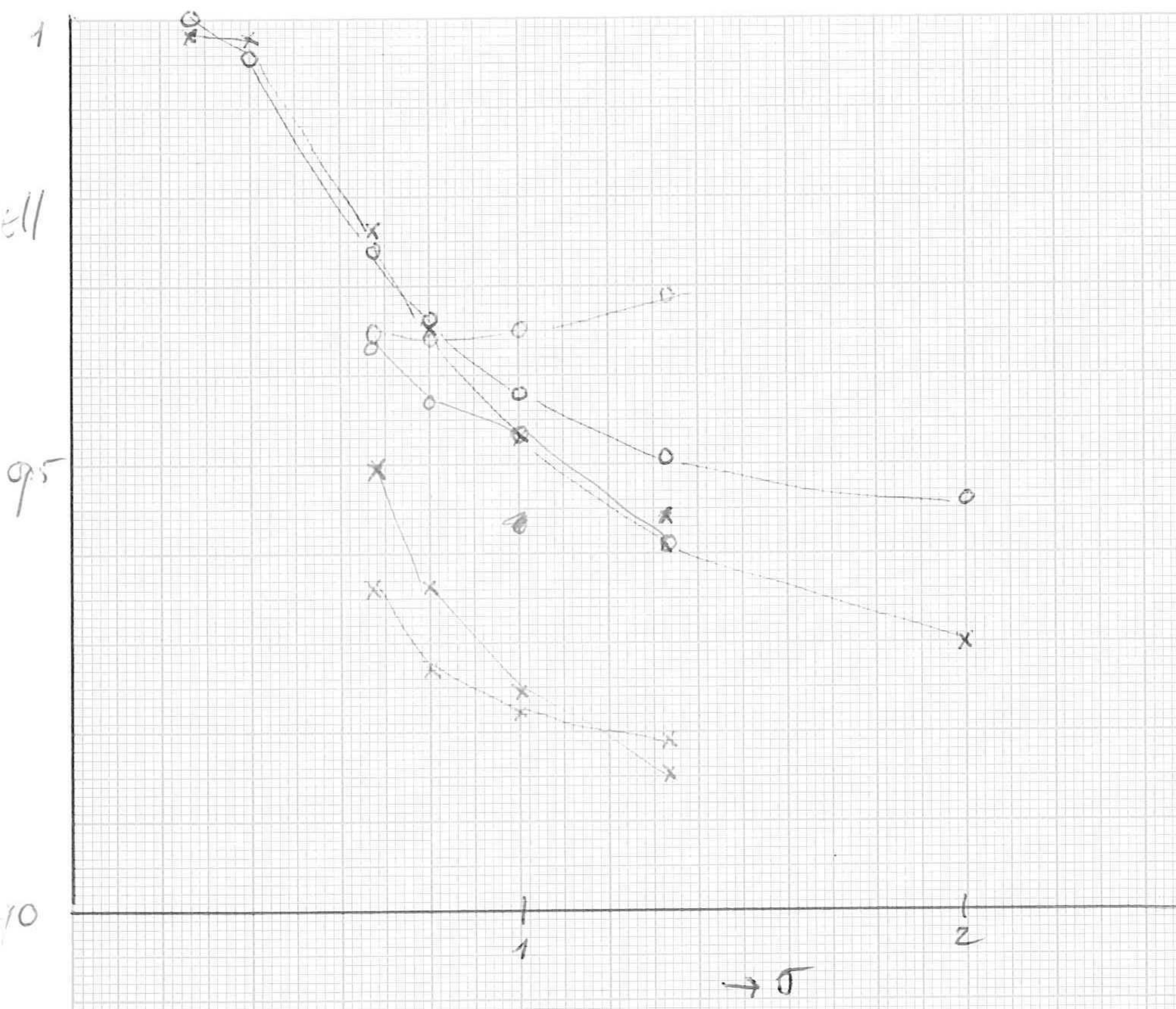


$ScD = 0.4 \quad 0.5 \quad 0.6 \quad 0.25$   
 $ratio = \frac{1}{3}$   
 $ScB = 0.6$   
 x side  
 o top  
Unlocked ScD



$ScD = 0.6$   
 $ScB = 0.6$   
 $ratio = \frac{1}{3} \quad \frac{1}{5} \quad \frac{1}{15} \quad \frac{1}{2}$   
Unlocked "ratio"





$S_{CB} = 0.6$

ratio =  $\frac{1}{3}$

$S_{CB} = 0.6 \quad 0.4 \quad 0.2$

united  $S_{CB}$



Wad in die m...

Table 7.04.1 The results of the computations, per CASE

CASE	ScD	Ratio	ScB	$\Delta D$	$\Delta x$	sigma	epsilon	$\  \varphi_{ell_s}$	$\  \varphi_{ell_r}$	$\frac{\  \varphi_{ell_r} - \  \varphi_{ell_s}}{\  \varphi_{ell_s}}$
<u>CASE 1</u>	0.4	1.0	0.6	0.00	0.2	0.00	0.00	1.000	1.000	0.000
	0.8	0.9		0.04		0.20	3.14	.999	1.000	.001
	0.7	0.8		0.08		0.40	6.28	.995	.998	.003
	0.6	0.7		0.12		0.60	9.42	.988	.994	.006
	0.5	0.6		0.16		0.80	12.57	.978	.986	.008
	0.4	0.5		0.20		1.00	15.71	.963	.973	.010
	0.3	0.4		0.24		1.20	18.85	.948	.954	.006
	0.2	0.3		0.28		1.40	21.99	.941	.930	-.012
	0.1	0.2		0.32		1.60	25.13	.970	.927	-.044
		0.1		0.36		1.80	28.27	.883	.981	.111
<u>CASE 2</u>	0.6	1.0	0.6	0.00	0.2	0.00	0.00	1.000	1.000	0.000
	0.8	0.9		0.06		0.30	3.14	.999	1.000	.001
	1.0	0.8		0.12		0.60	6.28	.996	.999	.003
	1.2	0.7		0.18		0.90	9.42	.990	.997	.007
	1.4	0.6		0.24		1.20	12.57	.979	.992	.013
	1.6	0.5		0.30		1.50	15.71	.964	.983	.020
	1.8	0.4		0.36		1.80	18.85	.944	.964	.022
	2.0	0.3		0.42		2.10	21.99	.923	.934	.012
	2.2	0.2		0.48		2.40	25.13	.921	.889	-.035
	2.4	0.1		0.54		2.70	28.27	1.000	.929	-.071
<u>CASE 3</u>	0.6	1.3	0.6	0.4	1.5	0.26	2.79	.998	1.000	0.002
		1.0			1.0	.40	4.19	.998	.996	-.002
		.6			.6	.67	6.98	.976	.974	-.002
		.5			.5	.80	8.38	.965	.966	.001
		.4			.4	1.00	10.47	.953	.958	.005
		.3			.3	1.33	13.96	.941	.951	.011
		.2			.2	2.00	20.94	.930	.946	.017
<u>CASE 4</u>	0.5	1.3	0.6	0.33	0.6	0.556	6.98	.985	.980	-.005
		0.5			0.5	.667	8.38	.974	.970	-.004
		0.4			0.4	.833	10.47	.960	.958	-.002
		0.3			0.3	1.111	13.96	.946	.947	.001
<u>CASE 5</u>	0.4	1.3	0.6	0.266	0.6	.444	6.98	.995	.991	-.004
		0.5			0.5	.533	8.38	.987	.981	-.006
		0.4			0.4	.667	10.47	.974	.968	-.006
		0.3			0.3	.889	13.96	.958	.952	-.006
<u>CASE 6</u>	0.25	1.3	0.6	0.16	0.6	0.278	6.98	.994	.998	.004
		0.5			0.5	0.333	8.38	.998	1.000	.002
		0.4			0.4	0.417	10.47	1.000	.998	-.002
		0.3			0.3	0.556	13.96	.994	.988	-.006

ScB



Table 7.04.1 The results of the computations, per CASE, continued										
<u>CASE 7</u>	ScD	Ratio	ScB	$\Delta D$	$\Delta x$	sigma	epsilon	$\ \varphi_t\ _s$	$\ \varphi_t\ _\tau$	$\frac{\ \varphi_t\ _s}{\ \varphi_t\ _\tau} - 1$
	0.6	1:5	0.6	0.48	0.6	0.80	8.38	.991	.968	-.023
					0.5	0.96	10.05	.979	.950	-.030
					0.4	1.20	12.57	.963	.929	-.035
					0.3	1.60	16.76	.944	.908	-.038
<u>CASE 8</u>										
	0.6	1:1,5	0.6	0.2	0.6	0.33	3.49	.992	.997	.005
					0.5	0.40	4.19	.990	.996	.006
					0.4	0.50	5.24	.989	.996	.007
					0.3	0.67	6.98	.998	.996	.008
<u>CASE 9</u>										
	0.6	1:2	0.6	0.30	0.6	.50	5.24	.983	.990	.007
					0.5	.60	6.28	.977	.987	.010
					0.4	.75	7.85	.972	.985	.013
					0.3	1.00	10.47	.967	.983	.017
<u>CASE 10</u>										
	0.6	1:3	0.4	0.4	0.6	0.667	6.98	.949	.963	.015
					0.5	0.80	8.38	.936	.957	.022
					0.4	1.00	10.47	.924	.953	.031
					0.3	1.333	13.96	.915	.951	.039
<u>CASE 11</u>										
	0.6	1:3	0.2	0.4	0.6	0.667	6.98	.936	.965	.031
					0.5	0.80	8.38	.927	.964	.040
					0.4	1.00	10.47	.922	.965	.047
					0.3	1.333	13.96	.919	.969	.054
<u>CASE 12</u>										
	0.6	1:3	'∞'	0.4	0.6	0.667	6.98	1.168	1.176	.007
					0.5	0.80	8.38	1.166	1.175	.008
					0.4	1.00	10.47	1.163	1.174	.009
					0.3	1.333	13.96	1.160	1.172	.010
<u>CASE A</u>										
	0.6	1:3	0.6	0.4	0.3	1.333	extended upwave fieldlength .939			
<u>CASE B</u>										
	0.6	1:3	0.6	0.4	0.3	1.333	double stepsize .934			

Discussion of these testcases is done in the next chapter.



8.00            Introduction

The results of table ( 7.04.1 ) are interpreted in this chapter.

The predicted values of  $\|\varphi_t\|$  are displayed in graphs for both the Side View Model and the Top View Model.

The predictions of the two models, compared using the parameter  $\frac{\|\varphi_t\|_T}{\|\varphi_t\|_S} - 1$  , are also displayed in graphs.

The effect of the variation of the parameters ScD, Ratio and ScB is investigated.

This leads to a discussion of the results in the next chapter.

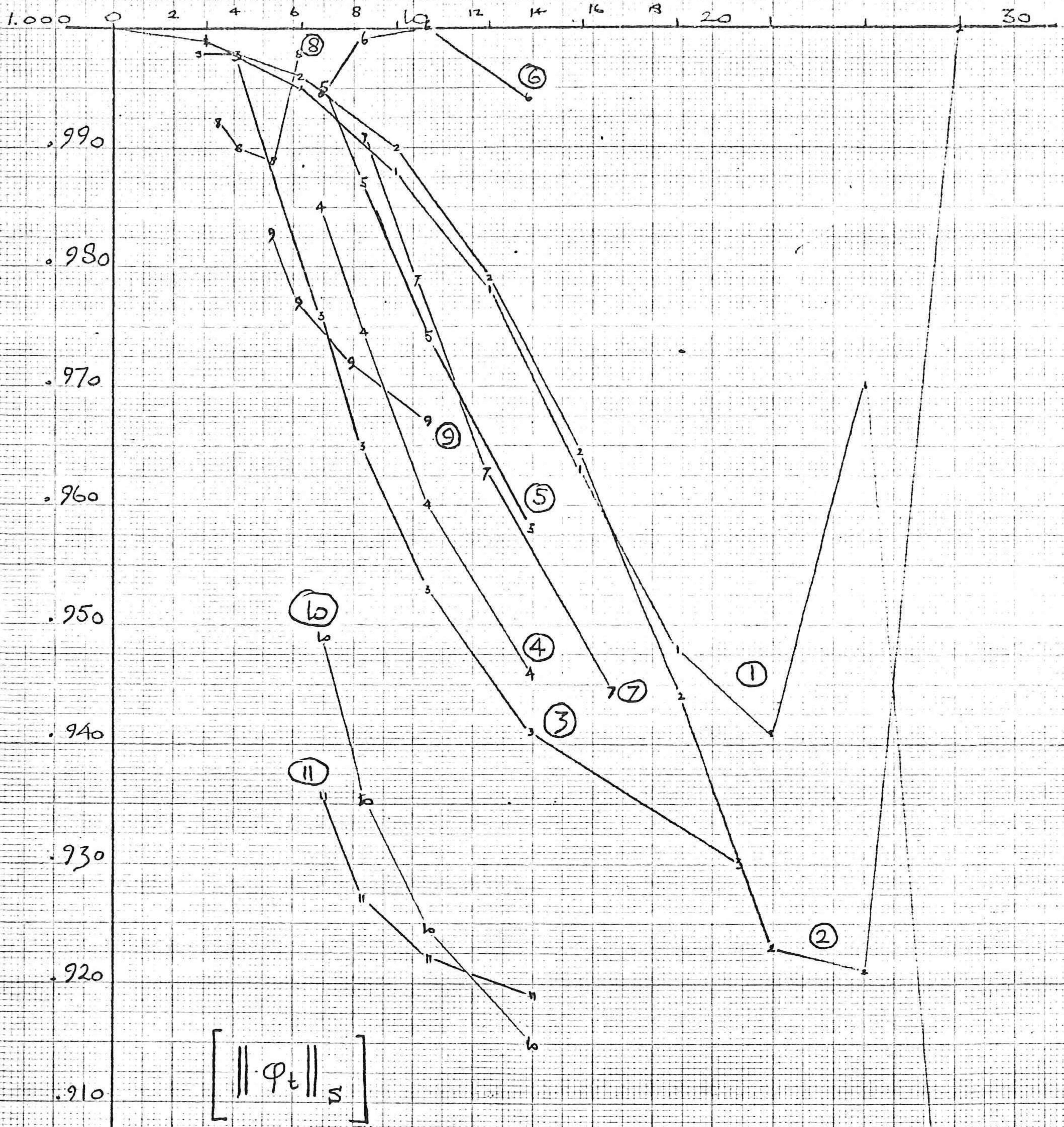
8.01            Presentation per model and per CASE

The predictions of  $\|\varphi_t\|$  have been presented in graph ( 8.01.1 ) for the Side View Model and in graph ( 8.01.2 ) for the Top View Model.

The results have also been grouped per CASE, as displayed in graph ( 8.01.3 ). For the ease of the comparison the curves have been untangled, providing better insight in the relative predictions per case.



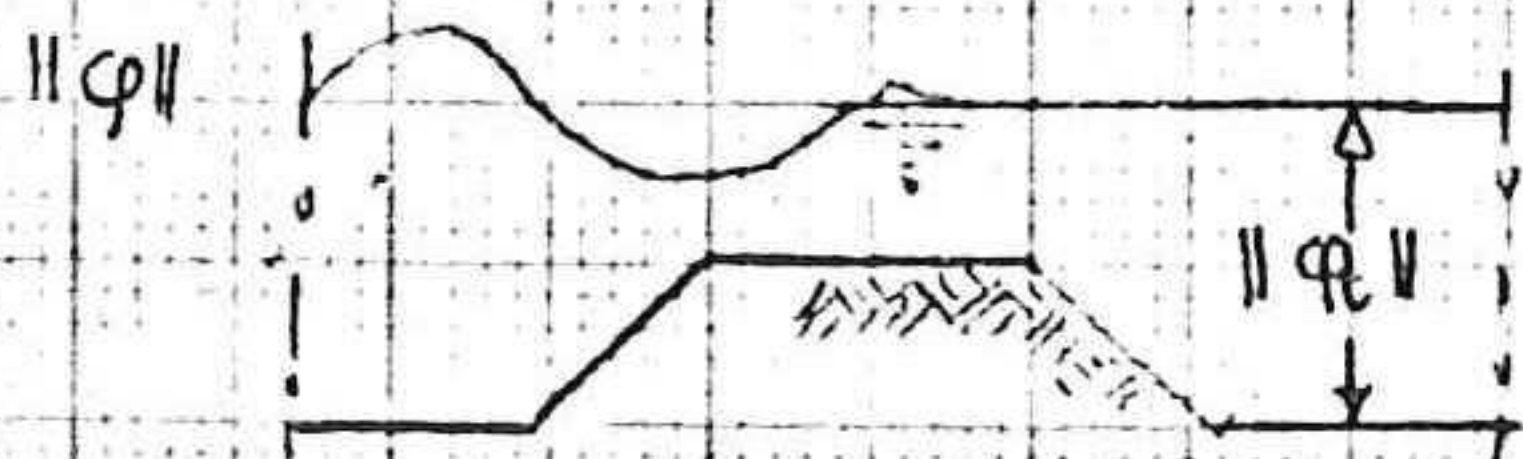
$$\left[ \varepsilon = \frac{\sigma \lambda_0}{D} \right]$$



$$\left[ \parallel \varphi_t \parallel_s \right]$$

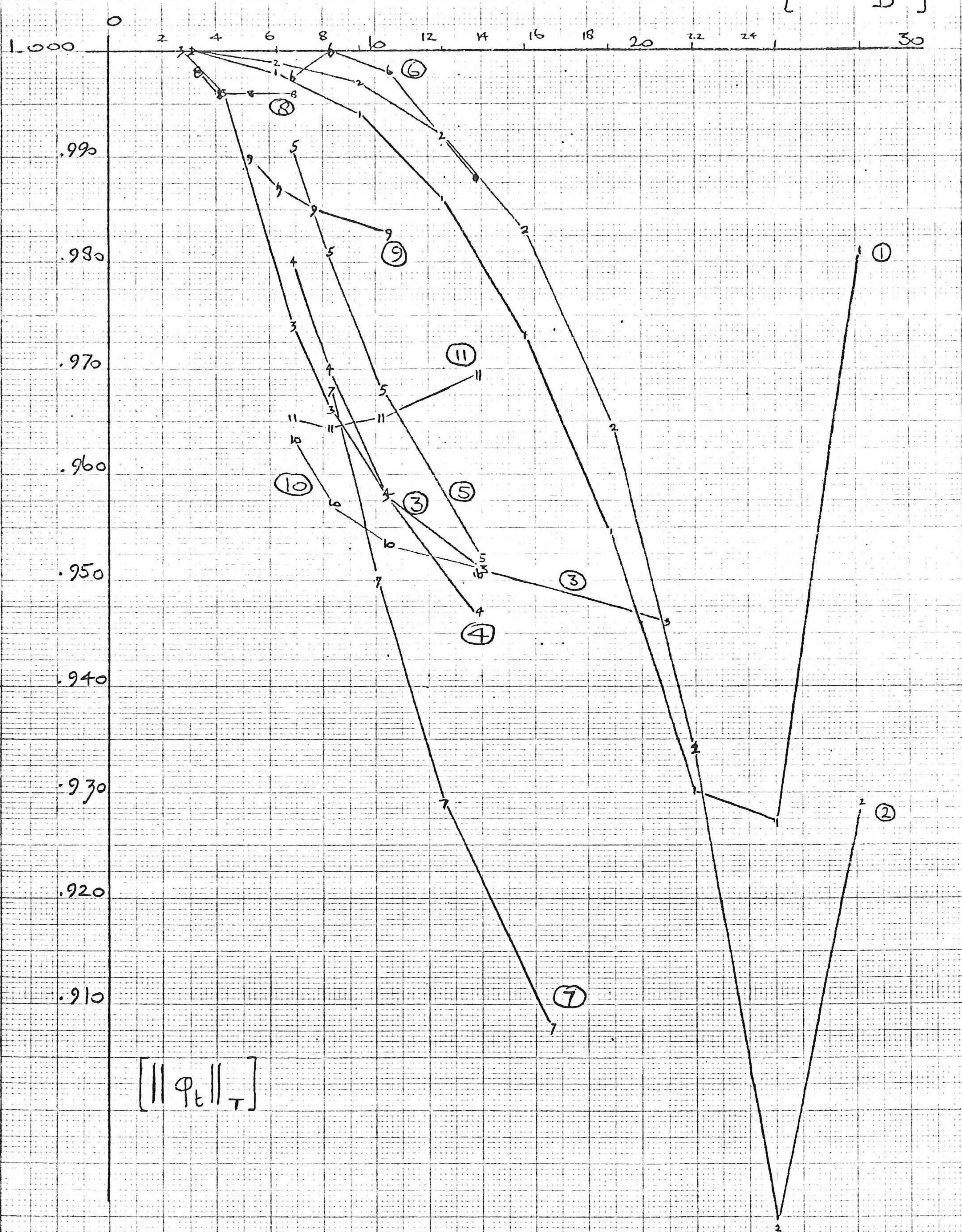
Graph 8.01.1

Predicted values, by the Side View Model, of the potential of the transmitted wave for the cases 1 through 11. ( see diagram 6.07.1 )





$$\left[ \epsilon = \frac{\sigma \lambda_0}{D} \right]$$



Graph 8.01.2

Predicted values, by the Top View Model,  
of the potential of the transmitted wave.



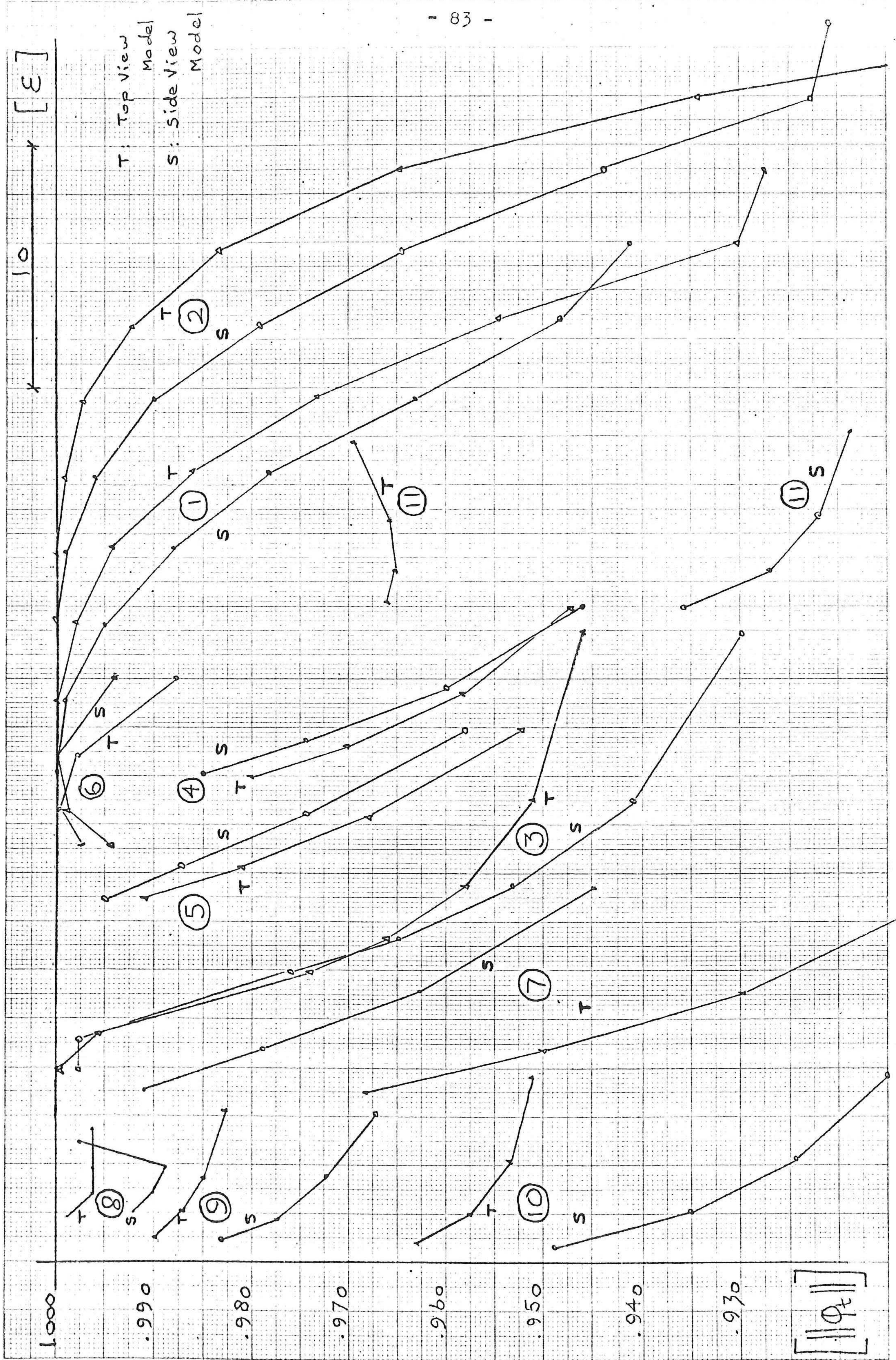


Figure 8.01.3 Comparative presentation of predictions per case.



8.02 Evaluation of the predictions of  $\|\varphi_t\|$

From the data in graphs ( 8.01.1 ) through ( 8.01.3 ) insight might be obtained concerning the propagation of a wave over a trapezoidal dam. As can be seen in graph ( 8.01.3 ) remarkable differences exist between the predictions by the two models. Since relatively little information is available of the characteristics of the propagation of waves over a trapezoidal dam it requires a special study, outside of the scope of this study, to further evaluate this information.

Only for CASE 1 and 2 it may be concluded that the behaviour which is predicted by the computation is in accordance with the behaviour described in other studies. Mobarek ( 1974 ) has done hydraulic scale model experiments for the wave propagation over a trapezoidal dam and found that the coefficient of transmission varies with the parameter Ratio as depicted in figure ( 8.02.1 ).

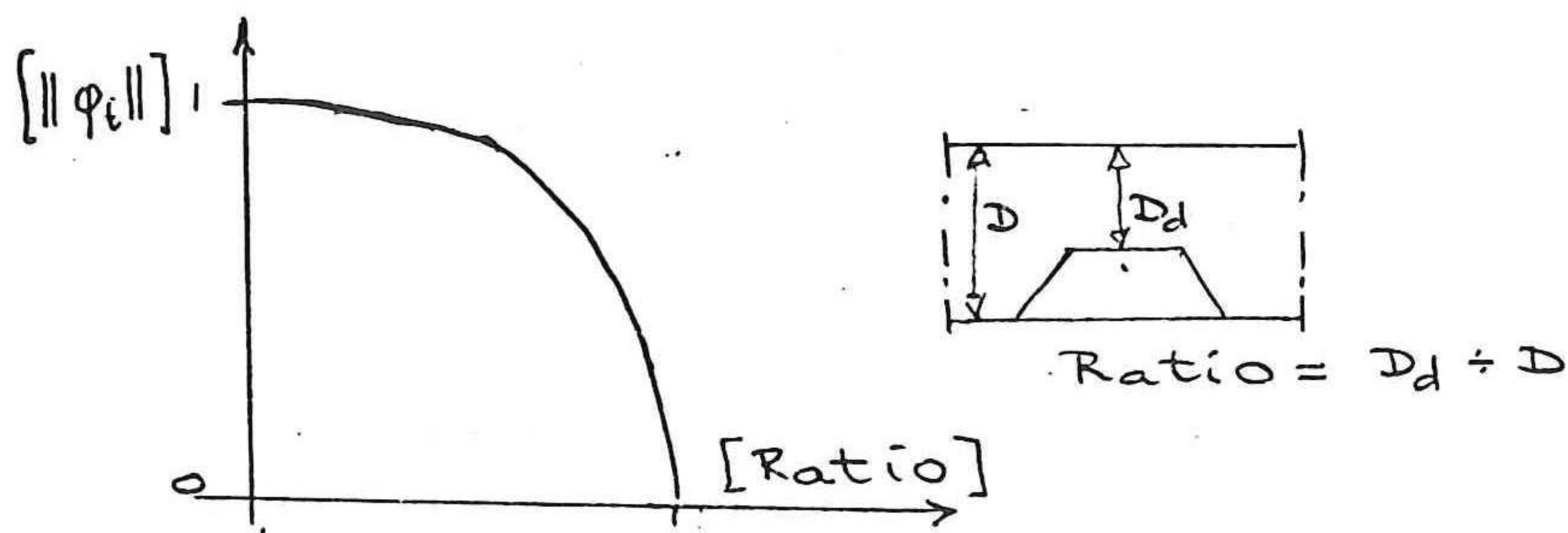


Figure 8.02.1

Schematic representation of the characteristic behaviour of transmission over a trapezoidal dam.

As can be seen in graph ( 8.01.3 ) this is indeed the form of the computed coefficients of transmission. Only for low values of Ratio do the predictions deviate from the



expected behaviour. (  $D_d$  less than 0.4 D, approximately )  
According to the data of Mobarek nonlinear behaviour will  
have to be taken into account for these values of Ratio,  
which means that the computer model could not be used  
for these cases.

8.03 Evaluation of the values of  $\frac{\|\varphi_t\|_r}{\|\varphi_t\|_s} - 1$ .

Graph ( 8.03.1 ) presents the comparison of the predicted  
values of  $\|\varphi_t\|$  per model using the parameter  $\frac{\|\varphi_t\|_r}{\|\varphi_t\|_s} - 1$   
mapped versus  $\varepsilon$ .

The errorband in the graph has magnitude  $4 \cdot 10^{-3}$ .

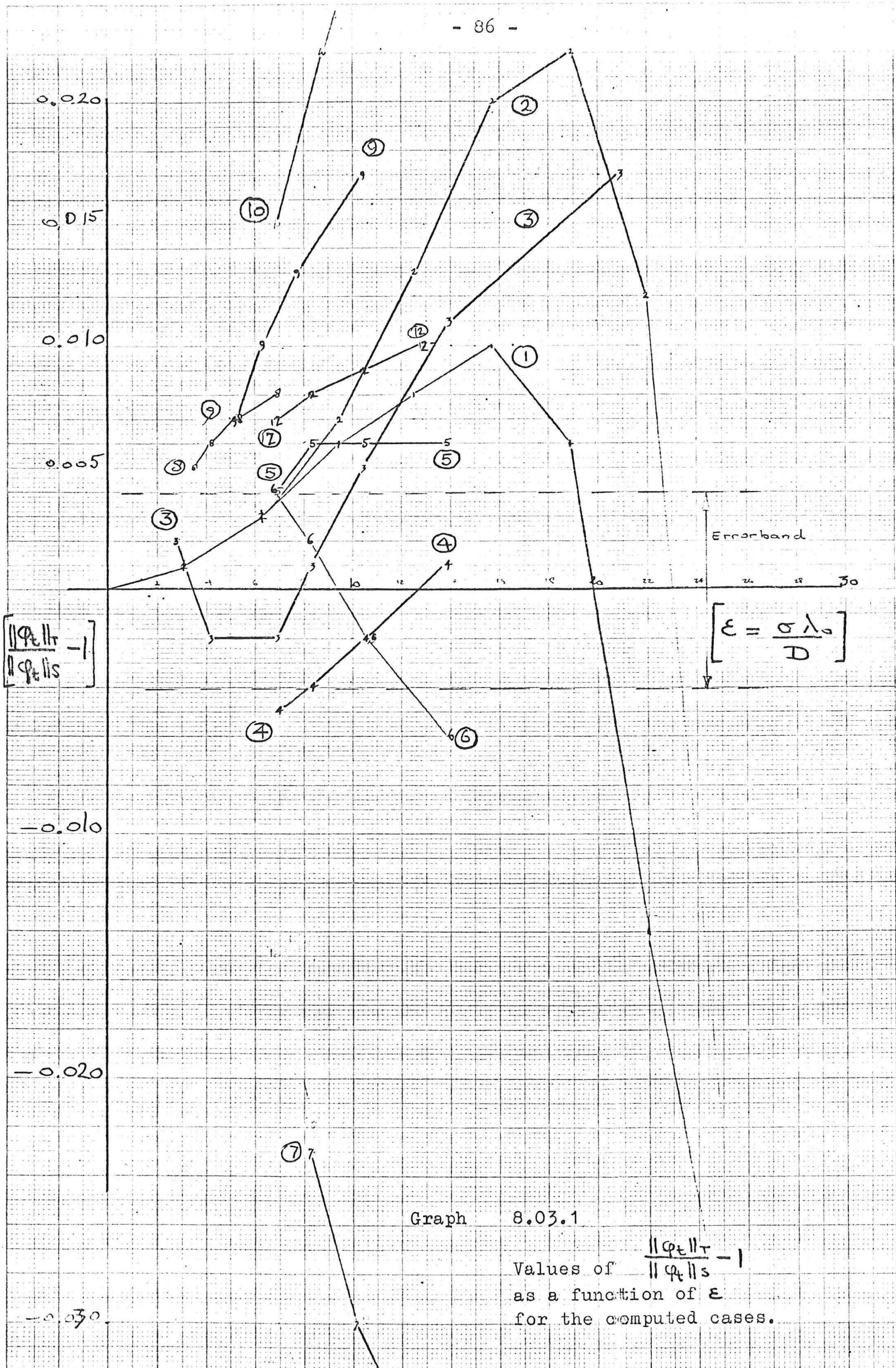
( See Appendix 6.)

Within the errorband the results of the two models will  
be called " equal ".

It is clear from the results in graph ( 8.03.1 ) that  
there is no straightforward relationship between the  
parameter  $\varepsilon$  and the difference in the predictions by the  
two models.

The results of the various CASEs will be compared in the  
next chapter.





Graph 8.03.1

Values of  $\frac{\|\varphi_t\|_T}{\|\varphi_t\|_S} - 1$  as a function of  $\epsilon$  for the computed cases.



9.00      Introduction

In this chapter the results will be discussed.

In the derivation of the refraction - diffraction equation ( see Appendix 3 ) the most important assumption made by Berkhoff was the ' near flatness ' of the bottom, for which the parametric requirement is:

$$\epsilon = \frac{\sigma \lambda_0}{D} \quad \text{must be small.}$$

The aim of this study is: to determine the limit of applicability of the refraction - diffraction equation.

For this purpose a computer program has been made by which the values of the wave potential at the water surface can be determined using the refraction - diffraction equation ( Top View Model ). Also the program computes the values of the wave potential using the set of equations from which the refraction - diffraction equation has been derived ( The Side View Model ).

This leads to the practical reinterpretation of the aim of this study: to determine when the predictions of the Side View Model and the Top View Model will be significantly different, whereby the expectation exists that the difference will be related to the value of  $\epsilon$ .

The results ( Graph 8.03.1 ) show that there is no direct relationship between the difference in the predictions and the value of  $\epsilon$ .



In this chapter the influence of the other parameters (  $ScD$ , Ratio and  $ScB$  ) will be investigated. These influences will be compared with the results of CASE 3.

#### 9.01 Results of CASE 3

A case in which the slope is the only parameter which is varied, is CASE 3 ( see diagram 6.07.1 ). The results for this case are for convenience repeated in graph ( 9.01.1 ).

The results of this case seem to confirm the expectations: with an increase of  $\xi$  the difference in the predictions of the two models increases.

The difference becomes significant, using the criterion defined in Appendix 6 , for a value of  $\xi = 10$  .

This result, however, is to be considered with great caution.

The results for the other cases ( graph 8.03.1 ) show that this behaviour is not at all characteristic.

In the next section the influence of the other parameters will be studied, for  $\xi$  having values of approximately  $\xi = 10$ .



$$\left[ \frac{\| \varphi_{t, \text{tr}} \|}{\| \varphi_{t, \text{lls}} \|} - 1 \right] \uparrow$$

0.008  
 0.006  
 0.004  
 0.002  
 0.002  
 -0.004

CASE 3  
 $S_{cD} = 0.6$   
 Ratio = 1:3  
 $S_{cB} = 0.6$

[B] →

20

15

10

5

Error band

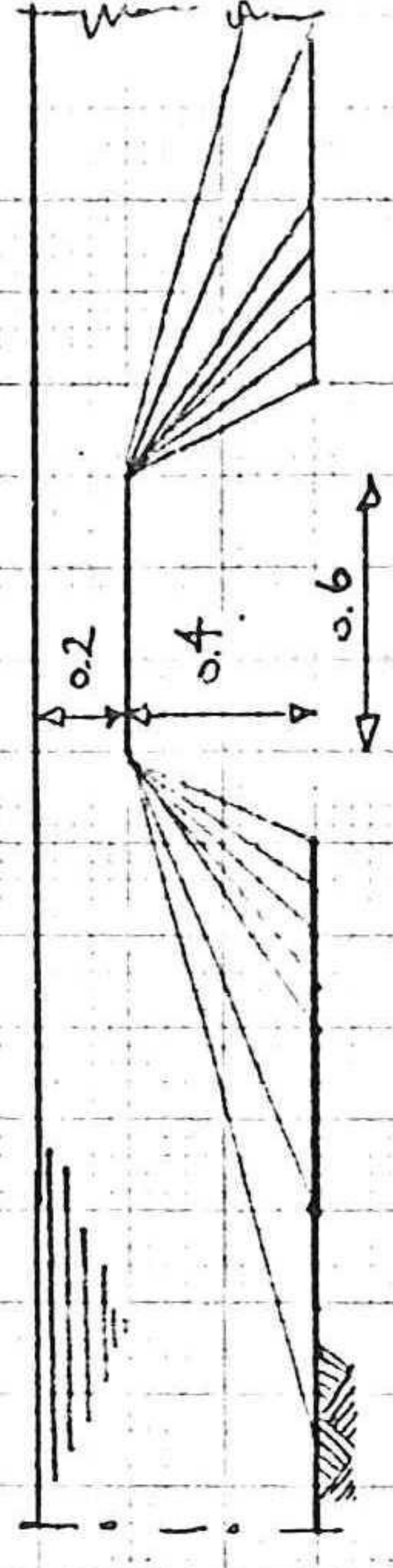


Figure 8.01.1 Presentation of the results of CASE 3



9.02 Influence of the parameters ScD, Ratio and ScB

The cases which differ only in the value of ScD are presented in graph ( 9.02.1 ).

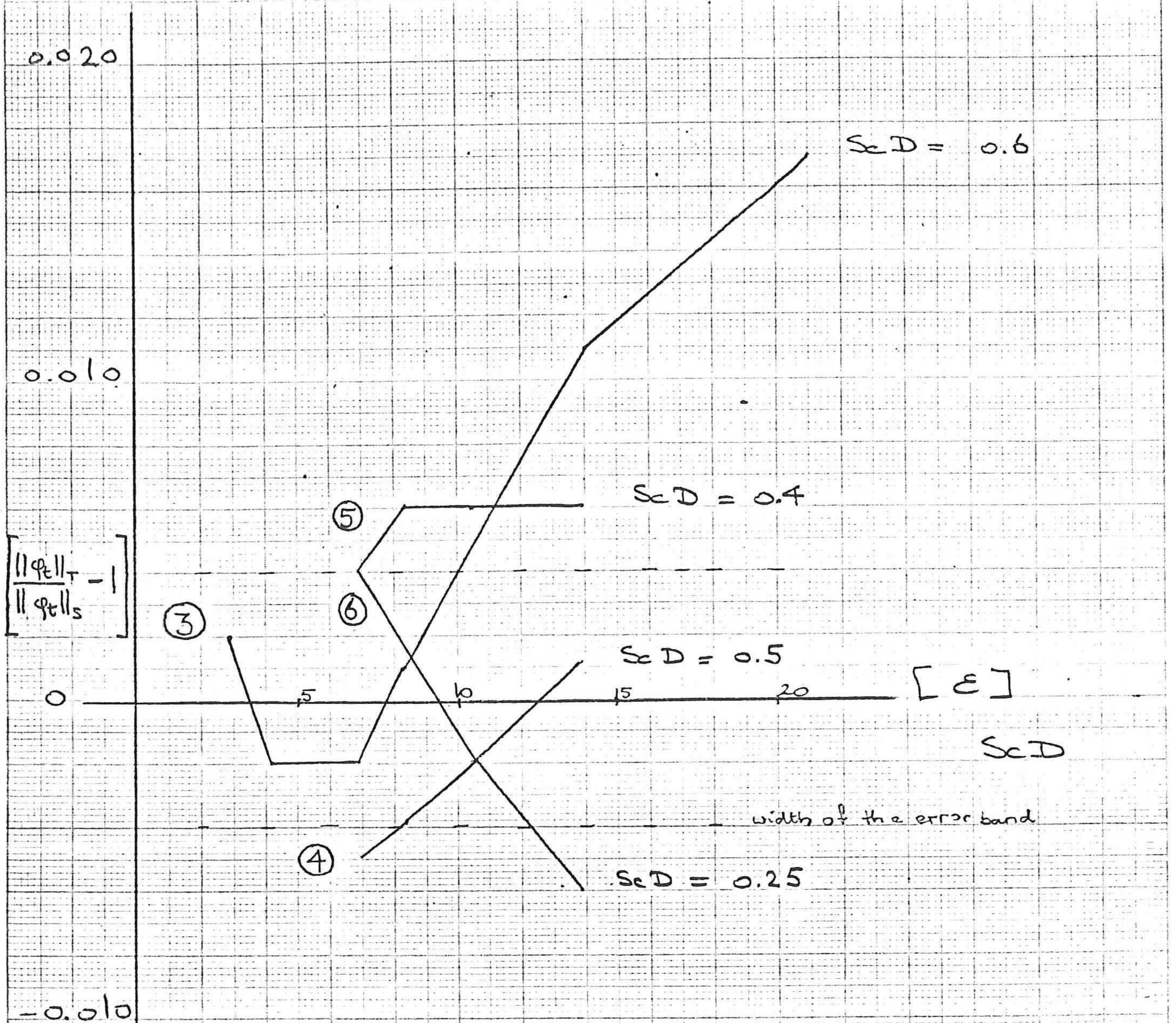
It is evident that variation of ScD greatly influences the results and that from the results in graph ( 9.02.1 ) no regularity can be discovered in the way the variation of the ScD affects the results.

The same is the case for those CASEs which only differ in the value of Ratio ( graph 9.02.2 ) respectively in the value of ScB ( graph 9.02.3 ).

Notably the results of graph ( 9.02.3 ) suggests that local effects have to be accounted for in the comparison of the two models ( and thus: in the search for the limits of applicability of the refraction - diffraction equation ).

The result of CASE 7 ( Ratio = 1.5 ) falls in the interval of nonlinear behaviour described in section 8.02, where the predictions for CASEs 1 and 2 show aberrational behaviour. Perhaps this is related to the fact that the result for this case differs greatly from the other results in graph ( 9.02.2 ).

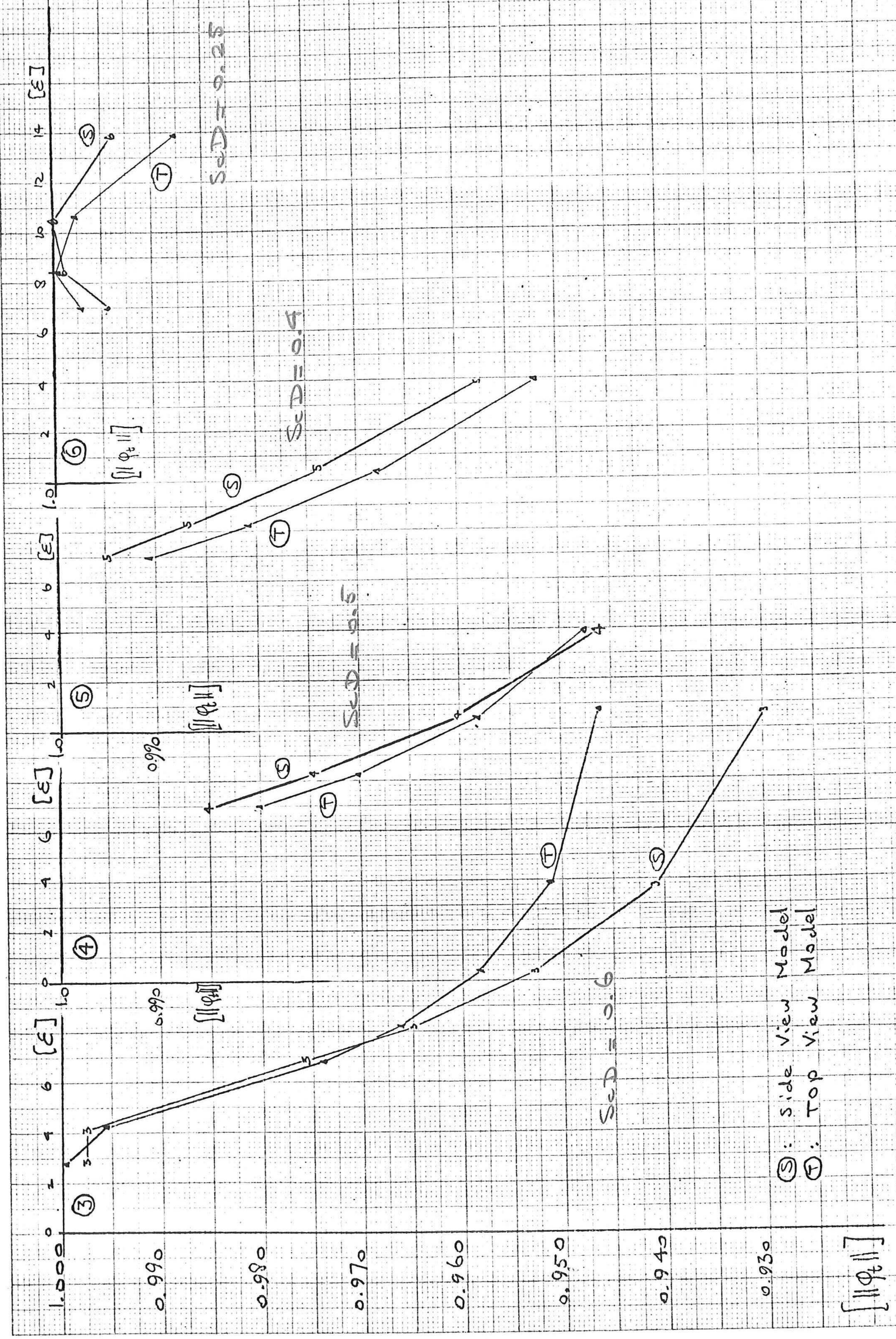




Graph 9.02.1

Results for CASEs 3, 4, 5 and 6  
( which differ only  
in the value of  $ScD$  )

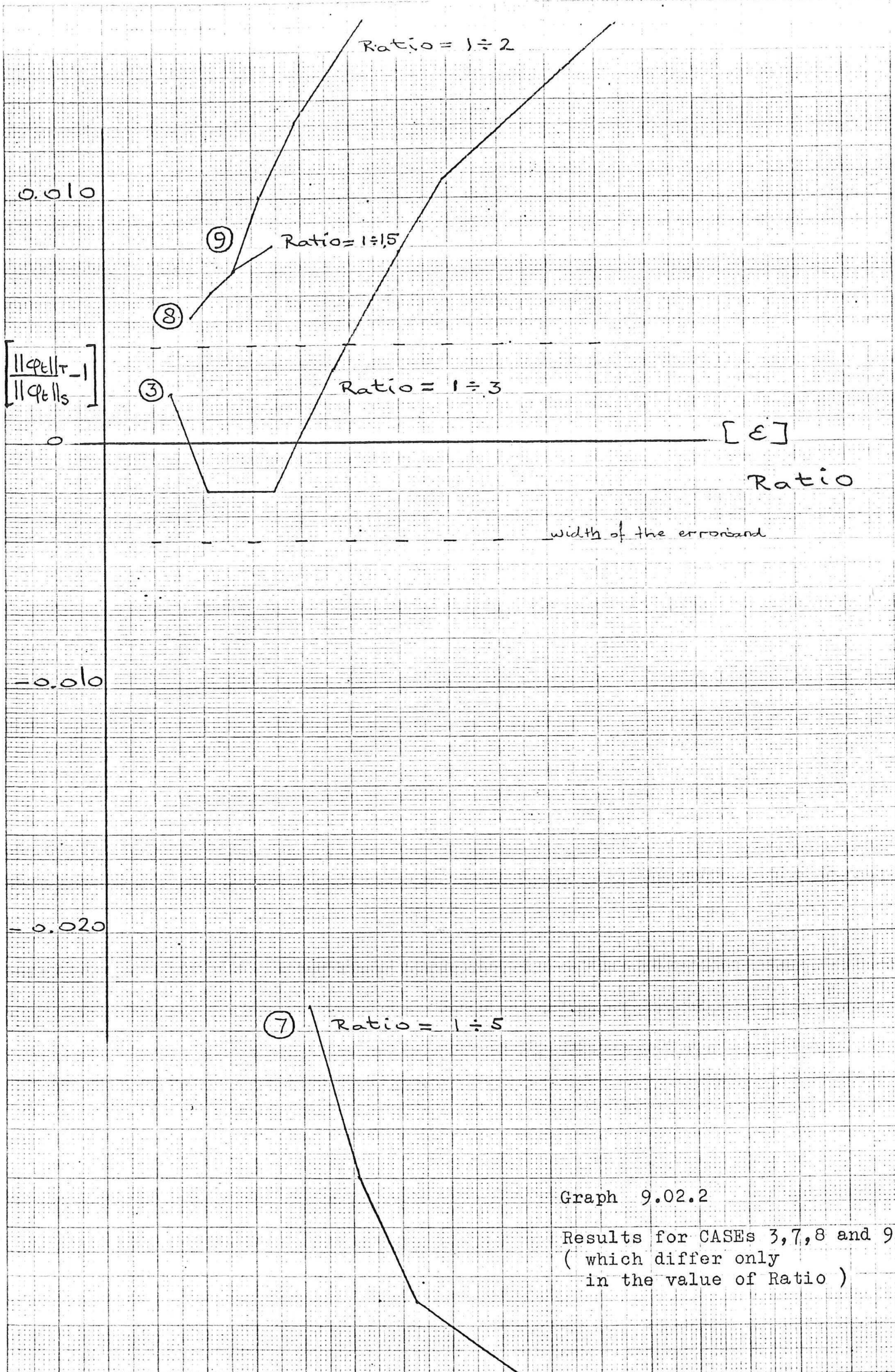




(S): Side View Model  
 (T): Top View Model

Graph 9.02.1A Compared values of  $||\varphi_e||$ , for variation of SED

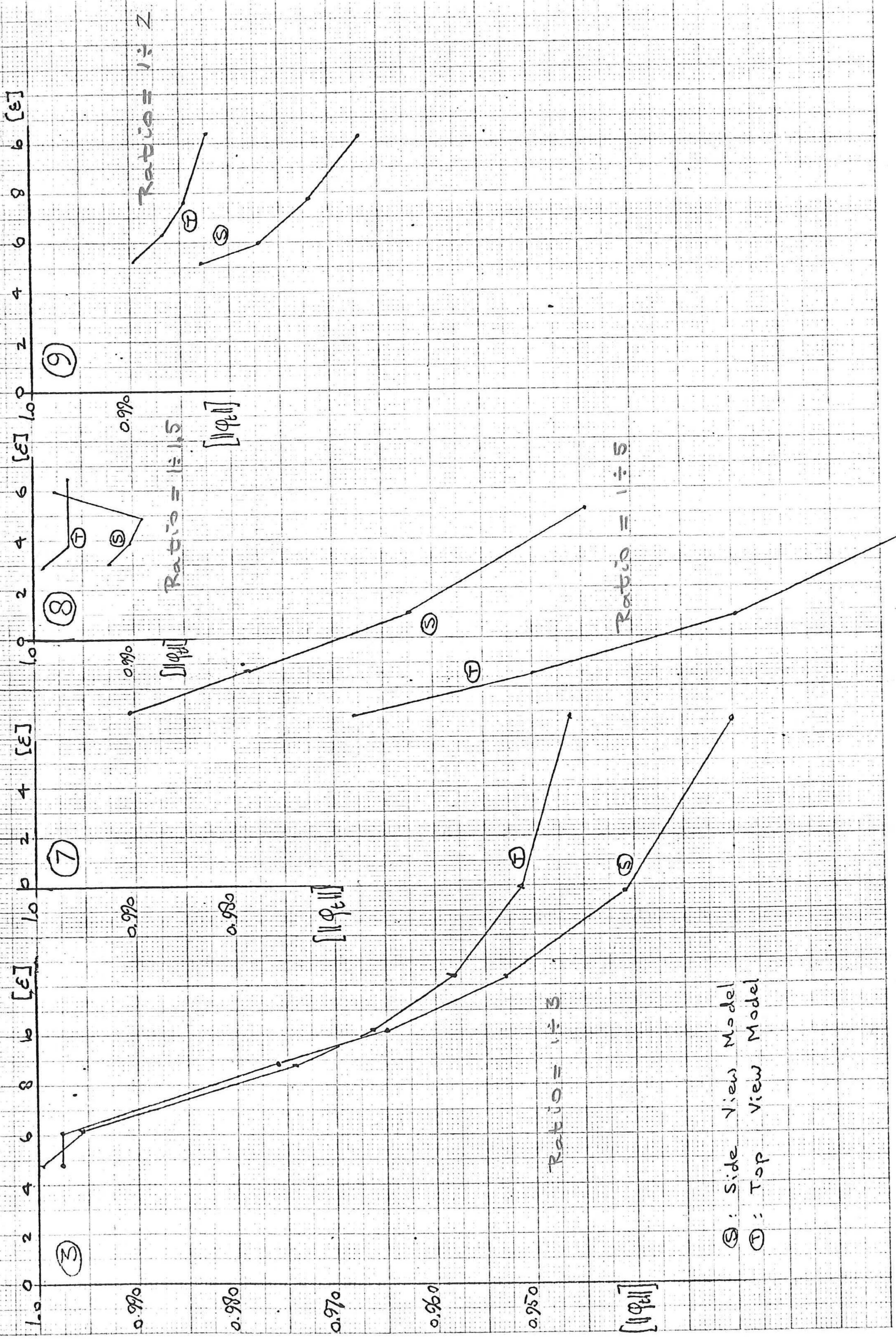




Graph 9.02.2

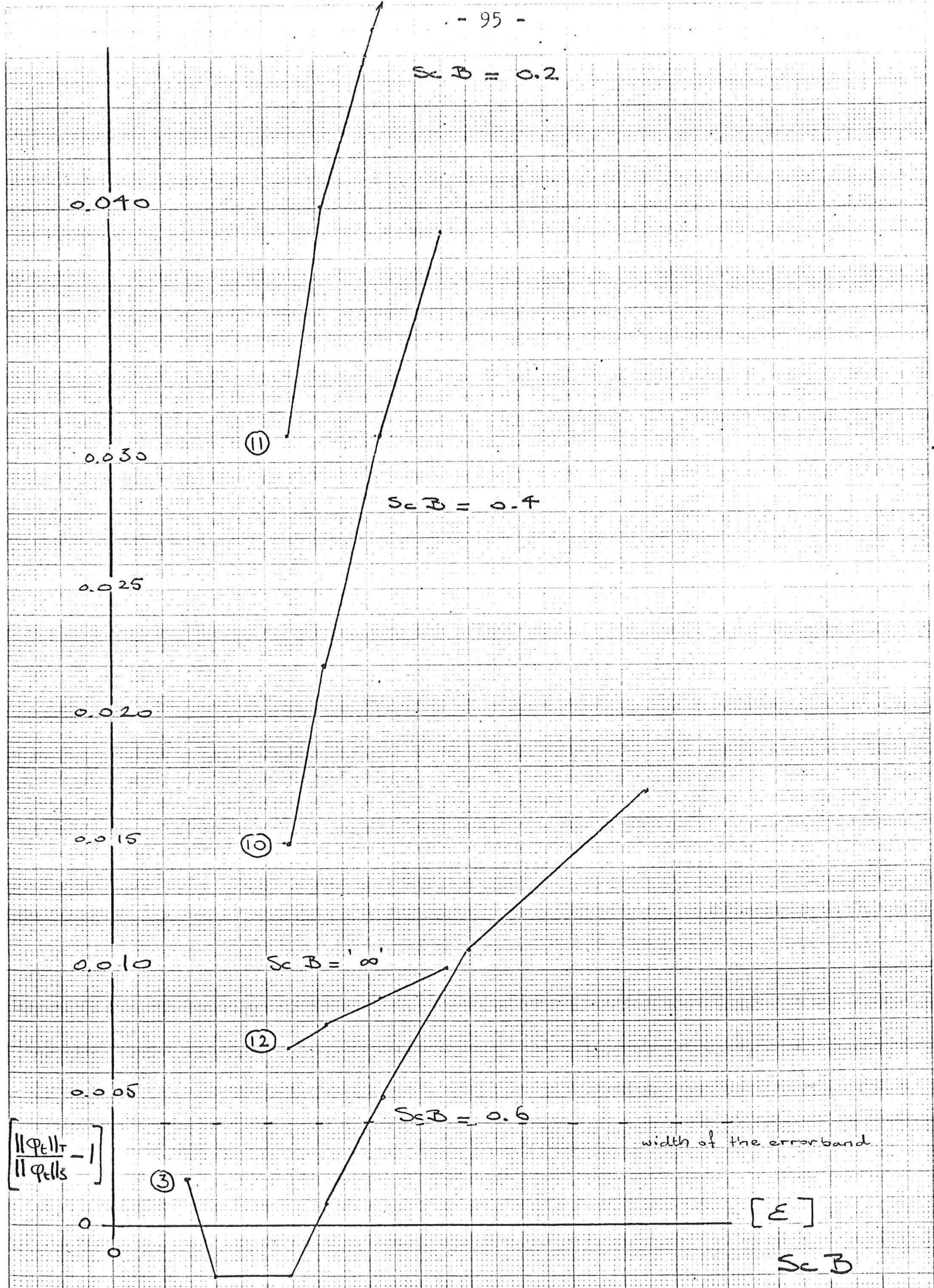
Results for CASEs 3, 7, 8 and 9  
( which differ only  
in the value of Ratio )





Graph 9.02.2A Compared values of  $1/\rho_t$  for variation of Ratio

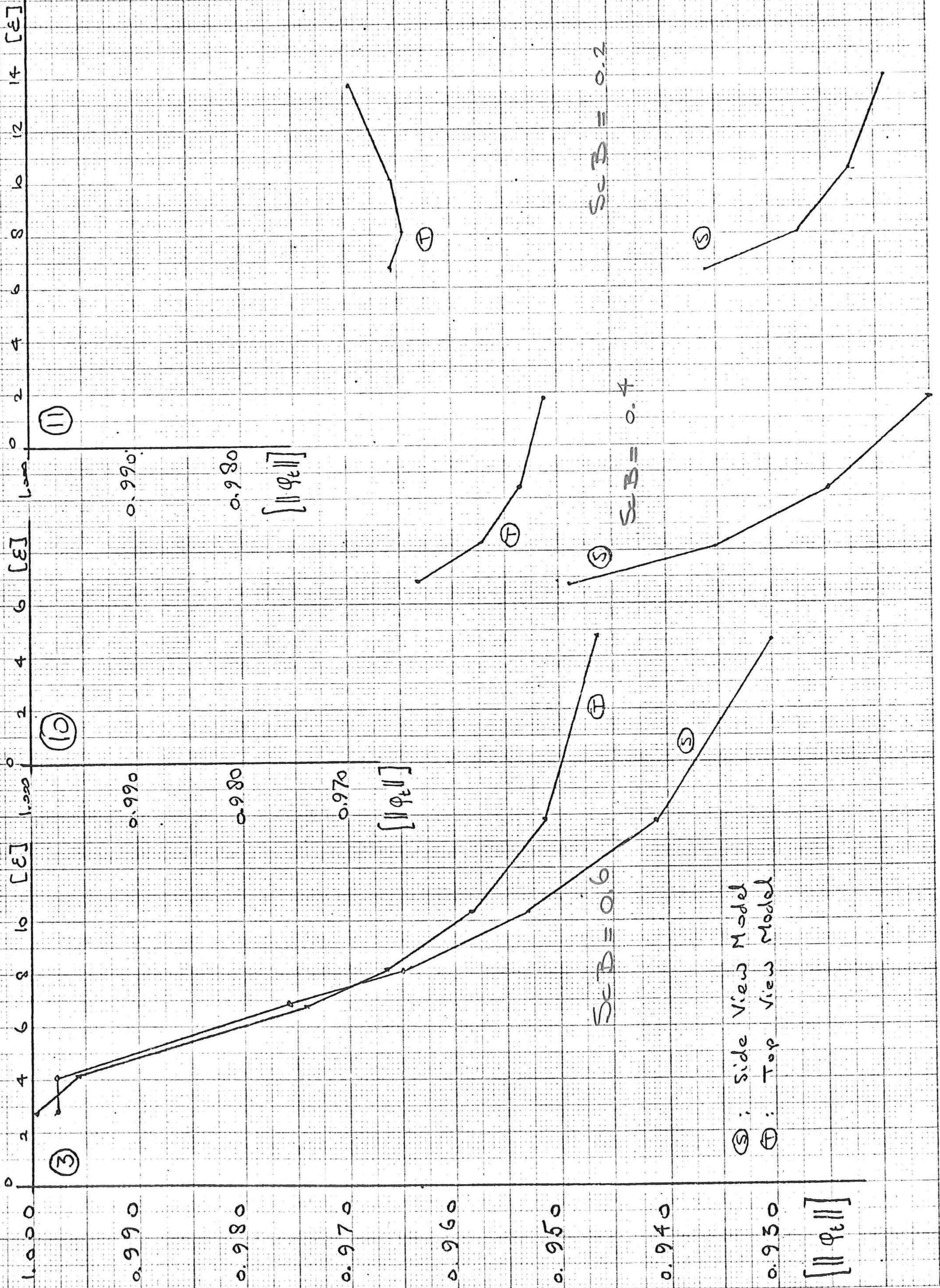




Graph 9.02.3

Results for CASES 3, 10, 11 and 12  
( which differ only in the value of ScB )





Graph 9.02.3A Compared values of  $\mu$  for variation of  $ScB$



This study has been undertaken in order to find out what is to be understood by a 'nearly flat bottom' which forms the basic assumption in Berkhoff's derivation of the refraction - diffraction equation.

That is, it is an inquiry into the limits of applicability of the refraction - diffraction equation.

For the study a computational program has been made which calculates the wave potential at the water surface using the refraction - diffraction equation ( Top View Model ) and also using the set of equations from which the refraction - diffraction equation has been derived, based on the Laplace equation ( Side View Model ).

The magnitude of the transmitted wave potential is used as the characteristic parameter:  $\|\varphi_t\|$ .

The results of the computation are compared in the form:

$$\frac{\|\varphi_t\|_T}{\|\varphi_t\|_S} - 1$$

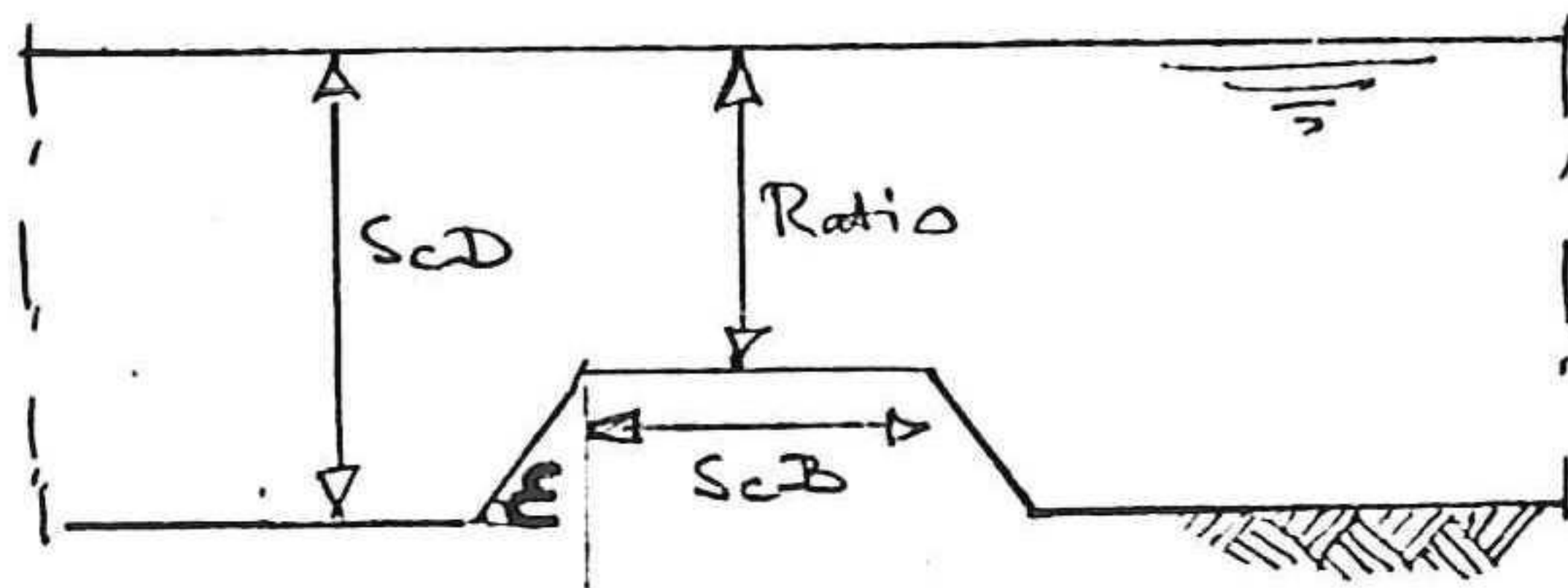
A measure of significance of the difference of the predictions is determined.

The expectation was that a correlation would exist between the characteristic parameter of the bottom slope,  $\xi$ , and the discrepancy in the predictions by the two models.

This correlation does indeed exist for one specific case ( CASE 3 of the examined cases ), but the behaviour is not at all characteristic.



It follows from the study that local parameters, characteristic for the bottom configuration, are of importance:



ScD : waterdepth  
Ratio : relative water-  
depth over dam  
ScB : crestwidth  
 $\xi$  : bottomslope

Due to the influence of these parameters it has not been possible to reach the goal of this study and to determine limits of applicability of the refraction - diffraction equation.

Further study will be required to find a more suitable formulation of the characteristic parameter of the bottom slope, incorporating local characteristics of the bottom profile.

A major spin off of this study has been the development of a standard element of AFEP which contains a two dimensional description of the refraction - diffraction equation with general first order boundary conditions ( complex number plane ).

This model can handle a very broad set of bottom configurations.



---

This study has produced a very versatile and generally useful computer program for the calculation of wave potential using either the Laplace equation or the refraction - diffraction equation with general first order boundary conditions in the complex plane.

It is recommended that the program be tested for its numerical behaviour ( stability, optimal stepsize, etc.) for the various general cases for which it may be used so that it can be used as a standard program, without requiring each user to go through the tedious and costly actions of numerical testing.

It is recommended that the characteristic parameter of the bottom profile,  $\epsilon$ , be reformulated in such a way that local behaviour of the bottom configuration is accounted for.

Once this has been done it will be possible to try to determine the limits of applicability of the refraction - diffraction equation.

It is recommended that the TOP VIEW MODEL ( ELM301 ) be checked against experimental results, two dimensionally.



It is recommended that the Top View Model be used to determine the characteristics of wave transmission over a submerged barrier for oblique wave incidence. The existing information is scarce, and describes mainly experimental results on scale models for simple cases ( of orthogonal wave incidence ).

Comparison of the two dimensional Top View Model against a three dimensional 'Laplace model' will provide more insight into the practicality of the two models, especially when compared to experimental results or field data.

It is recommended that the models used in this study be " polished ", making them fit for general usage.

- economisation of storage and time.
- mesh generation considering the required stepsize as follows from numerical requirements and the anticipated wavelength.
- adapting the boundary conditions for oblique waves.

It is also recommended that the initial phase of the students' encounters with the computer be closely guided. This greatly increases the effectiveness of operating with the computer and saves much time and money.



I am truly grateful for the freedom which I have been given in undertaking this study and I thank professor Battjes and professor Vreugdenhil for it.

Nico Booy has been a very valuable advisor and I have greatly appreciated his helpful criticism and comments.

Thank you very much, Guus Segal, for the cheerful help and support you gave when I fled the gloom of computer-mania.







Appendix 1

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Derivation of the free surface boundary condition

The surface of the fluid varies in place and time. The general description of a time and place dependent variation, for any function  $\Phi$  has the form:

$$\begin{aligned} \frac{d\Phi}{dt} &= \frac{\partial\Phi}{\partial t} + \frac{\partial\Phi}{\partial x} \cdot \frac{dx}{dt} + \frac{\partial\Phi}{\partial y} \cdot \frac{dy}{dt} + \frac{\partial\Phi}{\partial z} \cdot \frac{dz}{dt} = \\ &= \frac{\partial\Phi}{\partial t} + \nabla\Phi \cdot \vec{v} \end{aligned} \quad ( 2.1 )$$

In a field where the variation of momentum ( $\Phi =: \rho v_i$ ) is the characteristic parameter and where conservation of momentum may be assumed ( $\frac{d\rho v_i}{dt} = 0$ ) the equation expands to the Navier Stokes equation ( Hinze, 1975; R. Booy, 1978 ):

$$\frac{\partial \rho v_i}{\partial t} + (\nabla \rho \vec{v}_i) \vec{v}_j + \rho (\vec{v}_i \nabla) \vec{v}_i + \nabla p_i - \mu \nabla^2 v_i - \rho F = 0 \quad ( 2.2 )$$

where: p : pressure  
F : force field

- using
- 1)  $(\vec{v} \cdot \nabla) \vec{v} = \frac{1}{2} \nabla(\vec{v} \cdot \vec{v}) - \vec{v} \times (\nabla \times \vec{v})$
  - 2)  $\nabla \times \vec{v} \stackrel{As}{=} 0$  ( no rotation: ( 2.01.5) )
  - 3)  $F \stackrel{As}{=} -\nabla \Omega$  force potential

this becomes:

$$\frac{\partial \rho v_i}{\partial t} + (\nabla \rho \vec{v}_i) \vec{v}_j + \rho \left( \frac{1}{2} \nabla(\vec{v}_i \cdot \vec{v}_i) \right) + \nabla p_i + \rho \nabla \Omega = 0 \quad ( 2.3 )$$



using 4)  $\frac{d\rho}{dt} = \frac{\partial \rho}{\partial t} + \nabla \rho \cdot \vec{v} \stackrel{\text{AS}}{=} 0$  uniform and homogenous density distribution

5)  $\nabla \rho \stackrel{\text{AS}}{=} 0$

6)  $\vec{v} \stackrel{\text{AS}}{=} \nabla \phi$  ( follows from 2 )

7)  $\frac{\partial \nabla \phi}{\partial t} = \nabla \frac{\partial \phi}{\partial t}$

this becomes:

$$\nabla \frac{\partial \phi}{\partial t} + \nabla \left( \frac{1}{2} \nabla \phi \nabla \phi \right) - \frac{\nabla \rho}{\rho} + \nabla \Omega = 0 \quad ( 2.4 )$$

using 8)  $\int dr_3$  "zooming in on the surface"

9)  $\rho \stackrel{\text{AS}}{=} 0$

this becomes:

$$\frac{\partial \phi}{\partial t} + \frac{1}{2} \nabla \phi \nabla \phi + \nabla \Omega = f(t) \quad ( 2.5 )$$

using 10)  $f(t) \stackrel{\text{AS}}{=} 0$

11)  $\Omega = gz$  only gravity is considered for the potential force

12)  $z =: \eta$  :at the surface

this becomes:

$$\frac{\partial \phi}{\partial t} + \frac{1}{2} \nabla \phi \nabla \phi + g\eta = 0 \quad ( 2.6 )$$

:at the surface

" the Bernoulli equation "



Linearising the equation by disregarding the higher order variations of :

$$\frac{\partial \phi}{\partial t} + g\eta = 0 \quad (2.7)$$

at the average surface level.

$\eta$  describes the position of the surface ( $\eta(x, y, t)$ ) but is not known explicitly.

A description may be determined by assuming that a particle at the surface moves with the surface:

$$\frac{dz}{dt} = \frac{\partial \eta}{\partial t} + \frac{\partial \eta}{\partial x} \cdot \frac{dx}{dt} + \frac{\partial \eta}{\partial y} \cdot \frac{dy}{dt} \quad (2.8)$$

Disregarding the higher order ( non-linear ) variations:

$$\frac{\partial \eta}{\partial x} \cdot \frac{dx}{dt} \quad \text{and} \quad \frac{d\eta}{dy} \cdot \frac{dy}{dt}, \quad \text{and}$$

$$\text{using} \quad \frac{dz}{dt} = \frac{\partial \phi}{\partial z}$$

this becomes:

$$\frac{\partial \eta}{\partial t} - \frac{\partial \phi}{\partial z} = 0 \quad (2.9)$$

Combining the equation of the position of the surface ( 2.9 ) with the linearised Bernoulli equation ( 2.7 ):

$$(2.7): \quad \frac{\partial \phi}{\partial t} + g\eta = 0 \quad \Rightarrow \quad \frac{\partial^2 \phi}{\partial t^2} + g \frac{\partial \eta}{\partial t} = 0$$

$$(2.9): \quad \frac{\partial \eta}{\partial t} - \frac{\partial \phi}{\partial z} = 0$$

this produces:

$$\frac{\partial \phi}{\partial z} + \frac{1}{g} \cdot \frac{\partial^2 \phi}{\partial t^2} = 0$$

at the surface ( 2.10 )



This is the free surface boundary condition:

$$\frac{\partial \Phi}{\partial z} + \frac{1}{g} \frac{\partial^2 \Phi}{\partial t^2} = 0$$





Derivation of the refraction - diffraction equation

The equations for the fluid field are ( 2.04.1 ):

$$\begin{aligned} \frac{\partial \phi}{\partial z} + \frac{1}{g} \frac{\partial^2 \phi}{\partial t^2} &= 0 \\ \nabla^2 \phi &= 0 \\ \frac{\partial \phi}{\partial n} &= 0 \end{aligned}$$

D: mean water depth  
 $\sigma$ : bottom slope  
 $\lambda_0$ : deepwater wave-length

( 3.1 )

A solution exists for the case of a flat horizontal bottom ( 2.05.1 ):

$$\phi = \frac{ga}{\omega} \cdot \frac{\cosh(k(h+z))}{\cosh(kh)} \cdot e^{i(\omega t - kx)}$$

( 3.2 )

This solution can be considered to be the solution at the surface (  $\phi = \frac{ga}{\omega} \cdot e^{i(\omega t - kx)}$  )

and a function describing the distribution of this solution over the vertical:  $\frac{\cosh(k(h+z))}{\cosh(kh)}$ .

Berkhoff (1976) assumed that for a 'relatively flat' bottom this solution would still be approximately valid:

He assumed a solution of the form:

$$\phi(x, y, z, t) = a \cdot \frac{\cosh(k(h+z))}{\cosh(kh)} \cdot \phi(x, y, z) \cdot e^{i(\omega t - kx)}$$

( 3.3 )



The potential  $\varphi\langle x, y, z \rangle$  will be a very weak function of  $z$ , so it is assumed, most of the variation in the direction of the vertical already being gesccribed by the term  $\frac{\cosh(k(h+z))}{\cosh(kh)}$ .

The function  $\varphi\langle x, y, z \rangle$  may be expanded in a power series:

$$\varphi\langle x, y, z \rangle = \varphi_0\langle x, y \rangle + \sigma^2 z^2 \varphi_1\langle x, y \rangle + \sigma^4 z^4 \varphi_2\langle x, y \rangle + \dots \quad ( 3.4 )$$

At the surface:  $z=0$  and the power series will provide the original potential:  $\varphi\langle x, y, z=0 \rangle \equiv \varphi_0\langle x, y \rangle$

At the bottom  $z$  has its maximal value:  $z=D$  and neglection of the higher order terms will introduce an error of order of magnititude:  $\sigma^2 D^2$ .

Therefore, if  $\sigma^2 D^2 \ll 1$ , then:  $\varphi\langle x, y, z \rangle \approx \varphi_0\langle x, y \rangle$

Rescaling the equation, using "  $D$  " as a unit of distance along the bottom ( $\sigma^* := \frac{\sigma}{D}$ ) the criterion becomes:

$$\sigma^{*2} \ll 1 \quad ( 3.5 )$$

Using the free surface characteristic length to make the parameter dimensionless:

$$\varepsilon := \frac{\sigma}{D} \cdot \lambda_0 \quad ( 3.6 )$$

Using this the equations of the rescaled set become:

( see next page )



$$\frac{\partial \phi}{\partial z} + \frac{1}{g} \frac{\partial^2 \phi}{\partial t^2} = 0$$


$$\nabla^2 \phi = 0$$

$$\epsilon \left( \frac{\partial h}{\partial x} \cdot \frac{\partial \phi}{\partial x} + \frac{\partial h}{\partial y} \cdot \frac{\partial \phi}{\partial y} \right) + \frac{\partial \phi}{\partial z} = 0$$

( 3.7 )

using 1)  $\phi = a \cdot \frac{\cosh(k(h+z))}{\cosh(kh)} \cdot \phi_0(x, y) \cdot e^{i(\omega t - kx)}$

- 2) averaging for the mean energy flux by integration over the vertical from bottom to surface

Berkhoff obtains:

$$\nabla \cdot c c_g \nabla \phi_0 + k^2 c c_g \phi_0 = 0$$
( 3.8 )

the " Refraction - Diffraction equation "

The basic assumption ( thus: the major possible restriction of the validity of the equation ) is:

$$\sigma^2 D^2 \ll 1$$

which in scaled and dimensionless form becomes ( 3.6 ):

$$\epsilon^2 \ll 1$$

where:

$$\epsilon = \frac{\sigma \lambda_0}{D}$$





Listing of the computer program

This appendix contains the complete listing of the computer program used for this study.

The program, written in Fortran, is based on the Finite Element Package AFEP.

Next to the program listing a brief description of the purpose and workings of its parts is presented.

For a more complete understanding of the program also consult chapters 3 and 4, the AFEP User Manual and texts on Fortran IV if required.



JOB INITIATION  
"SIDE VIEW MODEL"

reservations for array space as required by FORTRAN.  
common blocks required by AFEP for the element matrix and subroutine PLSTEN

storing the size of each array in the first position of the array, as required by AFEP.

The variable, parameters of the test.

Only "NRUNS" has to be changed in line 3650 as well if change would be made.

```

010 //MMARDEF0 JOB 511,U,TIME=(10,39),REGION=2048K,MSGLEVEL=(2,0)
020 /*JOBPARM LIMES=9
030 // EXEC FORTXC3,PARM.GO=LET,EP=MAIN,SIZE=MAX-12K*
040 //FORT.SYSIN DD *
050 C
060 I IMPLICIT REAL *R(A-H,O-Z)
070 DIMENSION COOR(3000),KELEM(15000),KROUND(3000),KBNDDPT(500),
080 INTMAT(6),AMAT(9000),USOL(2000),RHS0(2000),RHS1(2000),
090 IUSER(40),USFR(40),IUSER(40),
100 V IUSER(2000),WORK(5000),DEPTH(2000),CUNTLN(16)
110 V IUSER(2000),RINPUT(99),KORDER(300),CUNTLN(16)
120 V COMMON /CACTL/ ILEFM,ITYPE,ISOORT,IMPELM,ICUNTI,IFIRSI,NOI,ISTART
130 /CPL0TC/ SCALE,MCONT,MCEL,ITURN,IMARK,NUMD
140 C
150 COOR(1) = 3000
160 KELEM(1) = 15000
170 KROUND(1) = 3000
180 KBNDDPT(1) = 500
190 INTMAT(1) = 6
200 AMAT(1) = 9000
210 USOL(1) = 2000
220 RHS0(1) = 2000
230 RHS1(1) = 2000
240 WORK(1) = 500
250 DEPTH(1) = 200
260 IUSER(1) = 40
270 IINPUT(1) = 40
280 RINPUT(1) = 99
290 KORDER(1) = 300
300 CUNTLN(1) = 16
310 C
320 SCU = 40-1
330 OMEGA = 100
340 RATIO = 300
350 ITESI = 1
360 NRUNS = 10
370 C
380 CALL INTAKE(IINPUT,RINPUT,SCU,ITESI,RATIO)
390 DO 1 I=1,NRUNS
400 CALL VARY(I,RINPUT,SCU,ITEST)
410 CALL MESS2(I,KELEM,COORD,KBOUND,KBNDDPT,IWORK,WORK,
420 IINPUT,RINPUT,3000,000)
430 CALL DEEP(KBOUND,COR,DEPTH,BORDER,M1,OMEGA,RATIO)
440 IF (I.EQ.1)
450 VCALL COME(2,BELEM,KROUND,INTMAT)
460 VCALL AUSER(IUSER,DEPTH,KBNDDPT,COORD,OMEGA,M1)
470 VCALL MATRIX(-1,AMAT,COORD,KELEM,KBOUND,WORK,IUSER,INMAT,IWORK)
480 C
490 CALL BUSER(IUSER,USER)
500 CALL RESIDEC(-1,RHS0,COORD,KELEM,KROUND,WORK,IUSER,AMAT,USOL)
510 CALL SOLVFC(AMAT,USOL,RHS0,INMAT,KBOUND)
520 IF (I.EQ.1) WRITE(6,*) (USER(9+J-1),J=1,17)
530 CALL GRAPH(12,7,USOL,COORD,KBNDDPT)
540 IPRINT(USOL,KBNDDPT,CUNTLN,1,1,6,2,2)
550 ICONT = 2*NRUNS
560 CALL PLOT(C,KOUTP,COORD,KBNDDPT,USOL,3000,2600,
570 CUNTLN,11)
580 CALL IPRINT(USOL,KBNDDPT,CUNTLN,2,1,6,2,2)
590 CALL PLOT(C,KOUTP,COORD,KBNDDPT,USOL,3000,2600,
600 CUNTLN,11)
610 C
620 CONTINUE
630 *END OF THE MAIN PROGRAM*
640 STOP
650 END

```

1



```

0720 "TEXT: THE REQUIRED SUBROUTINES"
0725
0730 SUBROUTINE INTAKE(IINPUT,RINPUT,SCD,ITEST,RATIO)
0735
0740 IMPLICIT REAL*(A-H,O-Z)
0745 DIMENSION IINPUT(1),RINPUT(1)
0750 READ(55,*) (IINPUT(J+5),J=1,91)
0755 READ(55,*) (RINPUT(J+5),J=1,30)
0760
0765 GO TO (1,2,3),ITEST
0770
0775 C 1
0780 RINPUT(9) = SCD * RINPUT(9)
0785 RINPUT(11) = SCD * RINPUT(11)
0790 RINPUT(13) = SCD * RINPUT(13)
0795 RINPUT(16) = SCD * RINPUT(16)
0800 RINPUT(18) = SCD * RINPUT(18)
0805 RINPUT(21) = SCD * RINPUT(21)
0810
0815 C 2
0820 RINPUT(9) = SCD * RINPUT(9)
0825 RINPUT(11) = SCD * RINPUT(11)
0830 RINPUT(13) = SCD * RINPUT(13)
0835 RINPUT(16) = SCD * RINPUT(16)
0840 RINPUT(18) = SCD * RINPUT(18)
0845 RINPUT(21) = SCD * RINPUT(21)
0850
0855 C 3
0860 RINPUT(9) = SCD * RINPUT(9)
0865 RINPUT(11) = SCD * RINPUT(11)
0870 RINPUT(13) = SCD * RINPUT(13)
0875 RINPUT(16) = SCD * RINPUT(16)
0880 RINPUT(18) = SCD * RINPUT(18)
0885 RINPUT(21) = SCD * RINPUT(21)
0890
0895 RETURN
0900
0905 SUBROUTINE VARY(I,RINPUT,SCD,ITEST)
0910
0915 IMPLICIT REAL*(A-H,O-Z)
0920 DIMENSION RINPUT(1),TOE(4)
0925
0930 TOE(1) = 80-I
0935 TOE(2) = 40-I
0940 TOE(3) = 20-I
0945 TOE(4) = 10-I
0950 IOEDAM = TOE(I)
0955
0960 GO TO (1,2,3),ITEST
0965
0970 C 1
0975 IF (I.GE.11) STOP
0980 TUPDAM = SCD*(-100+(I-1)*(10-1))
0985 RINPUT(13) = TUPDAM
0990 RINPUT(16) = TUPDAM
0995
0998 C 2
1000 RETURN
1005 IF (I.GE.5) STOP
1010 RINPUT(10) = 2.2 - TOEDAM
1015 RINPUT(31) = RINPUT(10)
1020 RINPUT(17) = 2.8 + IOEDAM
1025 RINPUT(24) = RINPUT(17)
1030
1035 C 3
1040 RETURN
1045 IF (I.GE.5) STOP
1050 RINPUT(10) = 2.2 - TOEDAM
1055 RINPUT(31) = RINPUT(10)
1060
1065 RETURN
1070 END

```

Subroutine "INTAKE" and "VARY" depend on the choice of ITEST.

These subroutines change the data required for the making of the mesh by mesh generator MESD2D. These data are read from dataset "STORE(INDATA)" by means of FTSSTF0X1 (see line 3210).

Observe that in this case "INTAKE" is called only once, and that all variations are made at each run by "VARY". (compare versus remark at line 4170).

The form of the field, the number of nodal points per boundary, is described in "STORE(INDATA)": see line 3210



- ① raises the height of the trapezoidal dam
- ② displaces the toe of the symmetrical dam
- ③ displaces the toe of "halfdam".



SUBROUTINE DEEP(KBNDPT, COOR, DEPTH, BORDER, NI, OMEGA, RATIO)

IMPLICIT REAL \*8(A-H, O-Z)  
DIMENSION KBNDPT(1), COOR(1), DEPTH(1), BORDER(1)  
N8=KBNDPT(5)

MBN1=KBNDPT(1+MBN) - KBNDPT(0+MBN) + 1  
MBN2=KBNDPT(6+MBN) - KBNDPT(1+MBN) + 1  
NI=MBN2+5  
DEPTH(2) = MBN2  
DEPTH(3) = NI

DO 10 I=1, NI  
BORDER(I) = COOR(2+KBNDPT(I+MBN1-1))+4  
DEPTH(I) = -COOR(2+KBNDPT(I+MBN1-1))+5  
CONTINUE

IF REQUIRED, THE ARRAYS (DEPTH), (BORDER) MAY BE PRINTED BY  
WRITE(6,\*) DEPTH, BORDER  
WRITE(6,\*) KBNDPT  
WRITE(10,\*) NI, OMEGA, RATIO  
WRITE(10,\*) (DEPTH(I), I=1, NI)  
WRITE(10,\*) (BORDER(I), I=1, NI)

RETURN  
END

SUBROUTINE AUSER(IUSER, USER, DEPTH, KBNDPT, COOR, OMEGA, NI)

IMPLICIT REAL \*8(A-H, O-Z)  
DIMENSION IUSER(1), USER(1), DEPTH(1), KBNDPT(1), F(15), COOR(1)  
COMMON /CACIL/ ID(7), ISTART

DO 1 I=6, 40  
IUSER(I) = 0  
USER(I) = 0

CONTINUE  
GRAV = 100  
SURFAS = - (OMEGA\*\*2)/GRAV

NEWTON-RAPHSON IS USED TO DETERMINE LOCAL WAVE-LENGTH:

M = NI-5  
DO 2 I = 6, NI, M  
ITER = 0  
AK = 1.5  
ITER = ITER+1

IF (ITER.LE.20) GO TO 20  
WRITE(6,11) ITER  
FORMAT(10,1) ITER  
GO TO 22

DELTA = (OMEGA\*\*2) - GRAV\*AK\*DTANH(AK\*DEPTH(1))  
AK = AK+DELTA/(GRAV\*(DTANH(AK\*DEPTH(1))+AK\*DEPTH(1)/  
COSH(AK\*DEPTH(1))\*DCOSH(AK\*DEPTH(1))))  
IF (ABS(DELTA).GE.10-5) GO TO 10

IF (I.EQ.6) A2B1 = AK  
A2B3 = AK

DEEP :

For each point along the border the depth is computed using the coordinates of surface and bottom boundaries.

The information is made available for the Top View Model using background array "STORE(LINK)" see lines 3200, 4080, 6420.

AUSER

1) computation of the value of "k" at the boundary (see page 15-); lines 1740-1950

2) print the values of the parameters characteristic for the problem lines 1960-2020

3) computation of the distribution of the wave potential over the height as follows from the analytical flat bottom solution (see page -), which is used in the Barkhoff equation (page -) lines 2060-2180

4) Filling IUSER and USER, as described in a Appendix 5.  
First the arrays are cleared by filling them with zeroes (lines 1680-1710) Then the pointers and values are inscribed lines 2220-2390







2570 SUBROUTINE INPLOT(USOL,KBNPDT,CUNTLN,ICHOIS,NUM1,NUM2,M1,M2)  
2580 FOR THIS CASE: NUM1=1, NUM2=6, M1=M2=12. SEE SUBROUTINE CALL

2590 IMPLICIT REAL \*8(A-H,O-Z)  
2600 DIMENSION USOL(1),KBNPDT(1),CUNTLN(1)  
2610 COMMON/CPLOT/ SCALE,MCNT,MCCEL,MCCEL,ITURN,IMARK,NUMD  
2620  
2630 M=N=KBNPDT(5)  
2640 N=KBNPDT(NUM1+NRN) - KBNPDT(0+NRN) + 1  
2650 NCEL = M\*M  
2660 M=KBNPDT(NUM2+NRN) - KBNPDT(NUM1+NRN) + 1  
2670 NCEL = M\*M2  
2680 SCALE = 1D2

2690 AMAX = 3D2  
2700 AMIN = -AMAX  
2710 MUS = USOL(3)  
2720 LIMIT = ICHOIS\*5  
2730 DO 1 J=LIMIT,MUS,2  
2740 A = USOL(J)  
2750 IF (A.GT.AMAX) AMAX=A  
2760 IF (A.LT.AMIN) AMIN=A  
2770  
2780 RANGE = AMAX - AMIN  
2790 STEP = RANGE/1D1  
2800 DO 2 J=1,11  
2810 CUNTLN(J+5) = AMIN + (J-1)\*STEP  
2820 CONTINUE  
2830 WRITE(6,\*) (CUNTLN(J+5),J=1,11)  
2840 RETURN  
2850 END

2860 SUBROUTINE GRAPH(NUM1,NUM2,USOL,COOR,KBNPDT)

2870 IMPLICIT REAL \*8 (A-H,O-Z)  
2880 DIMENSION USOL(1),COOR(1),KBNPDT(1)  
2890 REAL \*4 F1(100),F2(100),ABSFFIE(100)  
2900 INTEGER \*4 IARR(100)  
2910 NBJ = KBNPDT(5)  
2920 NBOUN = KBNPDT(12+NRN) - KBNPDT(7+NRN) + 1  
2930  
2940 DO 1 I=1,NBOUN  
2950 LOCUS = 2\*KBNPDT(NBN-I)  
2960 AIJ1 = USOL(4+LOCUS)  
2970 AIJ2 = USOL(5+LOCUS)

2980 F1(I) = AIJ1  
2990 F2(I) = AIJ2  
3000 ABSFFIE(I) = DSQRT(AIJ1\*\*2 + AIJ2\*\*2)  
3010 CONTINUE  
3020 CALL PRFIE(F1,IARR,NBOUN,1,NBOUN)  
3030 CALL PRFIE(F2,IARR,NBOUN,1,NBOUN)  
3040 CALL PRFIE(ABSFFIE,IARR,NBOUN,1,NBOUN)

3050 RETURN  
3060 END

3070 //GO. SYSJOB DD DISP=SHR,DSN=ZITA.BKON  
3080 // DD DSN=SYS2.FURTLA,DISP=SHR  
3090 // DD DSN=SYS2.HUMLTAR,DISP=SHR  
3100 // DD DSN=SYS2.PLOTIAR,DISP=SHR  
3110 // DD DSN=ZITA.PI.FI2,DISP=SHR  
3120 // DD DSN=WRDLVT.PROGK,DISP=SHR  
3130 //GO.FI2I001 DD DSN=ZITA.PI.FI3,DISP=SHR  
3140 //GO.FI2F001 DD DSN=WRBDEO.STORE(LINK1),DISP=SHR  
3150 //GO.FI2F001 DD DSN=WRBDEO.STORE(LINKTA),DISP=SHR  
3160 //GO.FI2F001 DD DSN=WRBDEO.USOL1,DISP=SHR  
3170 //GO.FI2F001 DD DSN=WRBDEO.USOL1,DISP=SHR  
3180 //GO.PLOTUS DD DSN=WRBDEO.LP,DISP=SHR  
3190 //GO.FI2F001 DD DSN=WRBDEO.LC,DISP=SHR  
3200 //GO.SYSIN DD \*

3210 IN 3 5 2 0  
3220 011110 2 2 2  
3230 147  
3240 148  
3250 148  
3260 148  
3270 143

"INPLOT":

Uses the number of points along the boundaries to determine the number of cells for subroutine PLOTEN (see AFET Manual)

The threshold levels of the potential are chosen at intervals of 10% between the maximum and minimum values. The calculated values of these threshold levels are stored in array 'CONTLN', as required by 'PLOTEN'.

GRAPH:

The subroutine 'PRFIE', stored in library WWDLVT.PROG (see line 3180), is used to present the values of the potential as determined by 'solve'.

The numbers have to be changed from real \*8 to real \*4.

Printed are the values:  
F1 (real part of the potential)  
F2 (imaginary part)  
||F|| =  $\sqrt{F1^2 + F2^2}$

The Libraries required for the job as required for AFET and PRFIE.

Input and output datasets:

STORE (LINK1) : see lines 1570, 1080, 6420  
STORE (LINKTA) : 790, 800  
LP (required by MESS2D)  
LC (required by PLOTEN)

NUMERICAL DATA : see AFET MANUAL, SUBROUTINE INIT

The numbers 147, 148 indicate the elements ELM 147, ELM 148. ELM 148 see page 44



JOB INITIATION & "TOP VIEW" MODEL

as for Side View Model.

```

10 EXEC FORTXCG,PARV,GOU=DET,EP=MAIN,SIZE=MAX-12K
//FORT:SYSIBM DD *
IMPLICIT REAL *(A-H,O-Z)
DIMENSION COOR(1000),KELEM(5000),KBOUND(2000),KENDPT(300),
INMAT(6),AMAT(25000),USOL(1000),RHSD(1000),
WORK(1000),RINPT(200),DEPTH(200),RORDER(200),CONTLN(16),
IINPUT(100),RINPUT(300),IUSER(100),USER(1000),
V COMMON /CACTL/ IELEM,ITYPE,ISOORT,INPELM,ICOUNT,IFIRST,NOI,ISIART
V COMMON /CPLOT/ SCALE,NCONT,ACEL,MCEL,ITURN,IMARK,NUMD

```

```

COOR(1) = 1070
KFLG(1) = 5000
AMAT(1) = 25000
WORK(1) = 600
KBOUND(1) = 200
INMAT(1) = 1000
KENDPT(1) = 300
USOL(1) = 1000
RHSD(1) = 1000
DEPTH(1) = 200
RORDER(1) = 1000
USER(1) = 100
IINPUT(1) = 100
RINPUT(1) = 300

```

```

NRUNS = 10
DO 2 I=1,NRUNS

```

```

CALL INCOPY(BORDER,DEPTH,N,OMEGA,RATIO)
CALL INTAKE(I,KELEM,COOR,KBOUND,KENDPT,IWORK,WORK)
CALL RES520(I,IINPUT,RINPUT,300,ODC)

```

```

V IF (I.EQ.1)
V CALL LOCAL(2,KELEM,KBOUND,INTMAT)
V CALL LOCAL(DEPTH,USER,COOR,N,OMEGA)
V CALL AUSER(I,USER,COOR,OMEGA)
V CALL MATRIA(-1,AMAT,COOR,KELEM,KBOUND,WORK,IUSER,USER,INTMAT,IWORK)

```

```

V CALL RUSER(I,USER)
V CALL RHSD(-1,RHSD,COOR,KELEM,KBOUND,WORK,IUSER,USER,AMAT,USOL)
V CALL SOLVE(A,AMAT,USOL,RHSD,INMAT,KBOUND)
V CALL GRAPH(NUM1,NUM2,USOL,COOR,KENDPT)
V CALL IMPLT(USOL,KENDPT,CONTLN,1,1,2,2,2)
NCONT = 2*NRUNS
V CALL PLOT(I,ROUTP,COOR,KENDPT(5)-6,KENDPT,USOL,3000,2600)
V CALL IMPLT(USOL,KENDPT,CONTLN,2,1,2,2,2)
V CALL PLOT(I,ROUTP,COOR,KENDPT(5)-6,KENDPT,USOL,3000,2600)

```

```

2 CONTINUE
"END OF THE MAIN PROGRAM"
STOP
END

```

" THEN : THE REQUIRED SUBROUTINES "

SUBROUTINE INCOPY(BORDER,DEPTH,N,OMEGA,RATIO) THESE DATA ARE TAKEN FROM THE " SIDE MODEL " USING THE BACKGROUND MEMORY.

```

IMPLICIT REAL *(A-H,O-Z)
DIMENSION BORDER(1),DEPTH(1)
READ(10,*) N1,OMEGA,RATIO
READ(10,*) (DEPTH(I),I=1,N1)
READ(10,*) (BORDER(I),I=1,N1)

```

```

N=M1-5
RETURN
END

```

INCOPY : see lines 1560 - 1580

" N ": the number of nodes along the border (5 after positions are subtracted from the array length.)



SUBROUTINE "INTAKE": prepares the data required by MESB2D, using the coordinates, along the border of the side view Model and the Top View Model, which have been generated in the Side View Model.  
 For each variation of the side View Model the corresponding border is used.

"border": the coordinates of the nodes along the border.

LOCAL: compare versus USER in the side View Model.

- ① computation of the local wave-number along the border lines 4650 → 4880
- ② calculation of ccg (see page 15) lines 4890 → 4990
- ③ using the property  $\frac{\partial}{\partial x} = 0$ , the values of  $k$ , ccg are attributed to the nodal points having the same x-coordinate, thus obtaining the local values of the parameters  $c_{cg}$  and  $k$  throughout the field, as required by element ELM301 (see page -) lines 5000 → 5040
- ④ If the depth has not changed, no computation is required: the values in the prior point are used lines 5060 → 5110
- ⑤ The number of parameters, and the values of the parameters "AKX", "ccg" may be printed line 5130

SUBROUTINE INTAKE(IINPUT,RINPUT,BORDER,I,N,RATIO)  
 THE DATA FOR "IINPUT" & "RINPUT" ARE FETCHED FROM THE BACKGROUND MEMORY.

4170  
4180  
4190  
4200  
4210  
4220  
4230  
4240  
4250  
4260  
4270  
4280  
4290  
4300  
4310  
4320  
4330  
4340  
4350  
4360  
4370  
4380  
4390  
4400  
4410  
4420  
4430  
4440  
4450  
4460  
4470  
4480  
4490  
4500  
4510  
4520  
4530  
4540  
4550  
4560  
4570  
4580  
4590  
4600  
4610  
4620  
4630  
4640  
4650  
4660  
4670  
4680  
4690  
4700  
4710  
4720  
4730  
4740  
4750  
4760  
4770  
4780  
4790  
4800  
4810  
4820  
4830  
4840  
4850  
4860  
4870  
4880  
4890  
4900  
4910  
4920  
4930  
4940  
4950  
4960  
4970  
4980  
4990

```

IMPLICIT REAL *8(A-H,O-Z)
DIMENSION IINPUT(1),RINPUT(1),BORDER(1)
IF (I.EQ.1) (IINPUT(J+5),J=1,51)
VREAD(55,*) (IINPUT(J+5),J=1,51)
IF (I.EQ.1) (RINPUT(J+5),J=1,2)
VREAD(55,*) (RINPUT(J+5),J=1,2)
THE SIDE-BOUNDARY ("BORDER") IS READ:
    
```

```

DO 11 J=1,N
RINPUT(2*J+6)=RORDER(J+5)
RINPUT(2*J+7)= 0D0
CONTINUE
FOLLOWED BY READING THE TOP-BOUNDARY:
DO 12 J=1,N
K= 0
RINPUT(2*J+2*N+6)=BORDER(K)
RINPUT(2*J+2*N+7)= 1D0
CONTINUE
    
```

```

IF REQUIRED THE ARRAY (RINPUT) MAY BE PRINTED BY:
WRITE(6,*) (RINPUT(J),J=1,233)
RETURN
END
    
```

```

SUBROUTINE LOCAL(DEPTH,USER,COORD,N,OMEGA)
IMPLICIT REAL *8(A-H,O-Z)
DIMENSION DEPTH(1),USER(1),COORD(1),AKK(300),CCG(300)
APARMS = 300R(4)
CALL PRIMAL(LOCAL,USER,5+2*NPARMS,-1,USER,IDUMMY)
CALL PRIMAL(LOCAL,AKK,5+NPARMS,-1,AKK,IDUMMY)
THIS CALL CHECKS THE SIZE OF ARRAYS "USER", "AKK".
WRITE(6,*) N,(DEPTH(J),J=1,N)
    
```

```

NBOUN = N
GRAV = 1D0
DO 111 I=1,NBOUN
I5= I+5
IF (DABS(DEPTH(I5))-DEPTH(I5-1)) LE 10-6) GO TO 15
NO ITERATION IS REQUIRED.
ELSE
NEWTON-RAPHSON IS USED TO DETERMINE LOCAL WAVE-LENGTH:
    
```

```

(I=I) AK = 0PPTH(I5) ; AK = AK*((I-1)*5)+1)
ITER = 0
ITER = ITER+1
IF (ITER LE 20) GO TO 12
WRITE(6,11) ITER
FORMAT(10THE ITERATION FAILS: ITER .GE. ,13,/)
GO TO 13
    
```

```

DELTA = (OMEGA**2) - GRAV*AK*DTANH(AK*H)
AK = AK+DELTA/(GRAV*(DTANH(AK*H)+(AK*H)/
(DCOSH(AK*H)**2)))
IF (DABS(DELTA).GE.1D-5) GO TO 1
    
```

```

THE LOCAL WAVE-NUMBER HAS BEEN DETERMINED: "AK".
LOCALLY THE FACTOR "CCG" IS:
CCG = GRAV*(5D-1)*(DSINH(AK*H)*DCOSH(AK*H)+(AK*H)/
(AK*(DCOSH(AK*H)**2)))
AN ALTERNATIVE:
CU = DSQRT(GRAV*DTANH(AK*H)/AK)
CGU = CU*(5D-1)*(1D0+2D0*AK*H/DSINH(2D0*AK*H))
IF (I LE 10) WRITE(6,*) H,AK,CCG,OMEGA,N
    
```



```

IF (I.LE.10) WRITE(6,*) H,AK,CCGL,OMEGA,N
THE CONDITIONS ALONG THE LENGTH OF THE FIELD ARE ATTRIBUTED,
PER STEP, TO THE WIDTH OF THE FIELD.
DO 14 J=1,5

```

```

13 GO TO 111
14 FOR CONSTANT DEPTH: LOCAL VALUES: SAME AS IN PRIOR POINT.
15 DO 16 J=1,5
16 AKK(5*I+J) = AK
   CCG(5*I+J) = CCGI

```

```

17 AKK(5*I+J) = AKK(5*(I-1)+J)
   CCG(5*I+J) = CCG(5*(I-1)+J)

```

```

111 CONTINUE

```

```

WRITE(6,*) NPARMS,AKK,CCG

```

```

DO 222 I=1,NPARMS
  I5=I+5
  USER(I5) = CCG(I5)
  USER(I5+NPARMS) = CCG(I5)*(AKK(I5)**2)
222 CONTINUE
  USER(5+2*NPARMS+1) = AKK(6)*2
  USER(5+2*NPARMS+2) = AKK(6)
  USER(5+2*NPARMS+3) = AKK(NBOUN*5+6)

```

```

PI=4*DATAV(107)
TPER = (200*PI)/OMEGA
WAVE = GRAV*(TPER**2)/(200*PI)
WAVE = 2*PI/AKK(6)

```

```

999 WRITE(6,999) TPER,WAVE,WAVE,AKK(6)
FORFAT(, WAVE PERIOD WAVELENGTH = ,F12.6, //
        , DEFN WAVE WAVELENGTH = ,F12.6, //
        , EXISTING WAVELENGTH = ,F12.6, //
        , EXISTING WAVELENGTH K = ,F12.6, //)
V
V
RETURN
END

```

```

SUBROUTINE AUSER(IUSER,USER,COOR,OMEGA)

```

```

IMPLICIT REAL *8(A-H,O-Z)
DIMENSION IUSER(1),USER(1),COOR(1)
COMMON /CACL/ ID(7),ISTART
DO 1 I=1,40
  IUSER(I) = 0
  CONTINUE
  IUSER(I) = 0
  NPARTS = NPARTS(4)
  ISTART = IPOINT
  IUSER(6) = 1
  IUSER(7) = 1
  IUSER(8) = 1
  IUSER(9) = 1
  IUSER(10) = 1
  IUSER(11) = 16
  IUSER(12) = 20
  IUSER(13) = 24
  IUSER(14) = 28
  IUSER(15) = 32
  IUSER(15) = 2*NPARMS
  IUSER(21) = -(LAST+1)
  IUSER(23) = -(LAST+2)
  IUSER(31) = -(LAST+3)
  WRITE(6,*) IUSER
  RETURN
END

```

```

SUBROUTINE RUSER(IUSER,USER)

```

```

IMPLICIT REAL *8(A-H,O-Z)
DIMENSION IUSER(1),USER(1)
COMMON /CACL/ ID(7),ISTART
IUSER(6) = 1
IUSER(7) = 1
IUSER(8) = 0
IUSER(9) = 0
IUSER(10) = 0
RETURN
END

```

⑥ The values of  $w_9$  and of

" $k^2 c c g$ " are stored in array "USER", as well as the parameters for the boundary conditions.  $2k, k, k$  as required for ELM301 resp ELM302 and ELR302 (See appendix 5) lines 516 → 5230

⑦ Some characteristic information for the problem is computed and printed lines 5250 → 5340

"AUSER"

See lines 2200 → 2410 and appendix 5

"RUSER"

see lines 2450 → 2560 and appendix 5







Input for subroutines IUSER and USER

The input for the boundary conditions ( the values of the parameters  $\dot{a}$ ,  $\dot{g}$  for  $\frac{d\phi}{dn} = -\dot{a}\phi + \dot{g}$  ) is provided using arrays " USER " and " IUSER " ( AFEP convention ).

IUSER: array containing " integer " numbers.

USER:           ,,           " real "           ,,

In order to save memory space these arrays operate in the following way:

- the position in array IUSER corresponds with a value of a boundary condition of a specific boundary.
- a shift in the positions in the arrays may be accounted for by using shift parameter ISTART ( which is also useful when other values are also stored in IUSER ).
- The number ( " # " ) in a position of array IUSER indicates the position(s) in array USER where the required value(s) is stored belonging to the specific parameter of the considered boundary. ( That is: # acts as a "pointer" )

The following convention is used:

- # < 0 : The value of the parameter is constant along the boundary. The value is stored in position "# " of array USER, and is to be used in all nodes belonging to the boundary.
- # = 0 : The value of the parameter is zero. ( No values are stored in array USER. )



#>0 : The value of the parameter varies along the boundary. The value for each consecutive node is stored in consecutive positions of array IUSER, beginning at position "# " for the first node.

Apart from this information concerning the values of the parameters for the boundary conditions, array IUSER contains some administrative information as well concerning the structure of array IUSER itself.

Apart from that, it also contains a set of "switches" which indicate if a boundary has only zero-values for its parameters ("0") or not ("1").

If all the values of the parameters are zero, the switch is "off" ("0"), the contributions of the boundary element matrices and vectors do not have to be computed for that boundary ( due to the natural boundary conditions of the finite element method ).

The administrative section of IUSER describes the "lay-out": the number of boundaries involved; the number of parameters per boundary; the position of the first "pointer" per boundary.

It will be shown how this is done using the boundary conditions of the Side View Model as an example.

The parameter ( a1, a2, g1, g2 ) of boundary "# " is denoted by: a1b#, a2b#, ...



ISTART = 10

A shift of 10 positions is indicated.

Apart from the standard 5 positions reserved for AFEP,

10 positions are used for the "administrative section":

- 5 positions containing the "switches", one for the field and four for the boundaries, indicating whether the element matrix is to be computed or not ( "0" ).
- 5 positions with pointers for the first positions occupied by the parameters.

<u>IUSER</u> :	<table border="1" style="border-collapse: collapse; text-align: center; width: 100px; height: 100px;"> <tr><td>0</td><td>1</td><td>2</td><td>3</td><td>4</td><td>5</td></tr> <tr><td>6</td><td>1</td><td>1</td><td>0</td><td>1</td><td>1</td></tr> <tr><td>11</td><td>16</td><td>20</td><td>24</td><td>28</td><td>32</td></tr> <tr><td>16</td><td>0</td><td>0</td><td>0</td><td>0</td><td></td></tr> <tr><td>21</td><td>0</td><td>9</td><td>0</td><td>-6</td><td></td></tr> <tr><td>24</td><td>0</td><td>0</td><td>0</td><td>0</td><td></td></tr> <tr><td>28</td><td>0</td><td>0</td><td>0</td><td>-7</td><td></td></tr> <tr><td>32</td><td>0</td><td>0</td><td>-8</td><td>0</td><td></td></tr> </table>	0	1	2	3	4	5	6	1	1	0	1	1	11	16	20	24	28	32	16	0	0	0	0		21	0	9	0	-6		24	0	0	0	0		28	0	0	0	-7		32	0	0	-8	0		<p>5 AFEP positions</p> <p>" switches "</p> <p>pointers per boundary</p> <p>pointers to the positions in array USER, per parameter, per boundary.</p> <p>Compare versus</p> <p>( 2.10.1 ),</p> <p>( 2.10.2 ),</p> <p>( 3.07.1 ) and</p> <p>( 3.07.2 ).</p>
0	1	2	3	4	5																																													
6	1	1	0	1	1																																													
11	16	20	24	28	32																																													
16	0	0	0	0																																														
21	0	9	0	-6																																														
24	0	0	0	0																																														
28	0	0	0	-7																																														
32	0	0	-8	0																																														
<p>field:</p> <p>boundary 1</p> <p>boundary 2</p> <p>boundary 3</p> <p>boundary 4</p>	<p style="text-align: center;">g1 g2 a1 a2</p>																																																	

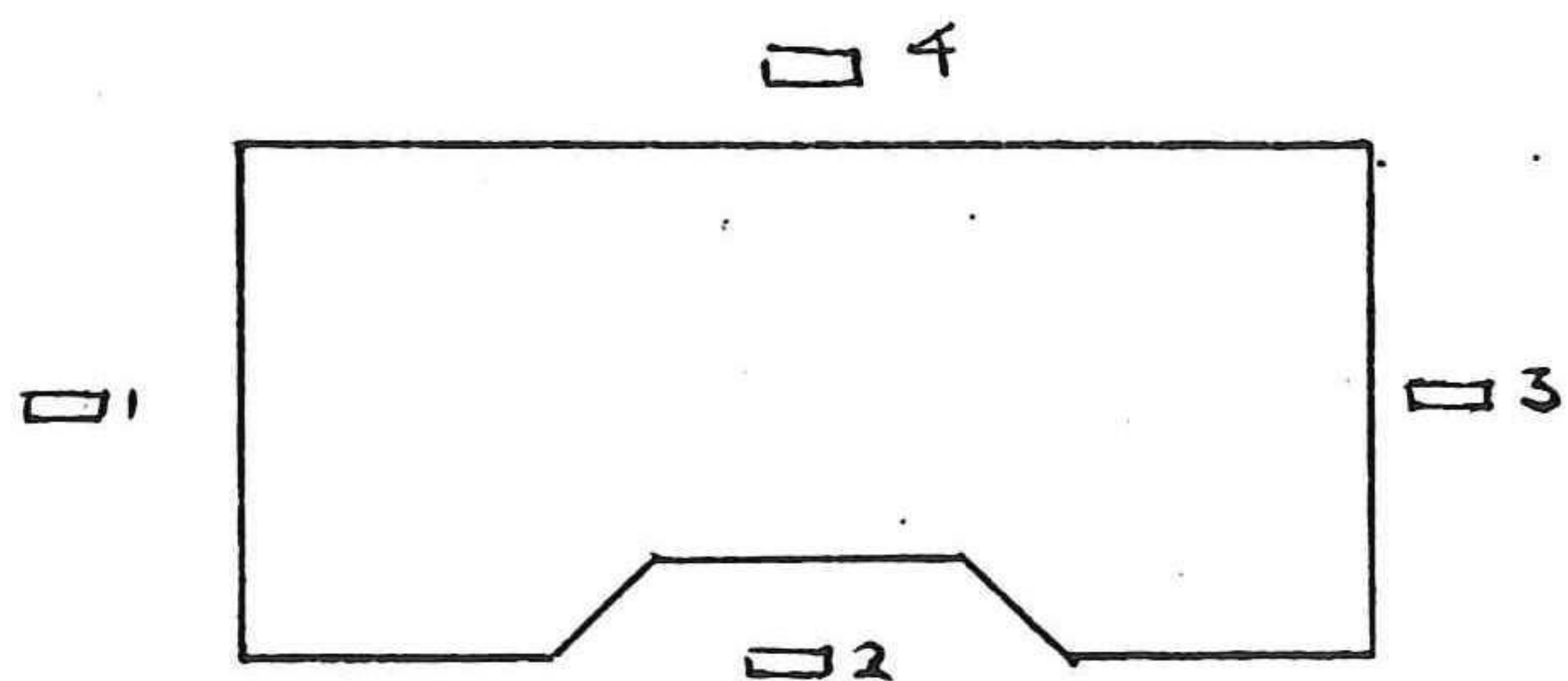
( AFEP requires a description of the parameter for the field condition. See the AFEP User Manual. )



USER :

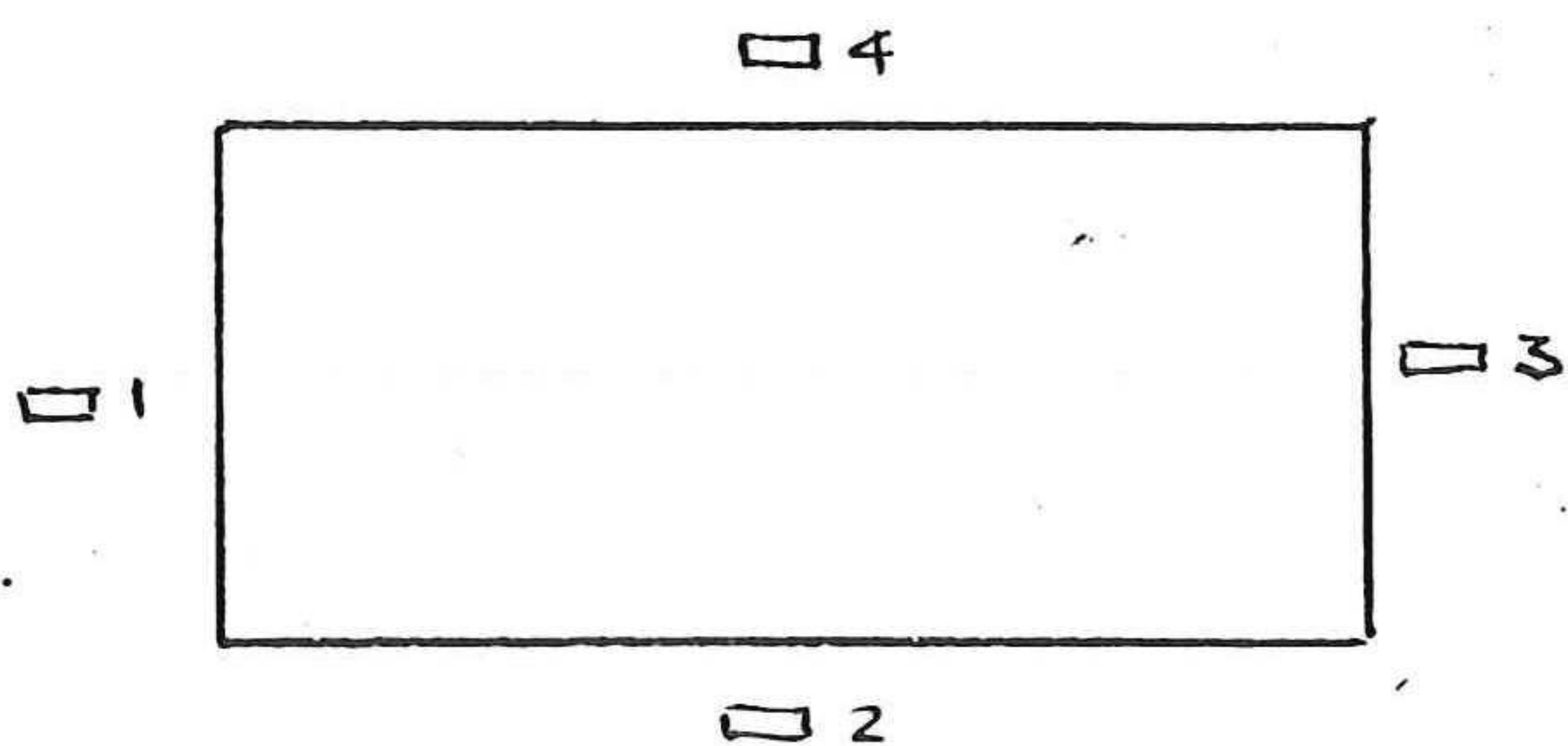
position:            1    }  
                       2    } AFEP reserved  
                       3    }  
                       4    }  
                       5    }  
                       6    a2b1 = k  
                       7    a2b3 = k  
                       8    a1b4 =  $\omega^2$  + grav  
                       9    g2b1 for node 1  
                     10                    2  
                       :                    3  
                       :                    .  
                       :                    .

Presented in another form, for both the Side View Model and the Top View Model:



( See ( 2.10.1 ) )

g1	g2	a1	a2	
0	0	0	0	field
$\omega \cdot k$	0	0	k	□ 1
0	0	0	0	□ 2
0	0	0	k	□ 3
0	0	$\frac{\omega^2}{g}$	0	□ 4



( See ( 2.10.2 ) )

g1	g2	a1	a2	
0	0	0	0	field
$\omega$	0	0	k	□ 1
0	0	0	0	□ 2
0	0	0	k	□ 3
0	0	0	0	□ 4

↙



Numerical aspects

In this appendix four facets will be dealt with:

- 1) verification of the location of the influx boundary,
- 2) verification of the steplength of the numerical computation,
- 3) ( corollary of (2) ): verification of the predictions of the Top View Model for some cases,
- 4) determination of a criterion for the significance of the difference in the predictions of the two models.

1) Position of the influx boundary

Only for the Side View Model is it required to check the location of the influx boundary, which has to be sufficiently far removed from the barrier. ( This follows from the properties of the model, allowing a redistribution of the potential along the vertical ).

CASE A has been computed to check the effect of the variation of the length of the wavefield upwave of the dam.

The only difference between CASE A and CASE 3 is the value of  $S_{CL}$  ( See figure 6.03.3 ) ).

The length of the wavefield is one unit longer, upwave of the dam, in CASE A. This amounts to an increase of approximately 50% compared to CASE 3

*50% longer?*



The results of the two cases are:

Table A6.1.1 effects of varied steplength	
	$\  \varphi_t \ $
CASE A	( 0.939 $\pm$ 0.002 )
CASE 3	( 0.941 $\pm$ 0.002 )

The difference between the two predictions is not significant.

The influx boundary is indeed sufficiently far removed from the dam.

## 2) Steplength

The computational accuracy depends upon the size of the steplength; the computational costs also: accuracy improves with decrease of steplength; costs go down with increase of steplength.

The steplength varies throughout the field of computation, the mesh being refined where the variations of the potential will be greater: over the dam. See figure ( 7.01.2 ).

The size of the steplength is verified by comparing CASE B versus CASE 3 for the Side View Model, and by comparing CASE 3 versus computations done using a one dimensional model bases on the refraction - diffraction equation, made by N. Booy. ( These computations also provide a possibility to verify the results of the Top View Model. This will be



discussed in paragraph (3) ).

Because of requirements from the program concerning memory space in the program, the comparison of CASE 3 versus CASE B has been done using a double steplength for CASE B. The numerical error may be determined as follows.

For a linear model the relation between the errors in the results and the steplength may be described by:

$$E_1 - E_2 = C \cdot (\Delta x_1^2 - \Delta x_2^2) \quad (A6.2.1)$$

gives:  $\alpha = C \cdot \Delta x^2 \quad (A6.2.2)$

where:  $E_i$  : result of computation with steplength  $x_i$

$\alpha$  : error in the result.

C : coefficient:

( See Kan & Segal, 1979 )

In this case:  $E_1 = .941$  ( CASE 3, 6<sup>th</sup> item, table ( 7.04.1 ) ) )

$E_2 = .934$  ( CASE 9 )

$\Delta x_1 = .1$

$\Delta x_2 = .21$

Substitution in ( 7.03.1 ) gives:

$$C = .2025$$

Substitution in ( 7.03.2 ) gives:

$$\alpha = (0.2025) \cdot (0.1)^2 = 0.002$$



The calculation shows that the steplength of the Side View Model has a suitable size.

The stepsize in the Top View Model is verified by computation done using the one dimensional model made by N. Booy, based on the refraction - diffraction equation.

These computations have been done using the same number of nodes as well as with the double number of nodes ( thus with half the steplength ) as has been used in CASE 3.

The results are identical for these three cases:

Table A6.2.3 Presentation of the results of the comparison of the results of the Top View Model versus the model by N. Booy with variation of steplength.		
Top View Model CASE 3	Booy model same steplength	Booy model half steplength
$\  \varphi_t \ $ : . 941	.941	.941

The comparison shows that the steplength of the Top View Model has a suitable size.

### 3 Verification of the Top View Model

It has been possible to check the values of  $\| \varphi_t \|$ , as predicted by the Top View Model, by using the model made by N. Booy to compute  $\| \varphi_t \|$  for identical configurations. The results were identical.

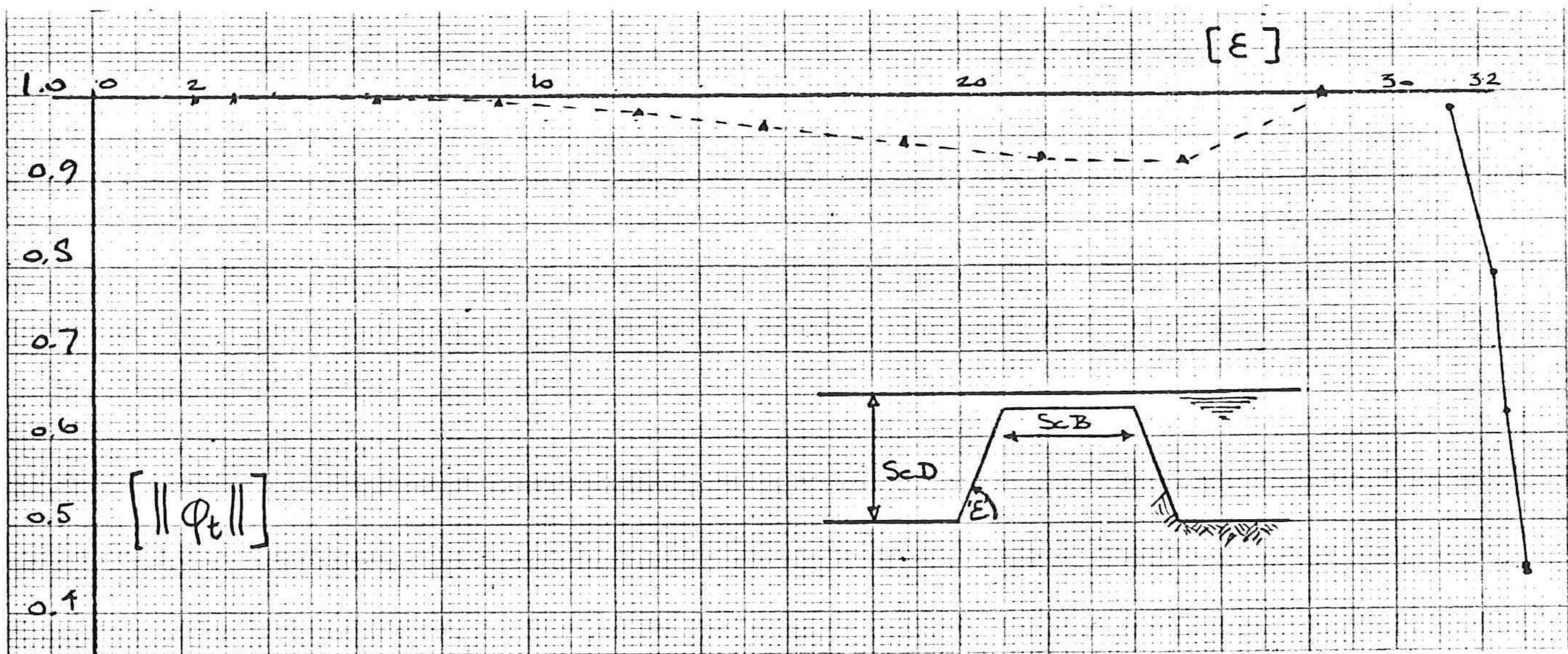


The model by Booy has also been used to provide some values of  $\|\phi_t\|$  for those combinations of parameters which lead to doubtful results ( See section 8.02 ).

The results are presented in the following table.

Table A6.3.1								
Values of $\ \phi_t\ $ computed using the model made by N. Booy. Compare versus diagram 6.07.1 and table 7.04.1, CASE 2.								
#	ScD	Ratio	ScB	$\Delta D$	$\Delta x$	sigma	epsilon	$\ \phi_t\ $
1	0.6	1:12	0.6	.55	0.2	2.75	31.27	.980
2		1:20		.570		2.85	32.30	.788
3		1:24		.575		2.875	32.59	.624
4		1:40		.585		2.925	33.25	.447
5		1:60		.590		2.95	33.54	.440
6		1:3		.40	2.0	0.2	2.27	.999

These results are also presented in the following graph.



Graph A6.3.2 Presentation of the results of table A6.3.1

The results of the computations with the model made by N. Booy confirm the predictions made using the Top View Model, even in the behaviour which does not conform with the expectations, as described in section 8.02, for low values of Ratio. Further investigation will be required to explain these results.



4) Accuracy of the values of  $\|\varphi\|$

As has been described in section ( 6.04 ),  $\|\varphi_t\|$  has been selected as characteristic parameter of the comparative calculation.

In this paragraph the error in determining  $\|\varphi_t\|$  will be determined. ( For comparison this will also be done for  $\|\varphi_c\|$  and  $\|\varphi_r\|$  ; it will be seen that the uncertainty in these parameters is greater. )

The results of the computations for this study are presented in the form depicted in figure ( 7.02.3 ).

The roundoff error in the results is, for all cases, one unit of the fifth significant digit.

The uncertainty in  $\|\varphi_t\|$

The uncertainty in  $\|\varphi_t\|$  can be determined directly from the graphs of  $\|\varphi\|$  which form the output of the computer program.

There is a slight fluctuation, a "fuzziness", in the values of  $\|\varphi_t\|$  ( See e.g. figure ( 7.02.3 ) ).

This fuzziness is a measure of the uncertainty of the values of  $\|\varphi_t\|$ .

For all cases the fuzziness has been determined from the graphs. The results are displayed in Table ( A6.4.1 ).



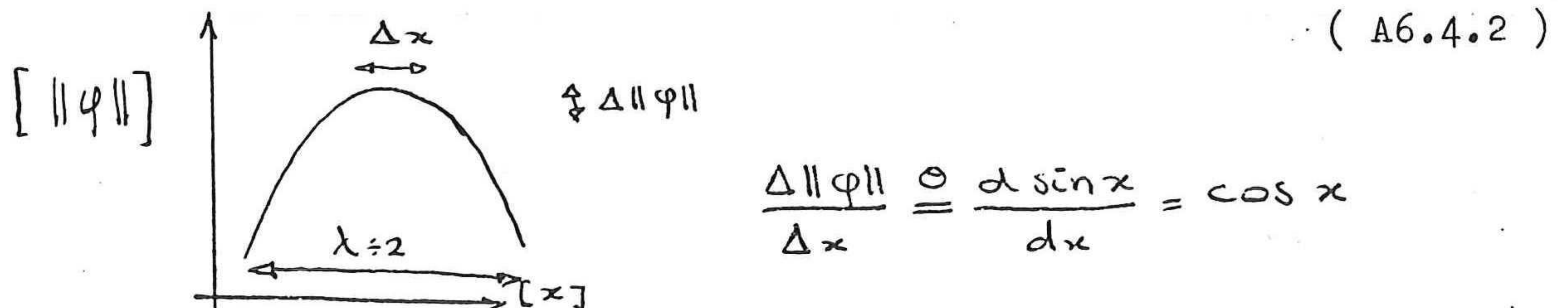
Table A6.4.1												
Order of magnitude of the fuzziness of per computed case. ( multiplied times $10^3$ ).												
CASE	1	2	3	4	5	6	7	8	9	10	11	12
Side View M.	.4	.5	.6	.2	.2	.4	.5	.4	.7	.5	.5	1.
Top View M.	.3	.5	.5	.4	.6	.4	.4	.4	1.	.4	.5	2.

The uncertainty in  $\|\varphi_i\|$  and  $\|\varphi_r\|$

The values of  $\|\varphi_i\|$  and  $\|\varphi_r\|$  have to be determined from the maximum and minimum values of  $\|\varphi\|$ , as described in section ( 6.05 ).

Because the locations of the maxima and minima will not necessarily coincide with the location of the nodes of the computation, an error will be introduced in the value of the maxima and minima.

An estimate of the size of this error may be made using the following sketch:



With this information the uncertainty in the value of at the location of the maxima and minima may be determined.



The steplength  $\Delta x$  can be expressed in the wavelength:

$$\Delta x \stackrel{\circ}{=} 0.11 \quad ; \quad \lambda \stackrel{\circ}{=} 3.71 \quad \Rightarrow \quad \Delta x \stackrel{\circ}{=} 0.030 \lambda$$

The uncertainty is:

$$\Delta \|\varphi\| = \cos 0 - \cos \frac{\Delta x}{\lambda} = 2 \cdot 10^{-3}$$

Error calculus shows that the error in the values of  $\|\varphi_i\|$  and  $\|\varphi_r\|$  will be described by:

$$\Delta \|\varphi_{i,r}\| = \frac{\Delta \|\hat{\varphi}\|}{\|\hat{\varphi}\|} + \frac{\Delta \|\check{\varphi}\|}{\|\check{\varphi}\|} \stackrel{\circ}{=} \frac{2 \cdot 10^{-3}}{1.5} + \frac{2 \cdot 10^{-3}}{0.5} \stackrel{\circ}{=} 5 \cdot 10^{-3}$$

( Using the order of magnitude of  $\|\hat{\varphi}\|$  and  $\|\check{\varphi}\|$  . )

It is clear that the uncertainty in  $\|\varphi_t\|$  is smaller than in  $\|\varphi_i\|$  and  $\|\varphi_r\|$  .

$$\begin{aligned} \text{The uncertainty in } \|\varphi_t\| &= 2 \cdot 10^{-3} \quad (\text{Table A6.4.1}) \\ \text{error } \|\varphi_t\| &= 2 \cdot 10^{-3} \quad (\text{Paragraph 2}) \end{aligned}$$

This leads to the criterion to be used for the uncertainty of  $\|\varphi_t\|$  :

$$\Delta \|\varphi_t\| = 2 \cdot 10^{-3}.$$

As a consequence the error in the parameter of comparison will be:

$$\Delta \left( \frac{\|\varphi_t\|_r}{\|\varphi_t\|_s} - 1 \right) = \frac{\Delta \|\varphi_t\|_r}{\|\varphi_t\|_r} + \frac{\Delta \|\varphi_t\|_s}{\|\varphi_t\|_s} \stackrel{\circ}{=} 4 \cdot 10^{-3}$$





The listings of the elements of the element vectors and element matrices for both the Laplace equation and the refraction - diffraction equation are presented on the next pages.

A description of the derivation of the elements is presented in chapter 3.



Jun 13 04:04 1981 .elm147 Page 1

```
1      subroutine elm147(coor,work,iuser,user
2      implicit real *8 (a-h,o-z)
3      dimension coor(1),work(1),user(1),iuse(1)
4      common /cwork/ index1(100),index2(100)
5      v      /cact1/ ielem,itype,isoort
6      100 11=index1(1)
7         12=index1(2)
8         13=index1(3)
9         x1=coor(4+2*11)
10        x2=coor(4+2*12)
11        x3=coor(4+2*13)
12        y1=coor(5+2*11)
13        y2=coor(5+2*12)
14        y3=coor(5+2*13)
15        e1=y2-y3
16        e2=y3-y1
17        e3=y1-y2
18        d1=x3-x2
19        d2=x1-x3
20        d3=x2-x1
21        h=1d0/(2d0+dabs(e2*d3-d2*e3))
22        a11=(e1*e1+d1*d1)*h
23        a12=(e1*e2+d1*d2)*h
24        a13=(e1*e3+d1*d3)*h
25        a22=(e2*e2+d2*d2)*h
26        a23=(e2*e3+d2*d3)*h
27        a33=(e3*e3+d3*d3)*h
28        work(0)=a11
29        work(3)=a12
30        work(10)=a13
31        work(15)=a12
32        work(20)=a22
33        work(22)=a23
34        work(30)=a13
35        work(32)=a23
36        work(34)=a33
37        do 102 i=0,30,12
38        do 102 j=1,5,2
39        work(i+j+0)=work(i+j-1)
40        work(i+j)=0d0
41        work(i+j+5)=0d0
42      102 continue
43      return
44      end
```



Jun 13 04:04 1981 elm143 Page 1

```
1      subroutine elm143(coor,work,iuser,user)
2      implicit real *8 (a-h,o-z)
3      real * 8 naam
4      dimension coor(1),work(1),iuser(1),user(1)
5      dimension x(2,2)
6      common /cwork/ index1(100),index2(100)
7      v      /cact1/ ielem,itype,isoort,inpelm,icount,ifirst,not,istart
8      data naam /4helem/
9      c      *element matrices for the laplace equation (boundary conditions)
10     c      itype=105 linear conforming element (mixed type)
11     c      *local parameters
12     c      x      array containing the coordinates
13     c      ih      help variable
14     c      n1,n2   help variables to compute h
15     c      h      factor h (see description)
16     c      iusrt   position in array user (iuser(nstart))
17     c      a1,a2   factors a1,a2 (see description)
18     c      *computation of x and h
19     do 10 i=1,2
20         ih=2*index1(i)
21         x(1,1)=coor(ih+4)
22     10    x(1,2)=coor(ih+5)
23         n1=x(1,1)-x(2,1)
24         n2=x(1,2)-x(2,2)
25         h=sqrt(n1*n1+n2*n2)
26         if(h.ne.0) goto 30
27     c      *error 1: volume is 0
28         call ertrap(naam,1,15,itype,ielem,0)
29     30    if(ifirst.eq.1) goto 60
30     c      *ifirst=0, test inpelm and icount
31         if(inpelm.eq.2) goto 40
32     c      *error 2: inpelm#2
33         call ertrap(naam,2,20,itype,ielem,inpelm)
34     40    if(icount.eq.4) goto 50
35     c      *error 3: icount#4
36         call ertrap(naam,3,20,itype,ielem,icount)
37     50    nstart=iuser(isoort+istart)
38     c      test nstart
39         if(nstart.gt.5) goto 55
40     c      *error 4: |iuser(nstart)|<6
41         call ertrap(naam,4,15,itype,ielem,0)
42     55    iusrt1 = iuser(nstart+2)
43         iusrt2 = iuser(nstart+3)
44         if (iusrt1) 1,2,3
45     1    isrt1 = labs(iusrt1)
46         a11 = user(isrt1)
47         a12 = a11
48
49         n1 = isrt1
50     2    go to 4
51         a11 = 0d0
52         a12 = 0d0
53         n1 = 0
54     3    isrt1 = iusrt1
55         n1 = isrt1 + work(4)
56     4    if (iusrt2) 5,6,7
```



```

57      5  isrt2 = 1a0s(1usr2)
58          a21 = user(isrt2)
59          a22 = a21
60              n2 = isrt2
61      go to 3
62      6  a21 = 0d0
63          a22 = 0d0
64              n2 = 0
65      go to 8
66      7  isrt2 = 1usr2
67          n2 = isrt2 + work(4)
68      8  call prinal(naam,5h1user,nstart+3,6,dummy,iuser)
69          nu = n1
70          if(n2.gt.n1) nu = n2
71      call prinal(naam,4huser,nu,-7,user,dummy)
72      c      the values of a1 and a2 have been attributed
73      if (1usr1.ne.0.or.1usr2.ne.0) go to 60
74      not=1
75      return
76      c
77      60 if (1usr1.le.0) go to 11
78          a11 = user(isrt1)
79              isrt1 = isrt1 + 1
80          a12 = user(isrt1)
81      11 if (1usr2.le.0) go to 22
82          a21 = user(isrt2)
83              isrt2 = isrt2 + 1
84          a22 = user(isrt2)
85      c
86      22 fac = h/24
87          work(6) = fac*(5*a11+2*a12)
88          work(7) = -fac*(5*a21+2*a22)
89          work(8) = fac*(2*a11+2*a12)
90          work(9) = -fac*(2*a21+2*a22)
91          work(10) = fac*(5*a21+2*a22)
92          work(11) = fac*(5*a11+2*a12)
93          work(12) = fac*(2*a21+2*a22)
94          work(13) = fac*(2*a11+2*a12)
95          work(14) = fac*(2*a11+2*a12)
96          work(15) = -fac*(2*a21+2*a22)
97          work(16) = fac*(2*a11+5*a12)
98          work(17) = -fac*(2*a21+5*a22)
99          work(18) = fac*(2*a21+2*a22)
100         work(19) = fac*(2*a11+2*a12)
101         work(20) = fac*(2*a21+5*a22)
102         work(21) = fac*(2*a11+5*a12)
103      c
104      c
105      return
106      c      *error messages
107      c      1: volume of element = 0
108      c      2: inplm # 2
109      c      3: 1count # 2
110      c      4: |iuser(iscort+istart)| < 6
111      c      5: sigma <= 0
112      c      6: length of array iuser too small

```

```

113      c      7: length of array user too small
114      end

```



Jun 13 04:04 1951 elr148 Page 1

```
1      subroutine elr148(coor,work,iuser,user)
2      implicit real *8 (a-h,o-z)
3      real * 8 naam
4      dimension coor(1),work(1),iuser(1),user(1)
5      dimension x(2,2)
6      common /cwork/ index1(10),index2(10)
7      v      /cact1/ ielem,itype,isoort,inpelm,icount,ifirst,not,istart
8      data naam /5nelrns/
9      c      *element vectors for the laplace equation (boundarycontions)
10     c      itype=104 linear conforming element (neumann type)
11     c      itype=105 linear conforming element (mixed type)
12     c      *local parameters
13     c      x      array containing the coordinates
14     c      in      help variable
15     c      n1,h2   help variables to compute h
16     c      h      factor h (see description)
17     c      iusrt   position in array user (iuser(isoort+istart))
18     c      g      array containing the functionvalues (derivatives)
19     c      *computation of x and h
20     do 10 i=1,2
21         in=2*index1(i)
22         x(i,1)=coor(in+4)
23     10   x(i,2)=coor(in+5)
24         n1=x(1,1)-x(2,1)
25         n2=x(1,2)-x(2,2)
26         n=sqrt(n1*n1+n2*n2)
27         if(n.ne.0) goto 30
28     c      *error 1: volume is 0
29     c      call entrap(naam,1,15,itype,ielem,0)
30     30   if(ifirst.eq.1) goto 60
31     c      *ifirst=0, test inpelm and icount
32     call prinai(naam,5,iuser,isoort+istart,6,dummy,iuser)
33         if(inpelm.eq.2) goto 40
34     c      *error 2: inpelm#2
35         call entrap(naam,2,20,itype,ielem,inpelm)
36     40   if(icount.eq.4) goto 50
37     c      *error 3: icount#4
38         call entrap(naam,3,20,itype,ielem,icount)
39     50   nstart=iuser(isoort+istart)
40     c      test nstart
41         if(nstart.gt.5) goto 55
42     c      *error 4: |iuser(nstart)|<6
43     c      call entrap(naam,4,15,itype,ielem,0)
44     55   iusrt1 = iuser(nstart)
45         iusrt2 = iuser(nstart+1)
46         if (iusrt1) 1,2,3
47     1     isrt1 = iabs(iusrt1)
48         j11 = user(isrt1)
49         j12 = j11
50         n1 = isrt1
51         go to 4
52     2     j11 = 0j0
53         j12 = 0j0
54         n1 = 0
55         go to 4
56     3     isrt1 = iusrt1
```



Jun 13 04:04 1981 air143 Page 2

```
57          n1 = isrt1 + work(4)
58      4  if (iusrt2) 5,6,7
59      5  isrt2 = 1205(iusrt2)
60          g21 = user(isrt2)
61          j22 = j21
62          n2 = isrt2
63      go to 8
64      6  j21 = 0d0
65          j22 = 0d0
66          n2 = 0
67      go to 3
68      7  isrt2 = iusrt2
69          n2 = isrt2 + work(4)
70      8  call prin1(nam,5,iuser,nstart+3,0,dummy,iuser)
71          nu = n1
72          if(n2.gt.n1) nu = n2
73      call prin1(nam,4,iuser,nu,-7,user,dummy)
74      c          the values of g1 and g2 have been attributed
75      c          for the case g1,g2 constant
76      if (iusrt1.ne.0.or.iusrt2.ne.0) go to 60
77      not=1
78      return
79      c
80      60  if (iusrt1.le.0) go to 11
81          g11 = user(isrt1)
82          isrt1 = isrt1 + 1
83          j12 = user(isrt1)
84      11  if (iusrt2.le.0) go to 22
85          g21 = user(isrt2)
86          isrt2 = isrt2 + 1
87          j22 = user(isrt2)
88      c
89      22  h=h/0
90          work(6)=h*(2*g11+g12)
91          work(7)=h*(2*g21+g22)
92          work(8)=h*(g11+2*g12)
93          work(9)=h*(g21+2*g22)
94      return
95      c  *error messages
96      c  1:  volume of element = 0
97      c  2:  inpelm # 2
98      c  3:  icount # 2
99      c  4:  |iuser(iscount+istart)|<6
100     c  5:  sigma <= 0
101     c  6:  length of array iuser too small
102     c  7:  length of array user too small
103     end
```



Jun 13 04:03 1981 elm300 Page 1

```
1      subroutine elm300(coor,work,iuser,user)
2
3      c
4      implicit real *8(a-h,o-z)
5      dimension coor(1),work(1),iuser(1),user(1),
6              v      t(3,3),sn(6,6),ind2(6)
7      common /cwork/ index1(100),index2(100)
8              v      /nwork/ inde(3)
9              v      /cact1/ ielem,itype,isoort,inpelm,icount,ifirst
10     c
11     npoint = coor(4)
12     nparms= npoint
13     c *** computation of e and delta in r2
14     do 600 i=1,8
15     600     do 600 j=1,3
16             t(i,j)=0
17             inde(1)=index1(1)
18             inde(2)=index1(3)
19             inde(3)=index1(4)
20             call elm301(coor,work,iuser,user)
21             do 610 i=1,6
22     610             do 610 j=1,6
23                     sn(i,j)=work(j+6*i-1)
24                     ind2(1)=1
25                     ind2(2)=2
26                     ind2(3)=5
27                     ind2(4)=6
28                     ind2(5)=7
29                     ind2(6)=3
30             do 620 i=1,6
31     620             do 620 j=1,6
32                     i1=ind2(i)
33                     j1=ind2(j)
34                     t(i1,j1)=t(i1,j1)+sn(i,j)
35             inde(1)=index1(1)
36             inde(2)=index1(2)
37             inde(3)=index1(3)
38             call elm301(coor,work,iuser,user)
39             do 710 i=1,6
40     710             do 710 j=1,6
41                     sn(i,j)=work(j+6*i-1)
42                     ind2(1)=1
43                     ind2(2)=2
44                     ind2(3)=3
45                     ind2(4)=4
46                     ind2(5)=5
47                     ind2(6)=6
48             do 720 i=1,6
49     720             do 720 j=1,6
50                     i1=ind2(i)
51                     j1=ind2(j)
52                     t(i1,j1)=t(i1,j1)+sn(i,j)
53             inde(1)=index1(1)
54             inde(2)=index1(2)
55             inde(3)=index1(4)
56             call elm301(coor,work,iuser,user)
57             do 810 i=1,6
```



Jun 13 04:03 1981 elm300 Page 2

```
57      do 810 j=1,6
58      810      sh(i,j)=work(j+6*i-1)
59              ind2(1)=1
60              ind2(2)=2
61              ind2(3)=3
62              ind2(4)=4
63              ind2(5)=7
64              ind2(6)=3
65      do 820 i=1,6
66      do 820 j=1,6
67              i1=ind2(i)
68              j1=ind2(j)
69      820      t(i1,j1)=t(i1,j1)+sh(i,j)
70              inde(1)=index1(2)
71              inde(2)=index1(3)
72              inde(3)=index1(4)
73      call elm301(coor,work,iuser,user)
74      do 910 i=1,6
75      do 910 j=1,6
76      910      sh(i,j)=work(j+6*i-1)
77              ind2(1)=3
78              ind2(2)=4
79              ind2(3)=5
80              ind2(4)=6
81              ind2(5)=7
82              ind2(6)=8
83      do 920 i=1,6
84      do 920 j=1,6
85              i1=ind2(i)
86              j1=ind2(j)
87      920      t(i1,j1)=t(i1,j1)+sh(i,j)
88      do 940 i=1,6
89      do 940 j=1,6
90              work(j+6*i-3)=t(i,j)/2
91      940      continue
92      return
93      c *** error messages
94      c *** 1: volume of element equals zero
95      end
```



Jun 13 04:03 1981 elr300 Page 1

```
1      subroutine elr300(coor,work,iuser,user)
2
3      c
4      implicit real *8(a-h,o-z)
5      dimension coor(1),work(1),iuser(1),user(1),
6      v          x(3,2),b(2),f1(4)
7      common /cwork/ index1(100),index2(100)
8      v          /cact1/ ielem,itype,isoort
9      100 do 110 i=1,2
10         ih=2*index1(i)
11         x(i,1)=coor(ih+4)
12         x(i,2)=coor(ih+5)
13     110     b(i)= user(9+index1(i))
14     p=dsqrt((x(2,1)-x(1,1))*(x(2,1)-x(1,1))+
15     v        (x(2,2)-x(1,2))*(x(2,2)-x(1,2)))
16     fac= p/6
17     f1 = user(3)
18     f2 = user(9)
19     c          user(3) & user(9) contain "f1" resp "f2"
20     work(6) = fac* (2*b(1)+b(2))+f1
21     work(7) = fac* (2*b(1)+b(2))+f2
22     work(3) = fac* (b(1)+2*b(2))+f1
23     work(9) = fac* (b(1)+2*b(2))+f2
24     return
25     end
```



```

subroutine elm301(coor,work,user,user)
c
implicit real*8(a-n,o-z)
dimension coor(1),work(1),user(1),user(1),
v          e(3,2),x(3,2),ind(5),s(3,3),g(3),h(3)
common /cacti/ ielem,itype,isoort,inpelm,icount,ifirst,not,istart
v          /nwork/index1(3)
data ind/1,2,3,1,2/, njan/4helem/
nparms = coor(4)
icount=6
do 10 i=1,3
  ih=2*index1(i)
  x(1,1)=coor(ih+4)
  x(1,2)=coor(ih+5)
  g(i)= user(index1(i)+5)
  this section of "user" contains: c*cg.
c 10  h(i)= user(index1(i)+5+nparms)
  this section of "user" contains: c*cg*k*k.
c
do 20 i=1,3
  e(1,1)=x(ind(i+1),2)-x(ind(1+2),2)
  e(1,2)=x(ind(1+2),1)-x(ind(i+1),1)
20  delta=e(1,2)*e(3,1)-e(1,1)*e(3,2)
  if (delta) 22,21,25
c *** error 1: volume is 0
21  call ertrap(naam,1,15,itype,ielem,0)
  return.
22  delta=-delta
25  nm1=icount+1
  nm2=3*(icount-1)
  ik =icount*icount+5
  do 35 i=0,ik
35  work(i)=0
  fac1= (g(1)+g(2)+g(3))/(5*delta)
  fac2=-delta/120
  do 110 i=1,3
  do 110 j=1,i
    if (i.eq.j) goto 100
    s(i,j)=(e(1,1)*e(j,1)+e(1,2)*e(j,2))*fac1
    s(j,i)=s(i,j)
    goto 110
100  s(1,j)=(e(1,1)*e(j,1)+e(1,2)*e(j,2))*fac1
110  continue
  s(1,1)=s(1,1)+(5*n(1)+2*n(2)+2*h(3))*fac2
  s(2,2)=s(2,2)+(2*n(1)+5*n(2)+2*h(3))*fac2
  s(3,3)=s(3,3)+(2*n(1)+2*n(2)+5*h(3))*fac2
  s(2,1)=s(2,1)+(2*h(1)+2*n(2)+1*h(3))*fac2
  s(3,1)=s(3,1)+(2*n(1)+1*n(2)+2*h(3))*fac2
  s(3,2)=s(3,2)+(1*n(1)+2*n(2)+2*h(3))*fac2
  s(1,2)=s(2,1)
  s(1,3)=s(3,1)
  s(2,3)=s(3,2)
  do 200 k=1,2
  k1=k*nm1-nm2
  do 200 i=1,3
  do 200 j=1,3
    ih=2*(i*6+j)+k1

```

```

57 200 work(ih)=s(i,j)
58  icount=3
59  return
60  end

```



Jun 13 04:04 1981 elm302 Page 1

```
1      subroutine elm302(coor,work,iuser,iuser)
2      implicit real *4 (a-n,o-z)
3      real * 8 naam
4      dimension coor(1),work(1),iuser(1),user(1),
5      v      x(2,2),b(2)
6      common /cwork/ index1(100),index2(100)
7      v      /cact1/ ielem,itype,isoort,inpelm,icount,ifirst,not,istart
8      data naam /4helem/
9      nparms = coor(4)
10     c      *element matrices for the berkhoff equation (boundaryconditions)
11     c      itype=302 linear conforming element (mixedtype)
12     c      *local parameters
13     c      x      array containing the coordinates
14     c      d      array containing the local values of c*cg*q
15     c      ih     help variable
16     c      h1,h2  help variables to compute h
17     c      n      factor h (see description)
18     c      iusrt  position in array user (iuser(nstart))
19     c      a1,a2  factors a1,a2 (see description)
20     c      *computation of x and h
21     do 10 i=1,2
22         ih=2*index1(i)
23         x(i,1)=coor(ih+4)
24         x(i,2)=coor(ih+5)
25     10     b(i) = user(5+index1(i))
26     c      in these positions of "user" the values of "c*cg" are stored.
27     c      see: subroutine "local".
28     n1=x(1,1)-x(2,i)
29     n2=x(1,2)-x(2,2)
30     n=dsqrt(n1*n1+n2*n2)
31     if(h.ne.0) goto 30
32     c      *error 1: volume is 0
33         call ertrap(naam,1,15,itype,ielem,0)
34     30     if(ifirst.eq.1) goto 60
35     c      *ifirst=0, test inpelm and icount
36         if(inpelm.eq.2) goto 40
37     c      *error 2: inpelm#2
38         call ertrap(naam,2,20,itype,ielem,inpelm)
39     40     if(icount.eq.4) goto 50
40     c      *error 3: icount#4
41         call ertrap(naam,3,20,itype,ielem,icount)
42     50     nstart=iuser(isoort+istart)
43     c      test nstart
44         if(nstart.gt.5) goto 55
45     c      *error 4: |iuser(nstart)|<6
46         call ertrap(naam,4,15,itype,ielem,0)
47     55     iusrt1 = iuser(nstart+2)
48         iusrt2 = iuser(nstart+3)
49         if (iusrt1) 1,2,3
50     1     isrt1 = iabs(iusrt1)
51         a11 = user(isrt1)
52         a12 = a11
53         n1 = isrt1
54         go to 4
55     2     a11 = 000
56         a12 = 000
```



```

57          n1 = 0
58          go to 4
59      3      isrt1 = isrt1
60          n1 = isrt1 + work(4)
61      4      if (iusrt2) 5,6,7
62      5      isrt2 = iabs(iusrt2)
63          a21 = user(isrt2)
64          a22 = a21
65          n2 = isrt2
66          go to 3
67      6      a21 = 0d0
68          a22 = 0d0
69          n2 = 0
70          go to 3
71      7      isrt2 = iusrt2
72          n2 = isrt2 + work(4)
73      8      call prin1al(naam,5niuser,nstart+3,b,dummy,iuser)
74          nu = n1
75          if(n2.gt.n1) nu = n2
76      call prin1al(naam,4huser,nu,-7,user,dummy)
77      c          the values of a1 and a2 have been attributed
78      c          for the case: a1,a2 constant
79      if (iusrt1.ne.0.or.iusrt2.ne.0) go to 60
80      not=1
81      return
82      c
83      60     if (iusrt1.le.0) go to 11
84          a11 = user(isrt1)
85          isrt1 = isrt1 + 1
86          a12 = user(isrt1)
87      11     if (iusrt2.le.0) go to 22
88          a21 = user(isrt2)
89          isrt2 = isrt2 + 1
90          a22 = user(isrt2)
91      c
92      22     fac= h/24
93          work(5) = fac*(5*a11*b(1)+2*a12*b(2))
94          work(7) = -fac*(5*a21*b(1)+2*a22*b(2))
95          work(6) = fac*(2*a11*b(1)+2*a12*b(2))
96          work(9) = -fac*(2*a21*b(1)+2*a22*b(2))
97          work(10) = fac*(5*a21*b(1)+2*a22*b(2))
98          work(11) = fac*(5*a11*b(1)+2*a12*b(2))
99          work(12) = fac*(2*a21*b(1)+2*a22*b(2))
100         work(13) = fac*(2*a11*b(1)+2*a12*b(2))
101         work(14) = fac*(2*a11*b(1)+2*a12*b(2))
102         work(15) = -fac*(2*a21*b(1)+2*a22*b(2))
103         work(16) = fac*(2*a11*b(1)+6*a12*b(2))
104         work(17) = -fac*(2*a21*b(1)+6*a22*b(2))
105         work(18) = fac*(2*a21*b(1)+2*a22*b(2))
106         work(19) = fac*(2*a11*b(1)+2*a12*b(2))
107         work(20) = fac*(2*a21*b(1)+6*a22*b(2))
108         work(21) = fac*(2*a11*b(1)+6*a12*b(2))
109      c
110      c
111      return
112      c      *error messages

```

```

113      c      1: volume of element = 0
114      c      2: inbalm # 2
115      c      3: icount # 2
116      c      4: |iuser(isoort+istart)| < 6
117      c      5: sigma <= 0
118      c      6: length of array iuser too small
119      c      7: length of array user too small
120      end

```



Jun 13 04:04 1981 eir302 Page 1

```

1      subroutine eir302(coor,work,iuser,user)
2      implicit real *8 (a-h,o-z)
3      real * 8 naam
4      dimension coor(1),work(1),iuser(1),user(1)
5      dimension x(2,2),p(2)
6      common /cwork/ index1(100),index2(100)
7      v      /cact1/ ielem,itype,isoort,inpelm,icount,ifirst,not,
8      data naam /5nelrns/                                istart
9      nparms = coor(4)
10     c      *element vectors for the berkoff equation (boundaryconditions)
11     c      itype=302 linear conforming element (mixed type)
12     c      *local parameters
13     c      x          array containing the coordinates
14     c      ih         help variable
15     c      n1,h2      help variables to compute h
16     c      n          factor n (see description)
17     c      iusrt      position in array user (iuser(isoort+istart))
18     c      g          array containing the functionvalues (derivatives)
19     c      *computation of x and h
20     do 10 i=1,2
21         ih=2*index1(i)
22         x(i,1)=coor(ih+4)
23         x(i,2)=coor(ih+5)
24     10     b(i) = user(5+index1(i))
25     c      this section of "user" contains: c*cg.
26     n1=x(1,1)-x(2,1)
27     n2=x(1,2)-x(2,2)
28     h=dsqrt(n1*h1+n2*h2)
29     if(h.ne.0) goto 30
30     c      *error 1: volume is 0
31     call entrap(naam,1,15,itype,ielem,0)
32     30     if(ifirst.eq.1) goto 60
33     c      *ifirst=0, test inpelm and icount
34     call prin1(naam,5,iuser,isoort+istart,6,dummy,iuser)
35     if(inpelm.eq.2) goto 40
36     c      *error 2: inpelm#2
37     call entrap(naam,2,20,itype,ielem,inpelm)
38     40     if(icount.eq.4) goto 50
39     c      *error 3: icount#4
40     call entrap(naam,3,20,itype,ielem,icount)
41     50     nstart=iuser(isoort+istart)
42     c      test nstart
43     if(nstart.gt.5) goto 55
44     c      *error 4: |iuser(nstart)|<6
45     c      call entrap(naam,4,15,itype,ielem,0)
46     55     iusrt1 = iuser(nstart)
47     iusrt2 = iuser(nstart+1)
48     if (iusrt1) 1,2,3
49     1     isrt1 = labs(iusrt1)
50         g11 = user(isrt1)
51         g12 = g11
52         n1 = isrt1
53     go to 4
54     2     g11 = 0d0
55         g12 = 0d0
56         n1 = 0

```



Jun 13 04:04 1981 eir302 Page 2

```
57      go to 4
58      3      isrt1 = 1usrt1
59                      n1 = isrt1 + work(4)
60      4      if (1usrt2) 5,6,7
61      5      isrt2 = 1abs(1usrt2)
62          g21 = user(isrt2)
63          g22 = g21
64                      n2 = isrt2
65      go to 3
66      6      g21 = 0d0
67          g22 = 0d0
68                      n2 = 0
69      go to 3
70      7      isrt2 = 1usrt2
71                      n2 = isrt2 + work(4)
72      8      call prin1(naam,5,1user,nstart+3,0,dummy,1user)
73          nu = n1
74          if(n2.gt.n1) nu = n2
75      call prin1(naam,4,user,nu,-7,user,dummy)
76      c          the values of g1 and g2 have been attributed
77      c          for the case g1,g2 constant
78      if (1usrt1.ne.0.or.1usrt2.ne.0) go to 60
79      not=1
80      return
81      c
82      60     if (1usrt1.le.0) go to 11
83          g11 = user(isrt1)
84                      isrt1 = isrt1 + 1
85          g12 = user(isrt1)
86      11     if (1usrt2.le.0) go to 22
87          g21 = user(isrt2)
88                      isrt2 = isrt2 + 1
89          g22 = user(isrt2)
90      c
91      22     fac=h/6
92          work(6)= fac*(2*g11*b(1)+g12*b(2))
93          work(7)= fac*(2*g21*b(1)+g22*b(2))
94          work(8)= fac*(g11*b(1)+2*g12*b(2))
95          work(9)= fac*(g21*b(1)+2*g22*b(2))
96      return
97      c      *error messages
98      c      1:  volume of element = 0
99      c      2:  inpelm # 2
100     c      3:  icount # 2
101     c      4:  |1user(isort+1start)|<6
102     c      5:  sigma <= 0
103     c      6:  length of array 1user too small
104     c      7:  length of array user too small
105     end
```