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Buoyancy response of a disc to an embedded planet: a cross-code comparison at high resolution

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ABSTRACT

In radiatively inefficient, laminar protoplanetary discs, embedded planets can excite a buoyancy response as gas gets deflected vertically near the planet. This results in vertical oscillations that drive a vortensity growth in the planet’s corotating region, speeding up inward migration in the type-I regime. We present a comparison between PLUTO/IDEFIX and FARGO3D using 3D, inviscid, adiabatic numerical simulations of planet–disc interaction that feature the buoyancy response of the disc, and show that PLUTO/IDEFIX struggle to resolve higher-order modes of the buoyancy-related oscillations, weakening vortensity growth, and the associated torque. We interpret this as a drawback of total-energy-conserving finite-volume schemes. Our results indicate that a very high resolution or high-order scheme is required in shock-capturing codes in order to adequately capture this effect.

Key words: accretion discs – hydrodynamics – methods: numerical – planet–disc interactions.

1 INTRODUCTION

The observation of a pair of planets embedded in the circumstellar disc around PDS 70 (Keppler et al. 2018; Haffert et al. 2019) has cemented the concept that planets form and grow in circumstellar discs. Their final radial location – once the disc has dispersed – is typically determined by their migration history through (and interaction with) the disc. This is especially important for low-mass planets, which can migrate rapidly (e.g. Ward 1997). For a recent review see Paardekooper et al. (2022).

Planet–disc interaction can lead to several different sources of torques, which in turn drive the planet inwards or outwards (Ward 1997). Typical examples are the Lindblad torques by planet-driven spiral wakes (Ward 1986; Ogilvie & Lubow 2002), the horseshoe drag by the planet’s corotating region (Ward 1991; Masset 2001; Paardekooper & Papaloizou 2009), and the dynamical corotation torque experienced by a migrating planet in a disc with a radial vortensity gradient (Paardekooper 2014; McNally et al. 2017).

Different physical effects can give rise to additional sources of torques, such as thermal lobes in the planet’s corotating region (Lega et al. 2014), thermal input due to accretion heating (Benítez-Llambay et al. 2015; Masset 2017), or overdensities around waves caused by the disc buoyancy response near the planet (Zhu, Stone & Rafikov 2012). The latter, which could potentially amount to a few tens of per cent of the total torque, is the focus of this work.

In a recent study, McNally et al. (2020; hereafter MN20) showed that the torque associated with the disc buoyancy response can be significant, speeding up inward migration of low-mass planets. The torque was explained by identifying overdense lobes centred at a height $z \sim 0.5\text{--}1 H$, where H is the disc scale height, near the

planet’s corotating region. Equivalently, the associated torque can be attributed to a vortensity enhancement in the planet’s corotating region, weakening the stalling effect of the (outward) corotation torque (McNally et al. 2017).

Modeling the buoyancy response of the disc can be challenging from a numerical point of view. As MN20 showed, resolving the vortensity evolution, the associated dynamical corotation torque, and the planet’s migration track requires high-resolution 3D global models of planet–disc interaction that are integrated for thousands of orbits. Such calculations are very computationally expensive, and especially so when additional physics such as radiation transport are included (Yun et al. 2022). Therefore, great care should be exercised to ensure that the problem is sufficiently resolved by a given numerical scheme.

Both MN20 and Yun et al. (2022) modeled the effects of buoyancy-related torques on planet migration using FARGO3D (Benítez-Llambay & Masset 2016), a code developed primarily to study highly supersonic shear flows such as protoplanetary discs. FARGO3D uses an operator-split upwind scheme on a staggered mesh, a method that has been successfully applied to planet–disc interaction for many years. Nevertheless, this is not the only approach that has been used to model planet–disc interaction: shock-capturing, Godunov-scheme codes such as NIRVANA-III (Ziegler 2004), PLUTO (Mignone et al. 2007), and ATHENA++ (Stone et al. 2008) have been used extensively to study protoplanetary discs and planet–disc interaction as well. It is therefore worth investigating how different families of codes handle the subtle effects of the buoyancy response on planet migration.

In this study we compare the results of the global simulations of MN20 to functionally identical models using PLUTO. Our primary aim is to identify the source of the buoyancy-related torque, and the reason the two codes give significantly different results. We complement these models with local shearing-box simulations using IDEFIX (Lesur

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et al. 2023) and FARGO3D. Our goal there is to find the physical regime where these two codes recover similar solutions.

Our physical framework and numerical setup are described in Section 2. We present our results in Sections 3 and 4, and discuss our findings in Section 5. We then provide a summary and our conclusions in Section 6.

2 PHYSICS AND NUMERICS

In this section we describe the physical framework and our numerical methods. Our setup closely replicates that of MN20, so we briefly summarize the relevant equations here. We refer the reader to MN20 for a more detailed description.

2.1 Physics

We consider a disc of ideal gas with adiabatic index γ and mean molecular weight μ around a star with mass M_* . The gas has a volume density ρ , velocity field \mathbf{u} and pressure P given by the ideal gas law $P = (\gamma - 1)e$, where e is the internal energy density. In a cylindrical coordinate system $\{R, \phi, z\}$, the inviscid Navier–Stokes equations then read

$$\frac{\partial \rho}{\partial t} + \mathbf{u} \cdot \nabla \rho = -\rho \nabla \cdot \mathbf{u}, \quad (1a)$$

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla P - \nabla(\Phi_* + \Phi_p), \quad (1b)$$

$$\frac{\partial e}{\partial t} + \mathbf{u} \cdot \nabla e = -\gamma e \nabla \cdot \mathbf{u}. \quad (1c)$$

Here, $\Phi_* = -GM_*/r$ is the potential of the star at spherical distance $r = \sqrt{R^2 + z^2}$ from the star, Φ_p is the potential of the planet, which we describe below, and G is the gravitational constant. We can now define the pressure scale height $H = c_{s,\text{iso}}/\Omega_K$, where $c_{s,\text{iso}} = \sqrt{P/\rho}$ is the isothermal sound speed and $\Omega_K = \sqrt{GM_*/R^3}$ is the Keplerian frequency of the gas. The gas temperature is then $T = \mu c_{s,\text{iso}}^2/\mathcal{R}$, where \mathcal{R} is the ideal gas constant.

Assuming $u_R = u_z = 0$, a vertically isothermal profile ($\partial T/\partial z = 0$), axisymmetry ($\partial/\partial \phi = 0$), and radial power-law profiles for the midplane density ρ_{mid} and temperature such that

$$\rho_{\text{mid}}(R) = \rho_0 \left(\frac{R}{R_0}\right)^p, \quad T(R) = T_0 \left(\frac{R}{R_0}\right)^q, \quad (2)$$

we can derive an equilibrium state for the disc following Nelson, Gressel & Umurhan (2013),

$$\rho^{\text{eq}}(R, z) = \rho_{\text{mid}}(R) \exp\left[-\frac{1}{h^2} \left(1 - \frac{R}{r}\right)\right], \quad (3)$$

$$u_\phi(R, z) = R\Omega_K \left[1 + (p+q)h^2 + q \left(1 - \frac{R}{r}\right)\right]^{1/2}, \quad (4)$$

where $h = H/R$ is the disc aspect ratio. Finally, we define the surface density $\Sigma(R) = \int_{-\infty}^{\infty} \rho(R, z) dz$.

We now embed a planet with mass M_p at distance R_p from the star. The planet is treated as a sphere with uniform density and radius $\epsilon = 0.1H_p$. We can then write the planetary potential at a distance $d = |\mathbf{r} - \mathbf{r}_p|$ from the planet using the prescription of Klahr & Kley (2006), where a cubic spline is used for $d < \epsilon$ to prevent numerical singularities (see also Yun et al. 2022)

$$\Phi_p = \begin{cases} -GM_p/d, & d \geq \epsilon \\ -GM_p(3\epsilon^2 - d^2)/(2d^3), & d < \epsilon. \end{cases} \quad (5)$$

The indirect term by the planet–star system orbiting about their common centre of mass is included. The planet is typically held on

a fixed circular orbit at R_p , in which case the indirect term by the disc on the star is ignored. In models where the planet is allowed to migrate, that term is included (see Crida et al. 2022). Since we neglect the effect of self-gravity, we include the torque correction term of Baruteau & Masset (2008) in models with migrating planets.

In addition to global models, we study a local version of the problem using the well-known shearing box formalism (see e.g. Hawley, Gammie & Balbus 1995). This formalism can be derived by considering a small box orbiting at the local Keplerian velocity at radius r_0 , neglecting any curvature effects (hence working in Cartesian coordinates) and expanding the stellar potential in a small parameter L/r_0 , where L is the size of the box. The local version of equations (1a), (1b), and (1c) then reads

$$\frac{\partial \rho}{\partial t} + \mathbf{u} \cdot \nabla \rho = -\rho \nabla \cdot \mathbf{u}, \quad (6a)$$

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -2\boldsymbol{\Omega}_0 \times \mathbf{u} - \frac{1}{\rho} \nabla P - \nabla(\Phi_0 + \Phi_p) + \nu \nabla^2 \mathbf{u}, \quad (6b)$$

$$\frac{\partial e}{\partial t} + \mathbf{u} \cdot \nabla e = -\gamma e \nabla \cdot \mathbf{u} + \chi \nabla \cdot \left[\rho \nabla \left(\frac{e}{\rho} \right) \right], \quad (6c)$$

where $\boldsymbol{\Omega}_0$ is the angular velocity of the shearing box and $\Phi_0 = -3\Omega_0^2 x^2/2 + \Omega_0^2 z^2/2$ is the tidal potential. We have also included a kinematic viscosity ν and thermal diffusivity χ . We use the same planet potential as in the global case (see equation 5), with the planet located at the origin of our coordinate frame. A vertically isothermal equilibrium solution is given by

$$\rho(z) = \rho_0 \exp\left(-\Omega_0^2 z^2 / (2c_{s,\text{iso}}^2)\right), \quad (7)$$

together with $P = c_{s,\text{iso}}^2 \rho$ and $\mathbf{u} = (0, -3\Omega_0 x/2, 0)^T$. The isothermal sound speed $c_{s,\text{iso}}$ is initially constant throughout the box.

2.2 Dynamical corotation torque

McNally et al. (2017) have shown that a migrating planet experiences a torque by the material inside the corotating region given by

$$\Gamma_h = 2\pi \left(1 - \frac{\varpi(R_p)}{\varpi_h}\right) \Sigma_p R_p^2 x_h \Omega_p \left(\frac{dR_p}{dt} - u_R\right), \quad (8)$$

where x_h is the horseshoe half-width (Paardekooper, Lesur & Papaloizou 2010) given by

$$x_h = \frac{1.1}{\gamma^{1/4}} \left(\frac{0.4}{\epsilon/H_p}\right)^{1/4} \sqrt{\frac{M_p}{h_p M_*}} R_p, \quad \epsilon = 0.6H_p. \quad (9)$$

Here, ϵ is chosen to match 3D models (Müller & Kley 2012), $\varpi(R_p)$ is the vortensity at R_p for an unperturbed disc, and ϖ_h is the characteristic vortensity enclosed in the planet's corotating region. For a vertically integrated disc the vortensity is given by $\varpi = (\nabla \times \mathbf{u})/\Sigma \cdot \hat{\mathbf{z}}$, while for a 3D disc the equivalent quantity was shown by Masset & Benítez-Llambay (2016) to be

$$\varpi = \left[\int_{-\infty}^{\infty} \frac{\rho}{(\nabla \times \mathbf{u}) \cdot \hat{\mathbf{z}}} dz \right]^{-1}. \quad (10)$$

It can then be shown that for both 2D and 3D, the unperturbed vortensity is approximately $\varpi_0(R) \approx \frac{1}{2}\Omega_K/\Sigma^{\text{eq}}$, with $\Sigma^{\text{eq}} = \int_{-\infty}^{\infty} \rho^{\text{eq}} dz$.

2.3 Numerics

We mainly use the PLUTOV4.4 code (Mignone et al. 2007), with comparison models using FARGO3D (Benítez-Llambay & Masset

2016) and IDEFIX (Lesur et al. 2023). PLUTO and IDEFIX use a conservative finite-volume shock-capturing approach by solving the Riemann problem across all cell interfaces. By contrast, FARGO3D uses a finite-difference method with a staggered mesh to achieve momentum conservation to machine precision. All codes employ the FARGO method by Masset (2000), implemented in PLUTO by Mignone et al. (2012). This enables considerable speedup by subtracting the background disc rotation before solving equations (1a), (1b), and (1c) relaxing the strict time-step limitation by the quickly-rotating inner boundary. We note that in the PLUTO and IDEFIX implementations, the conservative form of equations (1a), (1b), and (1c) is solved instead [see e.g. equations (1) and (24) in Mignone et al. (2007)].

The version of FARGO3D that we use includes the specific entropy formulation described in McNally, Nelson & Paardekooper (2019b; see appendix A therein), which better conserves entropy in the absence of shocks. In PLUTO, we use the ENTROPY_SWITCH option to achieve the same effect. By default, we use second-order accurate time-marching (RK2) and reconstruction (LINEAR) schemes in PLUTO, but we investigate several numerical combinations in our study (see Section 3.3). Both codes use the flux limiter by Van Leer (1974).

We use PLUTO to run global simulations identical to the FARGO3D setup of MN20 to facilitate a fair comparison between the two codes. In these models, our grid and physical setup is identical to that of MN20, with a grid spanning $r \in [0.8, 4.0]$ au, $\theta \in [\pi/2 - \arctan(5h), \pi/2]$, and $\phi \in [0, 2\pi]$ with $N_r \times N_\theta \times N_\phi = 805 \times 125 \times 3141$ for most models, achieving a resolution of 25 cells per scale height. We assume a constant aspect ratio $h = 0.05$ (i.e. $q = -1$) and a density profile that corresponds to $\rho_0 = 0.003 M_\odot \text{ au}^{-3}$ and $p = -3/2$, or $\Sigma = 3340 \text{ g cm}^{-2} (R/\text{au})^{-1/2}$, with $R_0 = 1$ au. The planet is held on a fixed circular orbit at $R_p = 2$ au for at least 200 orbits before it is allowed to migrate. Finally, we use wave-damping boundary zones at the top and radial sides of the domain similar to McNally et al. (2019a). In FARGO3D this is implemented as a linear viscous pressure (Stone & Norman 1992), while in PLUTO we use the wave-killing method described in de Val-Borro et al. (2006) with a damping time-scale of 0.1 local orbits.

In addition to our global studies, we employ the shearing box formalism to compare IDEFIX to FARGO3D. The unit of distance is taken to be $H_0 = c_{s,\text{iso}}(t = 0)/\Omega_0$, and the domain considered is $x \in (0, 10) H_0$, $y \in (-20, 20) H_0$, and $z \in (0, 3.873) H_0$ following Zhu et al. (2012). This domain is covered at a minimum of $N_x \times N_y \times N_z = 160 \times 640 \times 80$ grid cells, yielding a resolution of 16 cells per H_0 . Boundaries are taken to be periodic in y , and in z we employ outflow boundaries at the top of the domain, while a symmetry condition is forced at $z = 0$. Towards the outer boundary in x , over the last $2H_0$ we employ wave damping as described in de Val-Borro et al. (2006). At $x = 0$, we force the solution to be point-symmetric around the origin (see Zhu et al. 2012).

Our local models use the vanilla version of FARGO3D, without the specific entropy formulation. For IDEFIX, we use standard linear reconstruction with RK2 time-marching. In the case of non-zero diffusion coefficients, we use super-time-stepping (Alexiades, Amiez & Gremaud 1996) to alleviate constraints on the time-step from viscosity and thermal diffusion. We do not employ the FARGO method in IDEFIX. While this leads to smaller time-steps, we have checked that FARGO3D with and without the FARGO method gives identical results, which indicates that numerical diffusion associated with time-stepping is not a major issue in this problem.

3 GLOBAL MODELS: PLUTO VERSUS FARGO3D

In this section we present a comparison between our global models with PLUTO and the models by MN20 using FARGO3D. We begin with a comparison of our fiducial models, and then investigate the effect of different numerical schemes in PLUTO.

3.1 Fiducial models

In their study, MN20 demonstrated that the buoyancy response by the disc to the embedded planet drives a time-dependent vortensity enhancement in the planet's corotating region, resulting in an excess negative corotation torque and the speeding up of inward migration. To measure the torque acting on the planet we sum all contributions from the disc inside the active domain

$$\Gamma = GM_p \int_{\text{disc}} \frac{\mathbf{r} \times \mathbf{r}_p}{d^3} \rho dV, \quad \mathbf{d} = \mathbf{r} - \mathbf{r}_p, \quad (11)$$

and normalize the result by a reference value Γ_0 (e.g. Kley & Nelson 2012)

$$\Gamma_0 = \left(\frac{M_p}{h_p M_\star} \right)^2 \Sigma_p R_p^4 \Omega_p^2. \quad (12)$$

Fig. 1 shows a comparison between the two codes for models where the planet is either fixed or allowed to migrate. We find that the two codes agree very well for the fixed planet case in 3D (blue and orange dashed curves), but show significant differences when the planet is allowed to migrate (solid curves). In particular, the PLUTO model (orange) shows a much weaker torque than the FARGO3D model (blue). Nevertheless, the 3D PLUTO model shows a more negative torque than a corresponding 2D, vertically integrated model (shown in green), implying that the torque-generating mechanism found in the 3D models of MN20 is operating to an extent. We note that the 2D model has been run with PLUTO, but the torque and migration tracks agree very well with a 2D model using FARGO3D.

In Fig. 2 we compare the migration tracks for all models shown in Fig. 1. Our results remain consistent: the planet in the 3D PLUTO model migrates slower than that in the FARGO3D model, and slightly faster than the corresponding 2D model. We can therefore conclude that the vortensity generation by the disc buoyancy response is present, but too weak to speed up inward migration in our PLUTO models at the same rate as in FARGO3D.

Given that our results agree very well between codes for the fixed planet case (as well as in 2D, not shown here), we expect that the two codes resolve the buoyancy response to a different extent, and therefore the vortensity growth rate is different. Fig. 3 shows the vortensity structure in the corotating region of the planet for the FARGO3D (top) and PLUTO (bottom) models, showing they are very different between the two codes. Specifically, the vortensity enhancement in FARGO3D is concentrated at R_p , while in the case of PLUTO we see closed streamlines with excess vortensity at the edges of the corotating region. Fig. 4 shows a radial profile of the azimuthally averaged vortensity in the corotating region to better highlight the magnitude of the vortensity excess.

To understand why the vortensity structure in the corotating region is so different, we turn to investigating the buoyancy response of the disc to the planet. As material approaches the planet along the azimuthal direction at a height z_0 it receives a 'kick' due to the planet's potential, and therefore oscillates about z_0 with a frequency given by the Brunt-Väisälä frequency,

$$N(z) = \sqrt{\frac{g_z}{\gamma} \frac{\partial}{\partial z} \left[\ln \left(\frac{P}{\rho^\gamma} \right) \right]}, \quad g_z = -\frac{\partial \Phi_\star}{\partial z}, \quad (13)$$

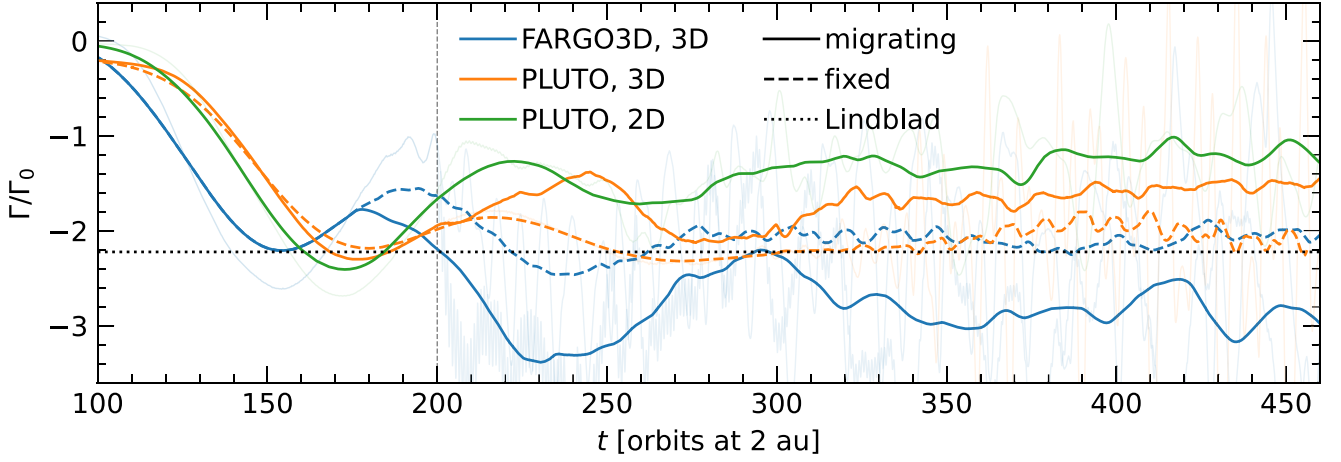


Figure 1. Torques acting on planets in our fiducial 3D models for FARGO3D (blue) and PLUTO (orange), for fixed (dashed) and migrating (solid) planets. The two codes agree very well for the fixed case, but differ substantially when the planet is allowed to migrate. The torque acting on the migrating planet in the 3D FARGO model is stronger than the fixed case, indicating feedback by the disc buoyancy response (MN20). For PLUTO, the 3D model shows a slightly stronger torque than a corresponding 2D setup (green), implying that this effect is present but not well resolved. The dots follow equation (14) in Paardekooper et al. (2010).

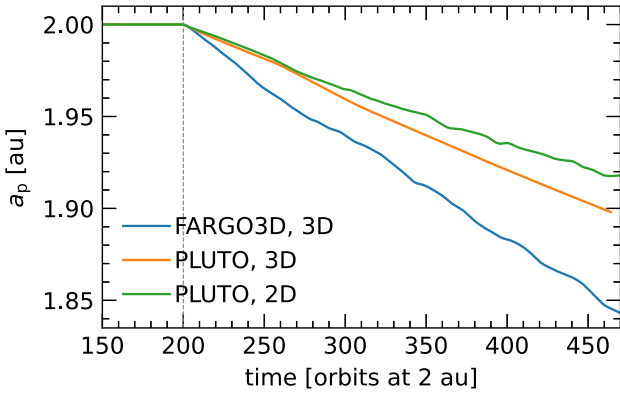


Figure 2. Migration tracks for the different models shown in Fig. 1. In the 3D model with FARGO3D (blue) the planet migrates inwards faster than a 2D model (green) due to the stronger negative corotation torque acting on it. Migration is much slower for the 3D PLUTO model (orange), suggesting that the effect that speeds up the planet in FARGO3D is not as strong in PLUTO.

with resonant modes being excited on lines of constant phase at azimuth ϕ_b given by

$$\phi_b(R, z) = \pm 2n\pi \sqrt{\frac{\gamma}{\gamma-1} \frac{\Omega_p - \Omega}{\Omega_K} \frac{H}{z} \left(1 + \frac{z^2}{R^2}\right)^{3/2}}, \quad (14)$$

where n is an integer (see also Zhu et al. 2015). We now plot the vertical velocity component of the gas $u_z(R, \phi)$ at a height $z = 2H$ after 200 planetary orbits for both codes and show the comparison in Fig. 5. We find that there are fewer modes visible for PLUTO, while FARGO3D shows more (albeit noisier) structure. We illustrate this more quantitatively in Fig. 6, where we plot $u_z(\phi)$ at $R = R_p - 2x_h \approx 0.96 R_p$ and $z = 2H$, showing that there are fewer and weaker oscillations resolved by PLUTO. This behaviour is consistent regardless of R and z , and most likely relates to the total variation diminishing (TVD) nature of Godunov schemes, which might damp wave-like behaviour overzealously to avoid spurious oscillations at the grid scale. By contrast, the algorithm used by FARGO3D results in possibly more accurate but also noisier structures.

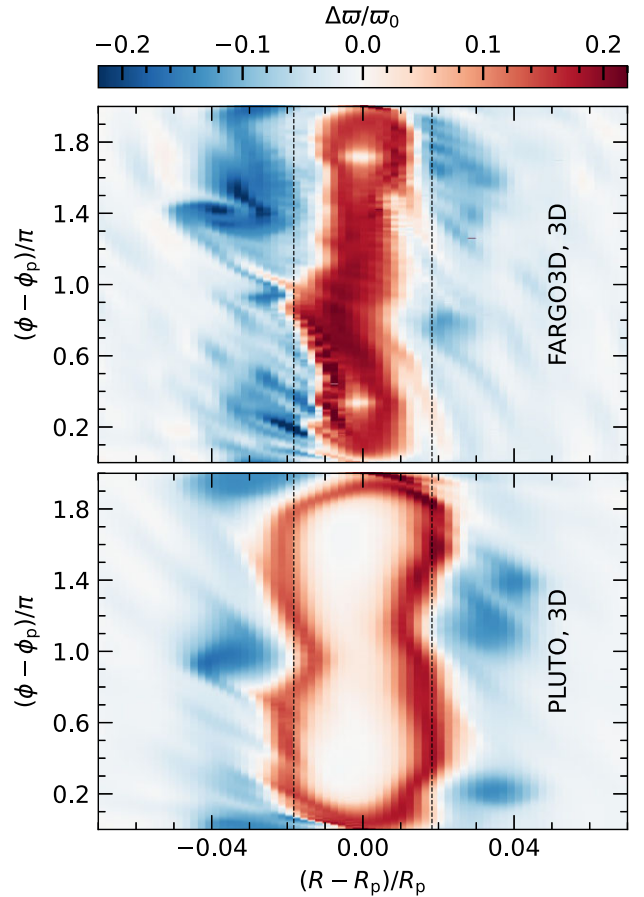


Figure 3. 2D heatmaps of the vortensity deviations in the planet's horseshoe region for the fixed 3D models at $t = 500 P_p$. Excess vortensity is concentrated about R_p for FARGO3D, but instead follows streamlines at the edge of the corotating region for PLUTO.

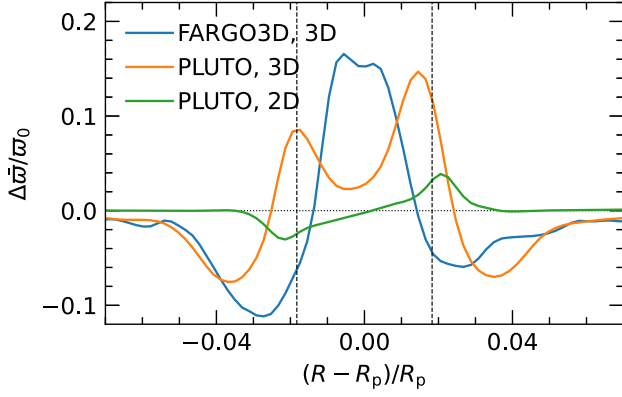


Figure 4. Azimuthally averaged vortensity deviations in the planet’s corotating region for fixed planets at $t = 500 P_p$. Even though there is a net ω growth in 3D models, their structure is quite different. Vertical dots mark $\pm x_h$.

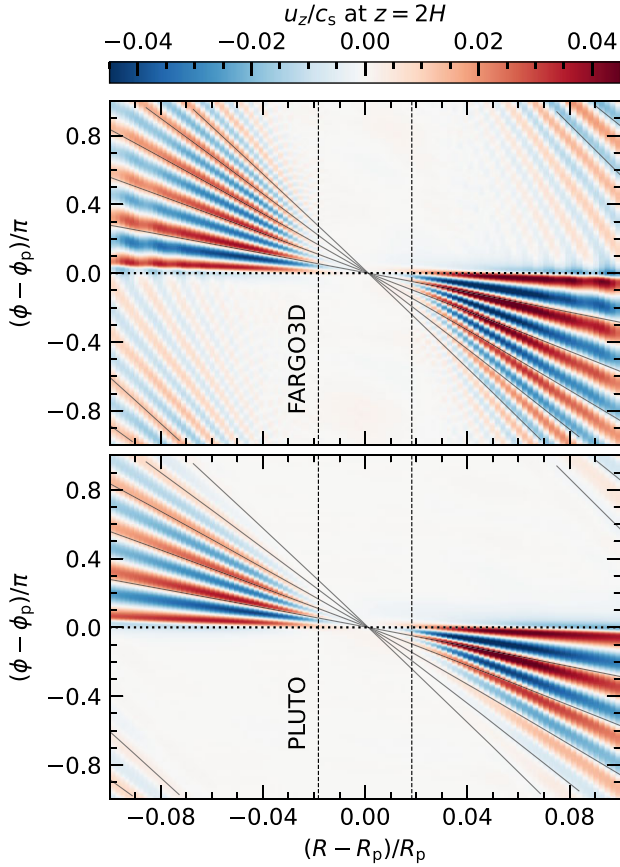


Figure 5. 2D heatmaps of the vertical gas velocity component u_z at a height $z = 2H$ for both codes, showing oscillations due to the disc buoyancy response. Diagonal black lines mark wavefronts of constant phase (see equation 14), and vertical black lines denote the limits of the corotating region $\pm x_h$. FARGO3D shows more buoyancy modes than PLUTO.

Assuming that these buoyancy modes are responsible for the vortensity growth in the corotating region, it becomes clear that the amplitude and number of modes resolved will decide the rate of vortensity growth as well as its radial structure. Given that higher-order modes $n \gg 1$ oscillate more tightly around the planet, they

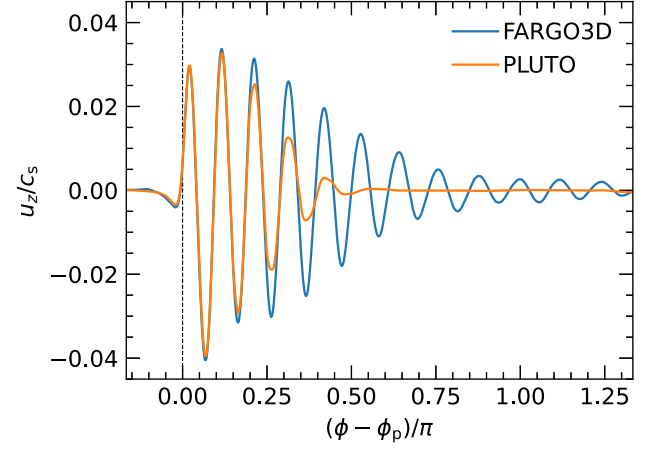


Figure 6. Gas u_z at $R = R_p - 2x_h$ and $z = 2H$, normalized to the local sound speed. Similar to Fig. 5, we see more oscillations in FARGO3D. The planet is marked by a vertical black line, and the undulations about it are due to the planet’s spiral arm.

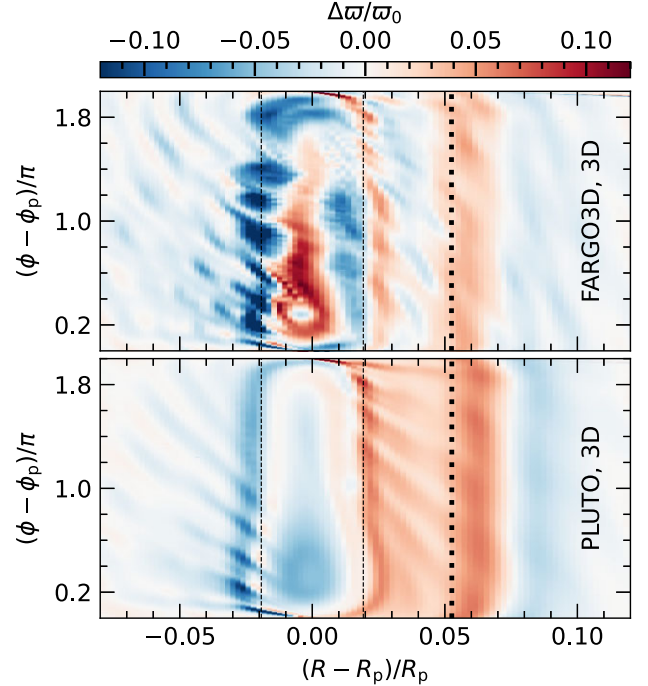


Figure 7. Similar to Fig. 3, for models with migrating planets at $R = 1.9$ au. A vortensity excess is clearly visible for FARGO3D. The dotted vertical line marks the planet’s initial location.

have the potential to both generate vortensity more efficiently inside the corotating region as well as closer to R_p . Since PLUTO resolves fewer modes, the vortensity growth is slower and more spread out in radius, resulting in a smaller excess that is confined to the edges of the corotating region. By contrast, in FARGO3D the vortensity excess is more concentrated at R_p . This can also be inferred from Fig. 4. As the planet migrates and the corotating region shrinks to a tadpole shape (Masset & Papaloizou 2003; Papaloizou et al. 2007), the excess trapped in the corotating region is higher in FARGO3D than in PLUTO (see Fig. 7), resulting in a stronger negative torque (as shown in Fig. 1).

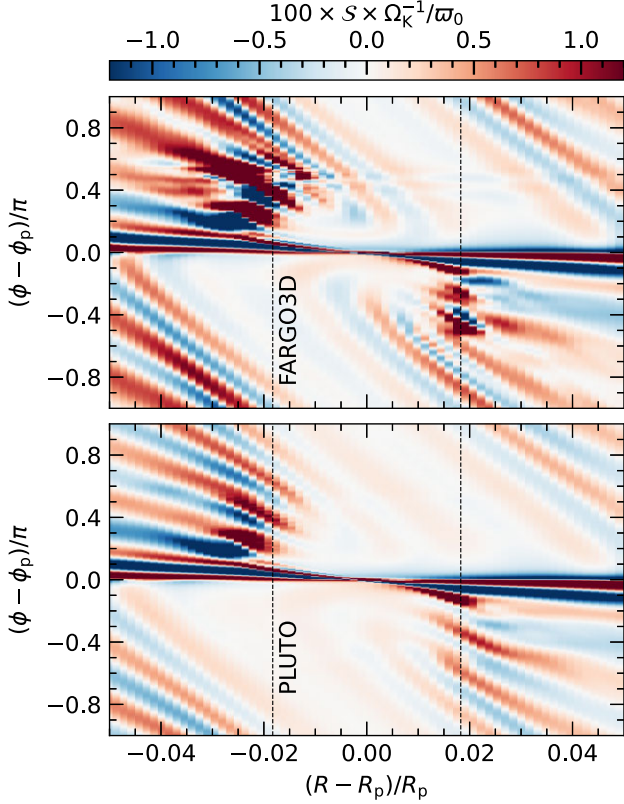


Figure 8. The vortensity source term through equation (16), normalized to a dimensionless quantity. We note the more saturated colors on the top panel (FARGO3D), indicating stronger vortensity generation. A horizontal ‘glitch’ at $\phi/\pi \pm 0.4$ is the result of spiral arms by emerging vortices.

3.2 Vortensity forcing

We now investigate whether the vortensity growth is tied to the buoyancy response. We start with the vortensity equation, obtained by taking the curl of the momentum equation (1b) and dividing by ρ

$$\frac{d\omega_\rho}{dt} = (\omega_\rho \cdot \nabla)\mathbf{u} + \frac{\nabla\rho \times \nabla P}{\rho^3}, \quad \omega_\rho = \frac{\nabla \times \mathbf{u}}{\rho}, \quad (15)$$

where $\omega_\rho = \omega_\rho \cdot \hat{z}$ is the potential vorticity and d/dt is the material derivative. The first and second term on the right-hand side relate to vortex stretching and baroclinic forcing, respectively. Given the vertically slanted nature of the buoyancy response (Zhu et al. 2012) and the presence of a global pressure gradient, both terms are expected to correlate with the buoyancy response and possibly contribute to vortensity growth.

Using equation (10), we can then approximate the vortensity growth rate as

$$\begin{aligned} \frac{d\varpi}{dt} &\approx \varpi^2 \int_{-\infty}^{\infty} \frac{1}{\omega_\rho^2} \left[(\omega_\rho \cdot \nabla)\mathbf{u} + \frac{\nabla\rho \times \nabla P}{\rho^3} \right] \cdot \hat{z} dz \\ &= S(R, \phi). \end{aligned} \quad (16)$$

We then plot this quantity S for both codes in Fig. 8, using high-resolution models from Section 3.3 for clearer results (the numerical and physical setups are otherwise exactly the same). We find that vortensity is generated along the planet-driven spiral shocks, which is to be expected, as well as along diagonals about the planet that coincide with the buoyancy modes in equation (14). This confirms that the vortensity generation is due to both ‘vortex stretching’

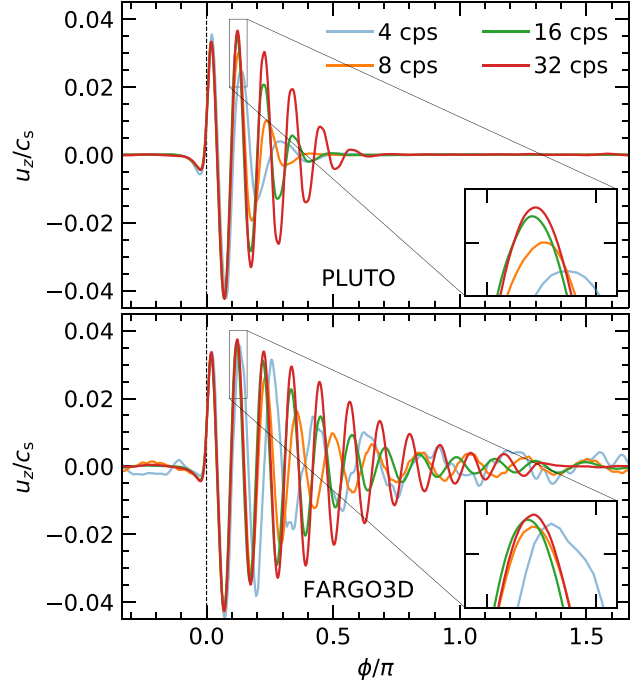


Figure 9. Gas u_z similar to Fig. 6 for different resolutions in units of cells per scale height (cps). Both codes resolve the $n = 1$ mode well at 16 cps (see insets), with FARGO3D showing the same amplitude for all resolutions (albeit a different phase for 4 cps). FARGO3D consistently resolves more modes than PLUTO, but also shows spurious oscillations at low resolutions (4–8 cps).

and baroclinic forcing by the disc’s buoyancy response inside the planet’s corotating region, with baroclinic forcing contributing to roughly 4–5 per cent of the total vortensity growth. We find stronger but significantly noisier vortensity generation for FARGO3D and larger n , consistent with our expectations above, supporting the idea that buoyancy modes might be overdamped in PLUTO (but also underdamped in FARGO3D). The source term also reaches closer to R_p for FARGO3D, which is consistent with the more concentrated vortensity excess in Fig. 4.

Based on our findings, we expect that the primary reasons for the two codes to differ are of numerical nature. However, it is still not clear which of the two codes offers the ‘more correct’ answer, as PLUTO tends to overdamp oscillations for this problem (and vice versa for FARGO3D). Since PLUTO offers a wide range of numerical options, we experiment with various configurations in the next section.

3.3 Different numerics

In the previous section we established that the differences between PLUTO and FARGO3D are primarily not physical, but instead most likely owe to the very different numerical methods used in each code. To investigate further, we run a series of tests between the two codes.

We begin with a resolution analysis, to test whether PLUTO recovers the behaviour in FARGO3D by simply increasing the resolution. For these tests we let the planet grow over one orbit and integrate for 10 planetary orbits. This allows enough time for the buoyancy patterns to be established in u_z , which we plot at $t = 10 P_p$.

Fig. 9 shows $u_z(\phi)$ at $R = R_p - 2x_h$ and $z = 2H$, similar to Fig. 6. We find that the $n = 1$ mode is already resolved at low resolution for both codes, and they both resolve more buoyancy modes as the resolution

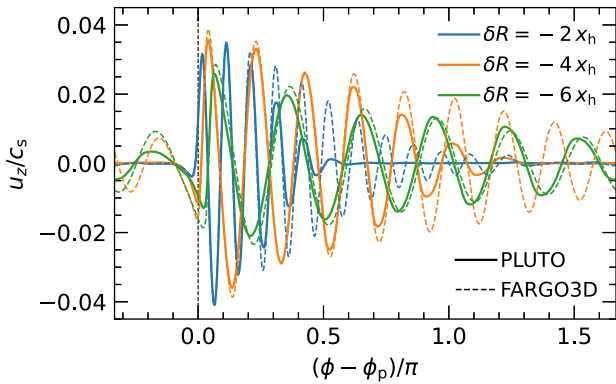


Figure 10. Gas u_z similar to Fig. 9, at different distances from the planet $\delta R = R - R_p$. FARGO3D once again shows stronger oscillations regardless of δR .

Table 1. Configuration flags used in PLUTO to test the effects of numerics, sorted by ascending accuracy.

Tag	Reconstruction	Time-stepping	Solver	Misc.
H2	LINEAR	RK2	h11c	(default)
HC2	LINEAR	RK2	h11c	CHAR_LIMITING
HL2.5	Lim03	RK2	h11c	CHAR_LIMITING
RW2.5	WENO3	RK2	roe	CHAR_LIMITING
RW3	WENO3	RK3	roe	CHAR_LIMITING
HW2.5	WENO3	RK2	h11c	CHAR_LIMITING
PPM	PARABOLIC	RK3	h11c	CHAR_LIMITING

The H2 setup is our default, and is used for every run unless otherwise stated.

increases. However, we also find that FARGO3D consistently resolves more and stronger modes at a given resolution compared to PLUTO, at the cost of some spurious oscillations far from the planet that become weaker with increasing resolution. In addition, FARGO3D achieves convergence for a given mode at a lower resolution than is necessary for PLUTO. Finally, we find that the phase of the buoyancy response is better captured at higher resolution, with modes packing together more tightly in both codes. Given that numerical noise in FARGO3D gives way to physical undulations at higher resolution, and the amplitude of the latter converges faster with increasing cell count in FARGO3D, we conclude that FARGO3D will resolve more and stronger buoyancy modes than PLUTO at any given resolution.

In Fig. 10 we then compare u_z at different radii or heights for both codes at a resolution of 16 cells per scale height. We find that both codes agree very well around the planetary spirals and $n = 1$ modes but are very different when the focus is buoyancy modes with $n > 1$, especially closer to R_p .

This behaviour implies that FARGO3D indeed recovers a more accurate solution for this problem than PLUTO, and especially so at higher resolution where the numerical shortcomings of FARGO3D reduce to the grid scale. With that in mind, we shift to PLUTO and run several tests with different numerical configurations at a resolution of 16 cells per scale height ($N_r \times N_\theta \times N_\phi = 528 \times 80 \times 2048$). We choose this lower resolution to significantly reduce computational costs, given that the results are qualitatively similar regardless of resolution.

From less to more accurate and computationally expensive, the numerical setups we used in PLUTO are summarized in Table 1. We note that not all models shown in this table are shown in plots. In that case, this implies that such a setup offers negligible gain compared to the setup right below it in the table.

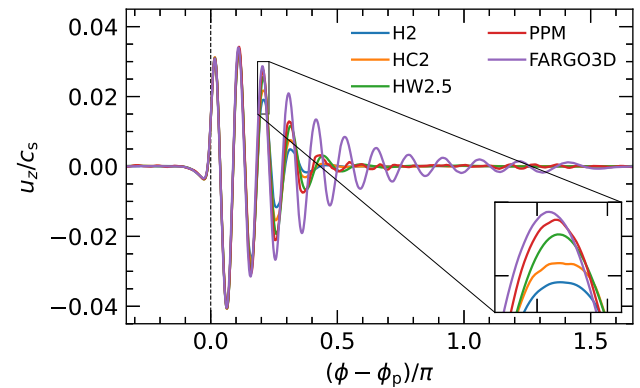


Figure 11. Gas u_z similar to Fig. 9 for different solvers in PLUTO. The third-order PPM setup achieves best results, but still not as good as FARGO3D.

Similar to above, we plot $u_z(\phi)$ at $R = R_p - 2x_h$ and $z = 2H$ in Fig. 11. All PLUTO models are compared against the corresponding FARGO3D model, showing that more accurate numerical setups result in more and stronger buoyancy modes. This supports the idea that FARGO3D recovers a more accurate solution, if one is willing to look past the numerical oscillations inherent in that code.

Interestingly, models HL2.5, RW2.5, RW3, and HW2.5 showed nearly identical results, highlighting the importance of a higher-order reconstruction rather than time-stepping. However, it is unexpected that the HW2.5 model offered slightly better results than the others listed in this paragraph, resolving each mode with an amplitude ≈ 0.7 percent larger even though model RW3 should be more accurate due to a less diffusive solver (roe instead of h11c) and a higher-order time-stepping scheme (RK3 instead of RK2). We did not investigate further on why that was the case.

3.4 Section summary

A comparison between PLUTO and FARGO3D showed that PLUTO results in a smooth linear flow around the planet but also does not resolve or incorrectly damps oscillations due to the disc buoyancy response by the planet. On the contrary, these modes are excited more accurately in FARGO3D, and their vortensity generating propensity is more correctly captured, at the cost of numerical noise being present. These patterns are consistent for very low to high resolution, with FARGO3D consistently recovering more and stronger buoyancy modes.

Using different numerical configurations in PLUTO, we showed that the reason for the differences between the two codes is tied to their numerical approach, which is fundamentally different for each code (shock-capturing, finite-volume Godunov schemes in PLUTO versus staggered mesh, finite-difference upwind methods in FARGO3D). PLUTO recovers more modes with higher-order schemes or more accurate solvers, but FARGO3D nevertheless achieves convergence for a given buoyancy mode at typically lower resolutions.

4 SHEARING BOX MODELS: IDEFIX VERSUS FARGO3D

In this section we describe the results using our local setup. We focus on a planet mass M_p such that $GM_p = 0.0058 \Omega_0^2 H_0^3$, following Zhu et al. (2012). The planet radius is set to $\epsilon = 0.7 H_0$, to make sure this length scale is resolved for all resolutions studied. Our standard resolution has an approximate cell size of $\Delta x = H_0/16$, so that even

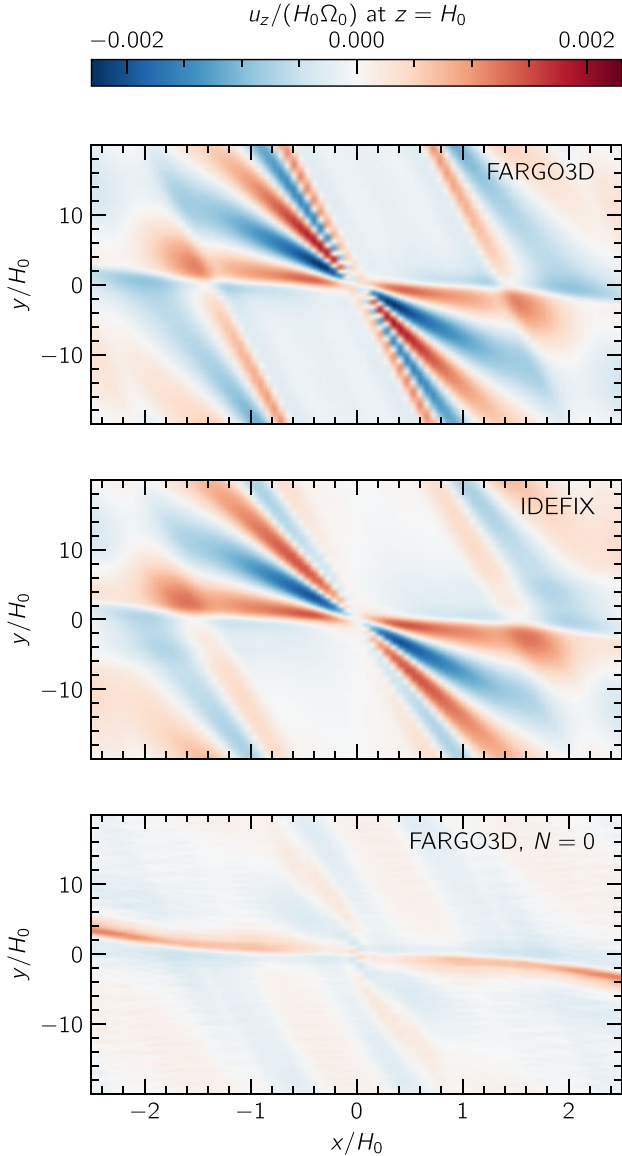


Figure 12. 2D heatmaps of the vertical gas velocity component u_z at a height $z = H_0$ at $\Omega_0 t = 20$, using FARGO3D (top panel), IDEFIX (middle panel), and FARGO3D with a polytropic initial condition (bottom panel). The top two panels show oscillations due to the disc buoyancy response.

when we go down a factor of 2 in resolution we still resolve the planet potential by several grid cells.

4.1 The case without explicit diffusion

In Fig. 12 we show the vertical component of the gas velocity at $\Omega_0 t = 20$ for our standard resolution of $\Delta x = H_0/16$, with no diffusion ($\nu = \chi = 0$). This time was chosen such that a well-defined buoyancy response has developed, but the distortions due to the periodic y boundaries remain limited. The top two panels show results obtained with an isothermal initial disc profile, using FARGO3D (top panel) and IDEFIX (middle panel). In the case of IDEFIX, the operator splitting technique used to integrate the vertical component of the stellar gravity makes it hard to start from exact numerical hydrostatic equilibrium (this is a well-known problem; see e.g. Bale et al. 2003; Paardekooper & Mellema 2006). Therefore the disc picks up a global

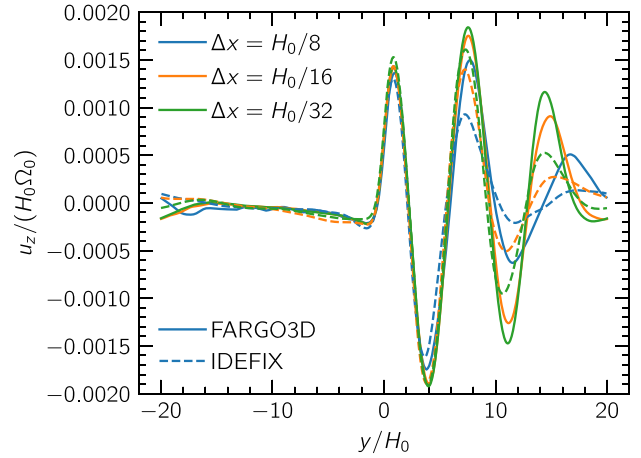


Figure 13. Vertical velocity at $\Omega_0 t = 20$ at $z = H_0$ and $x = -H_0/2$ for the two codes at different resolutions. Initial conditions are isothermal, and $\nu = \chi = 0$.

vertical oscillation that lasts throughout the simulation. While this could potentially limit the comparison between the codes, we here assume that this oscillation does not couple to the buoyancy-induced modes. This means we can subtract the average vertical velocity from the IDEFIX results and compare the residuals with FARGO3D. We note that the importance of this oscillation will be weaker for larger planet masses, as it is tied to the initial conditions and how well the initial profile is resolved.

In Fig. 12, local analogues of the buoyancy wave fronts can be clearly seen (see also Zhu et al. 2012). It is also clear that FARGO3D shows a stronger response, resolving more modes, just as in the global case. This comes at the expense of structure on the grid scale, as can be seen from the stair-like patterns in the top panel. IDEFIX shows a smoother structure, possibly due its TVD nature. In the bottom panel, we show the result for an initial condition that has constant entropy P/ρ^γ , which eliminates any buoyancy response. While this setup does lead to non-zero vertical velocities, the amplitude is very much reduced and the characteristic pattern of the buoyancy response has almost vanished. Any residuals may be partly due to the fact that the numerical equilibrium is not exactly isentropic, and they can be further reduced by increasing the resolution.

The stronger buoyancy response in FARGO3D is further illustrated in Fig. 13, where we show the vertical velocity as a function of y at $z = H_0$ and $x = -H_0/2$. At the first maximum, the codes agree to within 10 per cent for all resolutions. The spread increases to 16 per cent at the first minimum, located around $y = 4H_0$, and at the second maximum the spread is 50 per cent, with the lowest resolution showing more damping, as expected. As time evolves, the buoyancy-induced pattern keeps extending, with more oscillations being added, but once an oscillation has appeared, its amplitude remains steady. We can use this to estimate the impact of the subtraction of the global vertical oscillation in IDEFIX. By looking at slightly different times, we probe different phases of the vertical oscillation, and because the buoyancy-induced pattern is steady, any difference in the measured amplitudes is due to the subtraction process. From this, we estimate that we can measure the amplitudes confidently to approximately 10 per cent accuracy. With this in mind, we can say that at the first maximum the codes agree at all resolutions, while significant differences start to arise at the first minimum. Comparing FARGO3D and IDEFIX for equal resolutions, we find a difference at the second maximum of 40 per cent for $\Delta x = H_0/8$, 20 per cent for $\Delta x = H_0/16$,

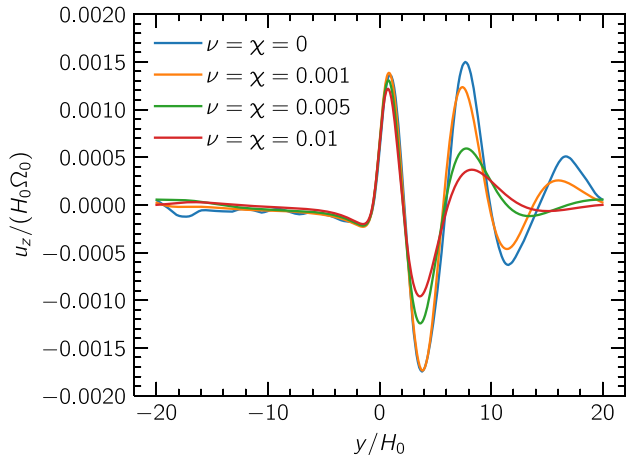


Figure 14. Vertical velocity at $\Omega_0 t = 20$ at $z = H_0$ and $x = -H_0/2$ for different levels of diffusion (units of ν and χ are $H_0^2 \Omega_0$). All results were obtained with FARGO3D at a resolution of $\Delta x = H_0/16$.

and 13 percent for $\Delta x = H_0/32$. It is plausible that if we would increase the resolution even further, FARGO3D and IDEFIX would agree at this second maximum to within our limit of 10 percent.

4.2 Adding explicit diffusion

We now consider cases where the diffusion coefficients are non-zero. For simplicity, we keep $\nu = \chi$, and therefore fix the Prandtl number to unity. This is purely for convenience, as in a protoplanetary disc it is usually the case that thermal diffusion (usually through radiation) dominates over viscous diffusion (through turbulence in the gas). Moreover, for the buoyancy response thermal diffusion is more important than viscous diffusion, and we found the results to be dominated by the value of χ rather than ν . Nevertheless, we want to work towards a case where numerical diffusion is subdominant in the buoyancy response, and hence keep the viscosity coefficient in line with thermal diffusion.

In Fig. 14 we show the vertical velocity, again at $\Omega_0 t = 20$, for different levels of diffusion. All results were obtained with FARGO3D at a resolution of $\Delta x = H_0/16$. Note that since the units of ν and χ are $H_0^2 \Omega_0 \approx c_s H_0$, the magnitude of ν in these units are the same as α in the usual α prescription for viscosity (Shakura & Sunyaev 1973). As expected, diffusion damps the buoyancy response, and while the first maximum is relatively unaffected, subsequent extrema are damped, to the extent that for $\nu = \chi = 0.01 H_0^2 \Omega_0$ there is only one more maximum established before the buoyancy response vanishes. Note also that the wiggly structure present in the inviscid run at $x < 0$ is effectively removed even for modest levels of diffusion.

Comparing Figs 13 and 14, we see that for our standard resolution of $\Delta x = H_0/16$, inviscid IDEFIX results are comparable to diffusive FARGO3D results with $\nu = \chi = 0.001 H_0^2 \Omega_0$. That is, for this specific problem setup (both numerical and physical), at this specific location in the domain, numerical diffusion in IDEFIX is similar to having $\nu = \chi = 0.001 H_0^2 \Omega_0$. At a lower resolution of $\Delta x = H_0/8$, the IDEFIX result is comparable to having $\nu = \chi = 0.002 H_0^2 \Omega_0$.

Perhaps a partial explanation for this behaviour lies in the TVD nature of shock-capturing methods such as IDEFIX and PLUTO. The concept of TVD is related to shock handling and the avoiding of the Gibbs phenomenon. Methods that are TVD will not show artificial post-shock oscillations. However, flow structures subject to the background shear will lead to an increase in TV, which is basically

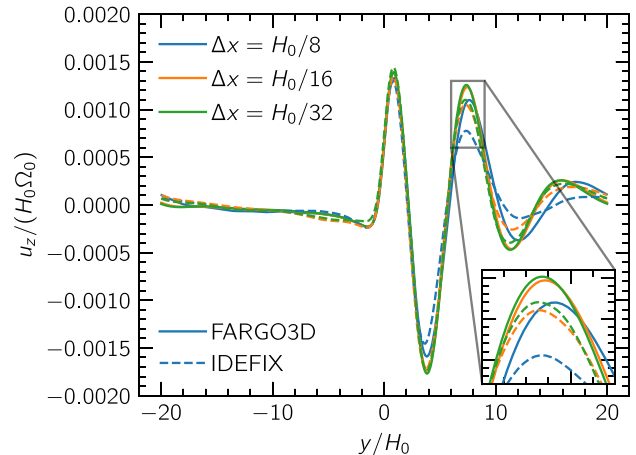


Figure 15. Vertical velocity at $\Omega_0 t = 20$ at $z = H_0$ and $x = -H_0/2$ for $\nu = \chi = 0.001 H_0^2 \Omega_0$. Shown are results using FARGO3D (solid curves) and IDEFIX (dashed curves) at different resolutions.

a measure of the integral of $|\partial W / \partial x|$, where W is any flow variable. Hence, IDEFIX will damp strongly sheared structures such as the higher order buoyancy response, even though no shocks are present (the buoyancy response basically consists of density variations at constant pressure). FARGO3D, on the other hand, explicitly adds artificial viscosity whenever a convergent flow is detected (see e.g. Stone & Norman 1992). When such a flow is absent, as is largely the case of the buoyancy response, FARGO3D will not apply any artificial viscosity and the density structures are sheared out until the grid scale becomes very apparent, as can be seen in the top panel of Fig. 12.

We now fix $\nu = \chi = 0.001 H_0^2 \Omega_0$, and compare FARGO3D and IDEFIX for different resolutions in Fig. 15, where we again show the vertical gas velocity at $\Omega_0 t = 20$, $x = -H_0/2$, and $z = H_0$. We first look at the codes individually, and look for signs of convergence with resolution. Since oscillations will be sheared apart more at larger y , they become more difficult to resolve and therefore a higher resolution is required. Focusing on the second maximum, the inset in Fig. 15 shows that both the solid curves (FARGO3D) and the dashed curves (IDEFIX) appear to converge to the same solution at our highest two resolutions. When comparing FARGO3D to IDEFIX, we need to keep in mind that the correction for the global vertical oscillation introduces an uncertainty of ~ 10 percent. Within that uncertainty, both codes actually agree on the amplitude of the second maximum. This is consistent with our observation above, that the inviscid IDEFIX results looked similar to FARGO3D results with diffusion coefficients $\nu = \chi = 0.001 H_0^2 \Omega_0$. This agreement is of course what *should* happen, in the regime where physical diffusion dominates over numerical diffusion, but this serves as a check that both codes are actually solving the same physical problem.

4.3 Vortensity forcing

We finally consider vortensity forcing in the local model, see equation (16). In Fig. 16, we show heat maps of this source term for three different runs, all with $\nu = \chi = 0$ at $\Omega_0 t = 20$ with $\Delta x = H_0/16$. The vortensity source consists of the vortex stretching term $(\omega_\rho \cdot \nabla) \mathbf{u}$ and the baroclinic forcing $(\nabla \rho \times \nabla p) / \rho^3$. In the isentropic case (bottom panel of Fig. 16), we found that the source is completely dominated by the vortex stretching term. This is expected, as this setup has constant entropy everywhere initially, and in the absence of thermal diffusion and shocks this means that entropy should remain

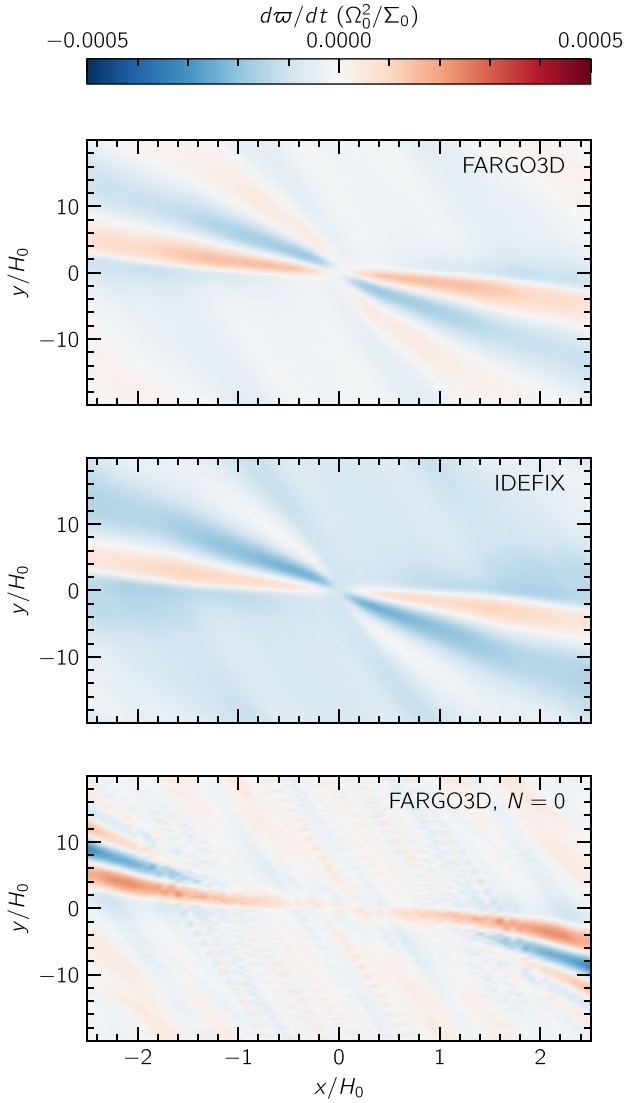


Figure 16. The vortensity source term, see equation (16). All simulations have $\nu = \chi = 0$ and $\Delta x = H_0/16$, and results are shown at $\Omega_0 t = 20$. Top panel: FARGO3D; middle panel: IDEFIX; bottom panel: FARGO3D with constant entropy initial condition.

constant at all times, rendering the $\nabla \rho \times \nabla P$ term identically zero. In reality, there is some numerical diffusion (in particular since we do not use the entropy-based variant of FARGO3D, see McNally et al. 2019b), as well as some converging flow structures that the code may interpret as shocks, which can lead to a non-zero vortensity source term. However, compared to the isothermal initial conditions, the resulting source is negligible. The stretching term mostly affects the spiral wave, with the corotation region largely unaffected.

In the top two panels of Fig. 16, we see evidence for vortensity generation by the buoyancy response. We found that this is again almost completely due to the stretching term, as the baroclinic forcing term is very localized to the region close to the planet. This is in contrast to the global simulations, where also the baroclinic forcing correlates with the buoyancy response, albeit at a small level compared to the stretching term. The lack of baroclinic forcing of vortensity away from the planet in the local model is because there is no pressure gradient to create a large enough $\nabla \rho \times \nabla P$ term. The only horizontal pressure variations occur close to the

planet (i.e. its atmosphere) and in the spiral wakes, which are essentially sound waves modified by rotation. The buoyancy response occurs at constant pressure, which strongly reduces the possibility of generating vortensity in the horseshoe region through baroclinic forcing. In the global models, on the other hand, there is a global radial pressure gradient, so that the density variations due to the buoyancy response combine with this radial pressure gradient to yield a non-zero $\nabla \rho \times \nabla P$ term. Both local and global models do agree that FARGO3D generates a stronger vortensity source than Riemann solvers such as IDEFIX and PLUTO, both through the vortex stretching term as well as the baroclinic forcing term.

4.4 Section summary

The local models have shown that the buoyancy response is damped more in IDEFIX compared to FARGO3D. In cases with explicit diffusion, both codes approach the same solution, but a detailed comparison is difficult because the IDEFIX solution includes a global vertical oscillation. Baroclinic forcing remains localized to a region very close to the planet because of the absence of a global pressure gradient, while vortex stretching is the main source of vortensity due to the buoyancy response.

5 DISCUSSION

In this section we discuss our results and their relevance to planet migration. We also note on the suitability of different codes to different physical problems.

5.1 Using u_z as a proxy for the buoyancy torque

A simplified analytical model by Lubow & Zhu (2014) showed that the width of the density ridges along buoyancy modes – which generate the related torque – depends on damping effects rather than pressure. This implies that a strict convergence test would require some explicit dissipation, as the density ridges would simply become narrower with increasing resolution otherwise, and then a comparison of torques at different resolutions and levels of dissipation. Given that the torque is only established after 3–4 libration time-scales, or ~ 200 – 300 orbits, such a detailed comparison would be prohibitively expensive from a computational point of view.

To circumvent this, we used u_z to measure the number and amplitude of buoyancy modes as a proxy for each code’s ability to adequately capture the buoyancy response, and assume that the resulting torque will be proportional to the latter. The advantage of this method is that it only requires the buoyancy pattern to be established, which only takes ~ 10 planetary orbits.

We nevertheless expect that in our inviscid models, where the only source of diffusion is numerical dissipation, modes will be narrower in azimuth at higher resolution. We indeed observe this mainly for PLUTO and IDEFIX and between resolutions of 4–16 cells per scale height (cps), with the width of each undulation narrowing further very slightly at 32 cps (see Figs 9 and 13). It is possible that the modes will not become infinitesimally thin due to nonlinear effects not predicted by the analytical model of Lubow & Zhu (2014).

5.2 How is planet migration affected?

Zhu et al. (2012) first demonstrated that an embedded planet can excite vigorous vertical motion in the form of buoyancy modes, driving a torque with magnitude comparable to the total torque acting on the planet. MN20 then showed how such a mechanism

can drive a vortensity excess in the planet’s corotating region, weakening the (outward) dynamical corotation torque (McNally et al. 2017) and speeding up inward migration. Their findings were limited to radiatively inefficient, quasi-adiabatic discs where cooling cannot damp these buoyancy oscillations, essentially maximizing their activity and showing that the planet can migrate inwards roughly twice as fast when this effect is considered.

Yun et al. (2022), however, showed that including a treatment of radiation transport by means of cooling via a flux-limited diffusion approach (see e.g. Levermore & Pomraning 1981) can significantly damp the buoyancy response when the time-scale associated with the buoyancy frequency in equation (13) is comparable to the thermal diffusion time-scale, suppressing it further as radiative diffusion becomes more efficient. This essentially quenches the buoyancy response for $z \gtrsim 2H$ even for optically thick discs at 1–2 au for their opacity model (Bell & Lin 1994).

Their study suggests that radiative diffusion could greatly limit the radial and vertical range where buoyancy modes can operate and drive vortensity growth near a planet. As a result, the related torque might not actually act on the planet except for in the innermost au (Ziampras et al., in preprint) and PLUTO/IDEFIX can still be safely used for planet–disc interaction in 3D geometries without much meaningful loss in quality.

5.3 Which code to trust?

As we discussed in Section 3.3, PLUTO/IDEFIX and FARGO3D use fundamentally different numerical approaches to solving equations (1a), (1b), and (1c). FARGO3D uses second-order accurate spatial reconstruction with a three-point stencil, but the time-stepping scheme is formally a single-step, first-order, operator-split approach. Nevertheless, the staggered mesh allows conservation of momenta to machine precision, which is crucial for smooth astrophysical flows.

On the other hand, PLUTO/IDEFIX are at least second-order accurate in both time and space, and can capture shocks while properly conserving the relevant quantities (mass, momentum, total energy). However, they are susceptible to the high-Mach problem (Trac & Pen 2004) which is present in highly supersonic flows such as protoplanetary discs. Essentially, more than 99 per cent of the total energy is kinetic, and therefore a strict conservation of total energy will cause any small errors during an energy update to propagate to the thermal energy. This means that even though total energy is conserved by construction, the ratio of kinetic to thermal energy will continuously change unpredictably, affecting the quality of results. This can be alleviated to an extent by instead conserving entropy (Ryu et al. 1993), which we, however, use for PLUTO models.

All in all, both codes have their shortcomings: PLUTO/IDEFIX are designed to capture shocks, while FARGO3D is best suited for linear, smooth, highly supersonic flows. Since this specific problem revolves around the excitation of linear waves, it makes sense that FARGO3D is more suitable here. When the focus is instead the planet-induced spiral shocks (Ogilvie & Lubow 2002; Rafikov 2002) it has been shown that even though both codes agree very well, the structure of a spiral shock will look smoother and more realistic in PLUTO/IDEFIX, while in FARGO3D one can see unphysical artifacts such as ‘ringing’ behind the shock front (see Fig. 17).

Of course, there exist problems where both codes are outshined by problem-specific approaches. When modeling hydrodynamic instabilities, for example, spectral codes such as SNOOPY¹ have

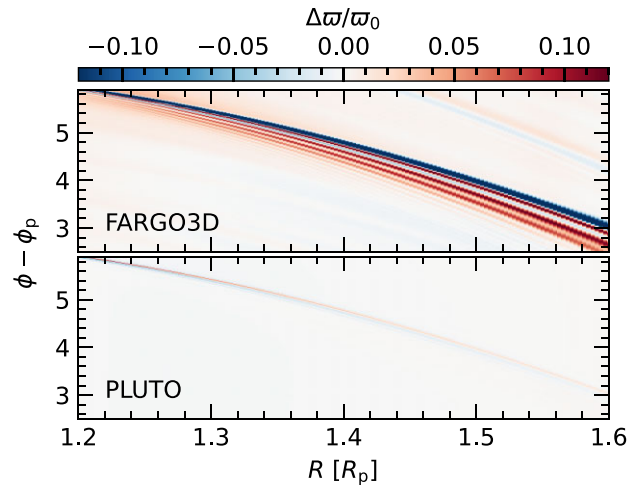


Figure 17. Perturbed vortensity around the spiral wake of the planet for vertically integrated, 2D models. The PLUTO solution shows perturbations only along the spiral shock as expected, while FARGO3D shows unphysical oscillations behind the shock front.

been shown to perform very well at capturing delicate effects that occur on very small length scales comparable to the grid scale (e.g. Cui & Latter 2022).

6 SUMMARY

We revisited the buoyancy response and the associated torque exerted on an embedded planet in an adiabatic protoplanetary disc (MN20). We compared the Godunov-scheme codes PLUTO (Mignone et al. 2007) and IDEFIX (Lesur et al. 2023) to the finite-difference code FARGO3D (Benítez-Llambay & Masset 2016).

We found that the two codes yield different results, and that the differences are inherent to the numerical schemes and not trivial. FARGO3D resolves more buoyancy modes, resulting in more efficient vortensity generation and a stronger negative torque, leading to a narrower high-vortensity corotating region and faster inward migration. With PLUTO, the vortensity profile in the corotating region is significantly different and the buoyancy-induced torque is substantially weaker for our fiducial model.

We also found that the buoyancy response in PLUTO is strongly dependent on resolution and the numerical configuration used (solver, reconstruction order, time-stepping order). With a sufficiently high-order (but also costly) numerical setup, PLUTO results resemble those of FARGO3D much more, implying that FARGO3D is more suitable for the study of this phenomenon.

Our local models confirmed these results, where we compared FARGO3D to IDEFIX. In this simplified geometry, again FARGO3D showed the stronger buoyancy response, keeping the density perturbations undamped even though they are taken apart by the background shear. IDEFIX, on the other hand, applies strong damping on structures close to the grid scale. In cases where we added explicit diffusion, both codes agreed on the solution at high enough resolution, with measured differences < 10 per cent. In the local models, the baroclinic forcing term vanishes away from the planet due to the lack of a global pressure gradient, which will lead to different evolution of the vortensity in the horseshoe region compared to global models.

We highlight, however, that the buoyancy-related torque might not be as present in realistic, radiative discs (Yun et al. 2022). This is

¹<https://ipag.osug.fr/~lesurg/snoopy.html>

heavily dependent on the disc model, and we will revisit this in future work. Nevertheless, given the niche of this effect, and the fact that both codes have been used successfully for planet–disc interaction and agree with each other in a wide variety of problems, we expect that this will not affect the codebase of finite-volume codes. Our intent is simply to raise awareness to the fact that sometimes one numerical approach will be more suitable than another for a particular problem.

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DATA AVAILABILITY

Data from our numerical models are available upon reasonable request to the corresponding author.

REFERENCES

- Alexiades V., Amiez G., Gremaud P.-A., 1996, *Commun. Numer. Methods Eng.*, 12, 31
- Bale D. S., LeVeque R. J., Mitran S., Rossmannith J. A., 2003, *SIAM J. Sci. Comput.*, 24, 955
- Baruteau C., Masset F., 2008, *ApJ*, 678, 483
- Bell K. R., Lin D. N. C., 1994, *ApJ*, 427, 987
- Benítez-Llambay P., Masset F. S., 2016, *ApJS*, 223, 11
- Benítez-Llambay P., Masset F., Koenigsberger G., Szulágyi J., 2015, *Nature*, 520, 63
- Crida A. et al. 2022, in Richard J., ed., SF2A-2022: Proceedings of the Annual meeting of the French Society of Astronomy and Astrophysics, p. 315
- Cui C., Latter H. N., 2022, *MNRAS*, 512, 1639
- de Val-Borro M. et al., 2006, *MNRAS*, 370, 529
- Haffert S. Y., Bohn A. J., de Boer J., Snellen I. A. G., Brinchmann J., Girard J. H., Keller C. U., Bacon R., 2019, *Nature Astron.*, 3, 749
- Hawley J. F., Gammie C. F., Balbus S. A., 1995, *ApJ*, 440, 742
- Hunter J. D., 2007, *Comput. Sci. Eng.*, 9, 90
- Keppler M. et al., 2018, *A&A*, 617, A44
- Klahr H., Kley W., 2006, *A&A*, 445, 747
- Kley W., Nelson R. P., 2012, *ARA&A*, 50, 211
- Lega E., Crida A., Bitsch B., Morbidelli A., 2014, *MNRAS*, 440, 683
- Lesur G. R. J., Baghdadi S., Wafflard-Fernandez G., Mouxion J., Robert C. M. T., Van den Bossche M., 2023, preprint ([arXiv:2304.13746](https://arxiv.org/abs/2304.13746))
- Levermore C. D., Pomraning G. C., 1981, *ApJ*, 248, 321
- Lubow S. H., Zhu Z., 2014, *ApJ*, 785, 32
- McNally C. P., Nelson R. P., Paardekooper S.-J., Gressel O., Lyra W., 2017, *MNRAS*, 472, 1565
- McNally C. P., Nelson R. P., Paardekooper S.-J., Benítez-Llambay P., 2019a, *MNRAS*, 484, 728
- McNally C. P., Nelson R. P., Paardekooper S.-J., 2019b, *MNRAS*, 489, L17
- McNally C. P., Nelson R. P., Paardekooper S.-J., Benítez-Llambay P., Gressel O., 2020, *MNRAS*, 493, 4382 (MN20)
- Masset F., 2000, *A&AS*, 141, 165
- Masset F. S., 2001, *ApJ*, 558, 453
- Masset F. S., 2017, *MNRAS*, 472, 4204
- Masset F. S., Benítez-Llambay P., 2016, *ApJ*, 817, 19
- Masset F. S., Papaloizou J. C. B., 2003, *ApJ*, 588, 494
- Mignone A., Bodo G., Massaglia S., Matsakos T., Tesileanu O., Zanni C., Ferrari A., 2007, *ApJS*, 170, 228
- Mignone A., Flock M., Stute M., Kolb S. M., Muscianisi G., 2012, *A&A*, 545, A152
- Müller T. W. A., Kley W., 2012, *A&A*, 539, A18
- Nelson R. P., Gressel O., Umurhan O. M., 2013, *MNRAS*, 435, 2610
- Ogilvie G. I., Lubow S. H., 2002, *MNRAS*, 330, 950
- Paardekooper S. J., 2014, *MNRAS*, 444, 2031
- Paardekooper S. J., Mellema G., 2006, *A&A*, 450, 1203
- Paardekooper S. J., Papaloizou J. C. B., 2009, *MNRAS*, 394, 2283
- Paardekooper S.-J., Lesur G., Papaloizou J. C. B., 2010, *ApJ*, 725, 146
- Paardekooper S.-J., Dong R., Duffell P., Fung J., Masset F. S., Ogilvie G., Tanaka H., 2022, preprint ([arXiv:2203.09595](https://arxiv.org/abs/2203.09595))
- Papaloizou J. C. B., Nelson R. P., Kley W., Masset F. S., Artymowicz P., 2007, in Reipurth B., Jewitt D., Keil K., eds, *Protostars and Planets V*. University of Arizona Press, Tucson, p. 655
- Rafikov R. R., 2002, *ApJ*, 569, 997
- Ryu D., Ostriker J. P., Kang H., Cen R., 1993, *ApJ*, 414, 1
- Shakura N. I., Sunyaev R. A., 1973, *A&A*, 500, 33
- Stone J. M., Norman M. L., 1992, *ApJS*, 80, 753
- Stone J. M., Gardiner T. A., Teuben P., Hawley J. F., Simon J. B., 2008, *ApJS*, 178, 137
- Trac H., Pen U.-L., 2004, *New Astron.*, 9, 443
- Van Leer B., 1974, *J. Comput. Phys.*, 14, 361
- Ward W. R., 1986, *Icarus*, 67, 164
- Ward W. R., 1991, Abstracts of the Lunar and Planetary Science Conference, 22, 1463
- Ward W. R., 1997, *Icarus*, 126, 261
- Yun H.-G., Kim W.-T., Bae J., Han C., 2022, *ApJ*, 938, 102
- Zhu Z., Stone J. M., Rafikov R. R., 2012, *ApJ*, 758, L42
- Zhu Z., Dong R., Stone J. M., Rafikov R. R., 2015, *ApJ*, 813, 88
- Ziegler U., 2004, *J. Comput. Phys.*, 196, 393

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