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# Nanopore electro-osmotic trap for the label-free study of single proteins and their conformations

Sonja Schmid<sup>1,3</sup>, Pierre Stömmmer<sup>2</sup>, Hendrik Dietz<sup>2</sup> and Cees Dekker<sup>1</sup>✉

**Many strategies have been pursued to trap and monitor single proteins over time to detect the molecular mechanisms of these essential nanomachines. Single-protein sensing with nanopores is particularly attractive because it allows label-free high-bandwidth detection on the basis of ion currents. Here we present the nanopore electro-osmotic trap (NEOtrap) that allows trapping and observing single proteins for hours with submillisecond time resolution. The NEOtrap is formed by docking a DNA-origami sphere onto a passivated solid-state nanopore, which seals off a nanocavity of a user-defined size and creates an electro-osmotic flow that traps nearby particles irrespective of their charge. We demonstrate the NEOtrap's ability to sensitively distinguish proteins on the basis of size and shape, and discriminate between nucleotide-dependent protein conformations, as exemplified by the chaperone protein Hsp90. Given the experimental simplicity and capacity for label-free single-protein detection over the broad bio-relevant time range, the NEOtrap opens new avenues to study the molecular kinetics underlying protein function.**

All vital functions in our cells are performed by proteins. Understanding protein function ultimately means understanding how these complex molecules behave over time, that is, how they adopt multiple energetically accessible conformations in solution in a thermal heat bath, often driven by additional external energy sources such as ATP, light, ion gradients and so on. Much progress has recently been made in obtaining high-resolution three-dimensional structures from cryo-electron microscopy, which is greatly informative but provides static snapshots, and hence the kinetics underlying protein function remain so far inaccessible. Single-molecule techniques can powerfully resolve molecular details that in bulk are masked by ensemble averaging, but they also face limitations such as a narrow temporal bandwidth (for example, in Förster resonance energy transfer (FRET)), interference from surface interactions (in atomic force microscopy) or the need for site-specific labelling (FRET dyes, handles for optical or magnetic tweezers). Scientists have developed new types of single-molecule traps aiming to overcome such shortcomings. For example, the anti-Brownian electro-kinetic trap<sup>1,2</sup> can hold a dye-labelled protein for seconds using real-time electro-kinetic feedback control. Electrostatic fluidic traps were shown to trap unfolded dye-labelled proteins under very low salt conditions, and to detect their charge<sup>3</sup>. Plasmonic nanoapertures can trap unlabelled single proteins<sup>4,5</sup>, but control and reproducible fabrication is challenging. Dielectrophoretic long-term trapping was demonstrated with silica beads but not on the level of individual protein molecules<sup>6</sup> while, only recently, first short-duration ( $\leq$ ms) protein detection has been shown<sup>7</sup>. Most of these methods require labelling of the proteins and sophisticated instrumentation, which limits their use to a handful of expert laboratories.

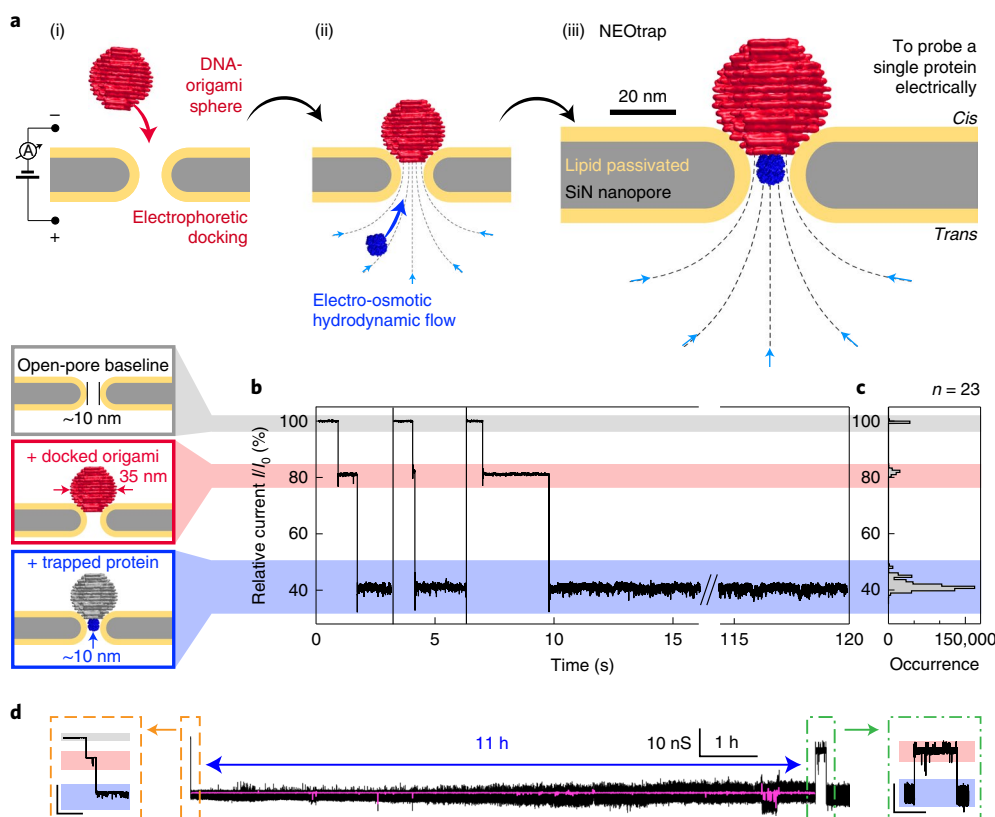
Here we present the nanopore electro-osmotic trap (NEOtrap), a nanopore-based device that can trap single unmodified proteins for up to hours, while recording conformation-sensitive currents at submillisecond resolution. We start by introducing the trapping concept that is based on docking a DNA-origami nanosphere onto

a solid-state nanopore and subsequent trapping of single proteins. Next, we demonstrate the single-molecule sensitivity and high resolution of the NEOtrap by discriminating between a variety of proteins on the basis of their mass and shape. Data for various pore sizes, voltages and ionic strengths inform on the trapping mechanism. Finally, we address protein function and present data that resolve nucleotide-dependent shifts in the conformations of the chaperone protein Hsp90, notably detected label-free and at the single-molecule level using the NEOtrap.

## Building the NEOtrap

Single-protein trapping using the NEOtrap occurs in two steps (Fig. 1), (1) in situ assembly of the molecular trap and (2) subsequent trapping of the protein of interest. The sensor consists of a solid-state nanopore, typically with a diameter  $d_{\text{pore}} \approx 20$  nm, which is transmission electron microscopy- (TEM-) drilled into a 20-nm-thick silicon-nitride membrane<sup>8</sup> and subsequently coated with a roughly 5-nm-thin lipid bilayer for passivation<sup>9</sup>, thus resulting in a nanopore with a diameter  $d_{\text{coated pore}} = d_{\text{pore}} - 10 \text{ nm} \approx 10 \text{ nm}$  (Fig. 1a). Voltage application (typically +100 mV) across the pore establishes a baseline ion current. The second ingredient of the NEOtrap is a charged and permeable nanostructure, in this case a 35-nm-diameter DNA-origami sphere<sup>10</sup>. Once such a negatively charged sphere diffuses close to the nanopore, the electric field drives it to the pore electrophoretically, where it gets docked onto the pore entrance<sup>11–13</sup> thus closing off a nanocavity. This leads to a characteristic and reproducible small current drop (around 20%), since origami nanostructures are permeable to ions<sup>11,14</sup>. The negative charge of the DNA phosphate backbone is screened by positive counter ions, and in the applied electric field, these counter ions migrate towards the cathode (upwards in Fig. 1a), which results in a strong hydrodynamic flux towards the pore. This nanofluidic effect is called electro-osmosis and is well studied in the context of surface charge in nanochannels and nanopores<sup>15–18</sup> (Supplementary Notes). Here, however, electro-osmosis is not surface driven—as

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**Fig. 1 | Working principle of the NEOtrap.** **a**, Illustration of the two-step trapping process: a DNA-origami sphere is electrophoretically docked onto a passivated solid-state nanopore, inducing hydrodynamic flow caused by electro-osmosis that facilitates protein trapping at the centre of the nanopore. 'Cis' and 'trans' reservoirs are defined in (iii). **b**, Current recordings start with a low-noise open-pore baseline (grey colour code); electrophoretic origami docking blocks roughly 20% of the current (red); the induced EOF traps a ClpP protein leading to further current blockage (blue). Three consecutive origami docking events followed by protein trapping (at 100 mV) are shown. In between these events, the trapped proteins were manually released by voltage inversion (−100 mV). The displayed relative current is normalized by  $I_{\text{op}}$ , the open-pore current (in the absence of a docked origami or a trapped protein) and 1 kHz low-pass filtered. More current traces and controls are shown in Supplementary Fig. 1. **c**, Histogram of 23 current traces such as those in **b**. **d**, One ClpP protein was trapped and recorded for 11 h. Insets show zoom views of ClpP trapping (orange, scale bar 10 nS versus 10 s) and spontaneous escape after 11 h (green, scale bar 10 nS versus 500 s); grey, red, blue shading as in **b**. The daylong recording was corrected for baseline drift and conductance increase over time (black, 1 kHz; pink, 100 Hz median). The unprocessed data are shown in Supplementary Fig. 2. The origami sphere in **a** and **b** was adapted with permission from ref. <sup>10</sup>, RSC.

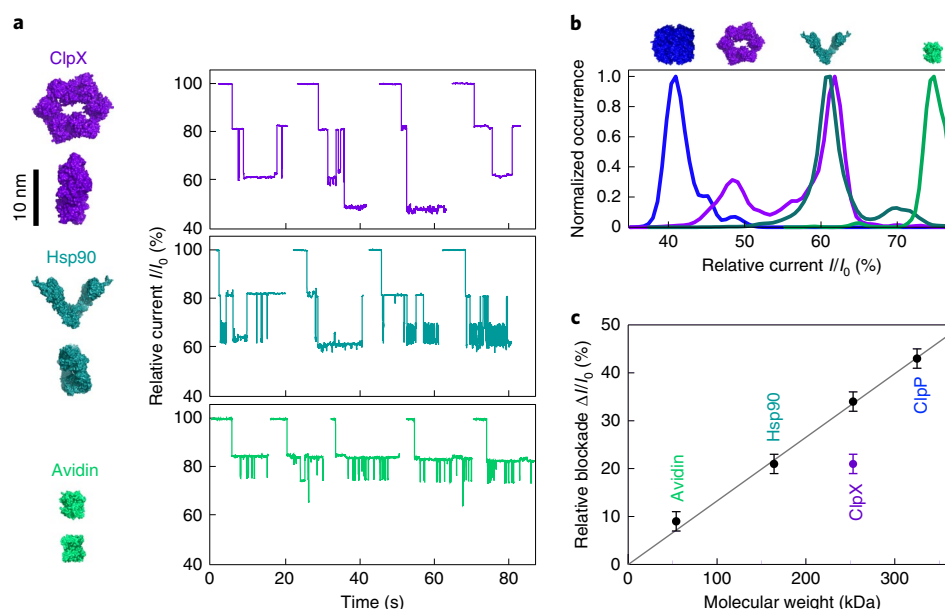
the lipid-bilayer surface is net neutral—but induced site-specifically by the DNA-origami structure that acts as a nanoporous negatively charged sponge. The electro-osmotic flow (EOF) drives proteins with the surrounding solvent towards the nanopore. As a result, a single protein can be trapped in the nanocavity formed by the coated nanopore that is capped by the docked origami sphere.

The trapping process can be monitored in real time (Fig. 1b): current traces show transitions from the open-pore current to a first partial blockade (due to the DNA sphere docking), and then to a second blockade associated with the trapping of a single protein into the cavity. This single-protein trapping is highly reproducible and comes with a characteristic current blockade, as shown in Fig. 1c and Supplementary Fig. 1. A single protein can be trapped reversibly for very long times. We demonstrated single-protein trapping up to 11 h (Fig. 1d and Supplementary Fig. 2). The NEOtrap is easy to assemble and can be reset simply by inverting the voltage, which releases the DNA-origami sphere and cleans up the trap, thus preparing it for a next round of assembly and trapping. Indeed, we routinely use a single nanopore for hundreds of single-protein trapping events in succession. Furthermore, the EOF caused by the DNA-origami sphere allows to trap proteins irrespective of their charge (Supplementary Table 1). The data shown in Fig. 1, for

example, were obtained for the negatively charged protein ClpP (net charge −55  $q_e$ ), which is drawn into the NEOtrap against the electrophoretic driving force acting on it, proving that the EOF dominates, and indicating that the NEOtrap can be used for a variety of charged biomolecules. Overall, the data show that the NEOtrap allows for easy and versatile long-term trapping of single proteins, which provides ample observation time of unmodified proteins.

### Distinguishing single proteins by size and shape

In Fig. 2, we present the ability of the NEOtrap to distinguish individual proteins on the basis of their current blockade signals observed in one and the same nanopore with a diameter  $d_{\text{pore}} = 20$  nm. We compare proteins of different mass (from 50 to 325 kDa), size (4.5 to 13 nm in diameter) and shape, namely avidin, Hsp90, ClpX and ClpP. All trapped proteins yielded characteristic signals, with a well-defined single peak in the blockade current histogram of most proteins (Fig. 2b)—except for ClpX, which produces two peaks, as discussed below. As expected from volume-exclusion considerations, larger proteins produced deeper current blockades, and a linear dependence was observed between the main blockade peak and the molecular weight of the proteins (Fig. 2c). An approximately linear dependence can be expected for globular proteins, as



**Fig. 2 | Mass- and shape-dependent single-protein identification with the NEOTrap.** **a**, Representative current recordings showing protein trapping events of ClpX, Hsp90 (apo form without nucleotides) and avidin for a 20 nm pore at 100 mV, with corresponding structures in front- and side-view (left images). The current is normalized by  $I_0$ , the open-pore current in the absence of a docked origami or a trapped protein and 1 kHz low-pass filtered. Although multiple avidins would fit into the NEOTrap cavity, only two occasions of double-trappings are observed, since there is only a small probability for this to occur within the relatively short duration of trapping events. **b**, Relative current histogram of 10–30 origami docking events with trapped ClpP, ClpX, Hsp90, avidin, colour-coded as in **c**. **c**, Relative current blockade  $\Delta I/I_0$  of each protein as a function of molecular weight. The line represents a linear fit through the black markers (slope of  $0.13 \pm 0.01\% \text{ kDa}^{-1}$ ; y-axis intercept of  $1.5 \pm 2.2\%$ ). The experiment was repeated with two pores in three separate experiments showing a similar linear mass dependence. For ClpX, two data points are observed, due to its disc shape, as discussed in the text. Protein Data Bank (PDB) codes ClpP, 1yg6; ClpX, 4i81; Hsp90, 2cg9; avidin, 4bx5.

larger-mass proteins occupy a larger volume in the electrolyte-filled pore, thus blocking the ion flow, accordingly. Vice versa, the relationship can also be used to estimate the mass of a trapped protein from the current blockade. While the capture rate was found to vary little in all cases, the trapped time (that is, inverse escape rate) increased exponentially as a function of molecular weight (Supplementary Fig. 3a), consistent with theoretical predictions (Supplementary Note 3). These kinetic rates yield estimated trapping potentials in the range of 10 to 20  $k_B T$  (Supplementary Fig. 3b).

The NEOTrap does not merely measure protein mass, but can also be used to sense protein shape. This can be seen from the data on the disc-shaped ClpX protein, where the current histogram (Fig. 2b) shows not one but two peaks,  $\Delta I/I_0 = 34 \pm 2\%$  and  $21 \pm 2\%$ , with the second peak at a much smaller blockade (that is, higher current level). We attribute the two current levels to two different orientations of the ClpX structure—parallel and perpendicular to the ion flow—in line with previous nanopore observations of orientation-specific ion-current signals by Yusko et al.<sup>19</sup>. Summing up, the data in Fig. 2 demonstrate the NEOTrap's potential to identify proteins on the basis of their mass, size and shape.

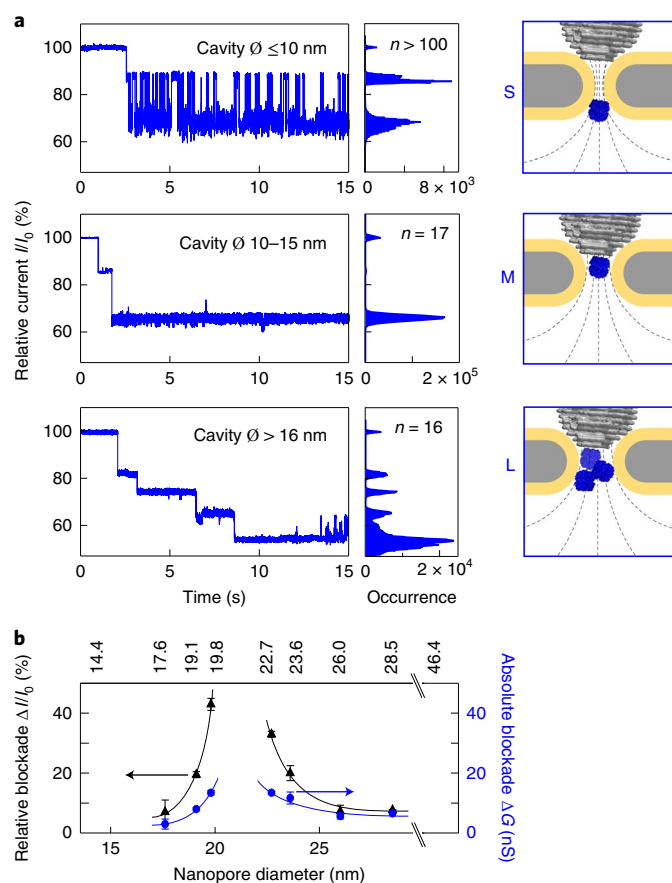
### Trapping dependence on pore size, voltage and ionic strength

The pore size has a strong effect on the observed single-protein trapping behaviour, as directly visible from the ClpP traces in Fig. 3a. Three regimes can be distinguished, namely small (S), medium (M) and large (L) sizes of the cavity after lipid-bilayer coating of the nanopores. A cavity with  $d_{\text{coated pore}} \leq 10 \text{ nm}$  (S) is too small to accommodate ClpP and many short-duration events are observed. These are attributed to ClpP proteins that bump onto the bottom rim of the cavity (in Fig. 3a), where they return to the same reservoir. Note also that the magnitude of EOF is weakest for these small

pore diameters (Supplementary Note 3). For cavity diameters of  $d_{\text{coated pore}} = 10 - 15 \text{ nm}$  (M), one ClpP can get stably trapped for hours (Fig. 1d). Since ClpP's size fits tightly into the cavity, the EOF is substantially blocked upon trapping and no additional (bumping) events are observed. The presence of a protein thus limits the trapping strength for capturing additional proteins, providing a mechanism for self-regulation of single-protein trapping in this regime. By contrast, for cavity sizes of  $d_{\text{coated pore}} > 16 \text{ nm}$  (L), multiple ClpP proteins are trapped consecutively, leading to a staircase behaviour in the current traces (Fig. 3a bottom). Here, each individual ClpP protein leads to the same conductance blockade of  $7.1 \pm 1.0 \text{ nS}$ , for up to three ClpP's that fit completely into the nanocavity, while the blockade for the fourth trapping was consistently smaller and exhibited larger variations, which can be attributed to end effects at the non-cylindrical pore opening<sup>20</sup>.

The cavity dimensions accessible for protein trapping are thus defined by the size of the protein itself (at the small limit), and by the translocation of the 35-nm-diameter origami sphere (at the large limit, Fig. 3b). As expected, the absolute conductance blockade per ClpP grew from the too-small-cavity regime (S) to the medium regime (M) that fully accommodates ClpP, thus causing the biggest blockade. However, the absolute conductance blockade dropped again from the medium to the large pore regime (L), from  $11.7 \pm 2.0 \text{ nS}$  at 23.6 nm pore diameter to  $5.7 \pm 1.0 \text{ nS}$  at 26 nm pore diameter. Microscopically, this may be explained by a sterically hindered ion flow when ClpP precisely fills the cavity, which causes an additional excluded volume for ions with associated water shells that adds to the total ClpP blockade.

The applied voltage affects the conductance blockade caused by the docked origami sphere differently compared to the one caused by the trapped protein, see Fig. 4a for data of two pore sizes and Fig. 4b for corresponding current traces. The blockade levels of the



**Fig. 3 | Pore size-dependence of NEOtrap signals.** **a**, Relative current traces and histograms showing ClpP trapping in three pore size regimes as illustrated on the right. From top to bottom, S shows a too small cavity leading to short trapping events where ClpP cannot enter the pore but merely blocks the entrance temporarily (here  $d_{\text{coated pore}} = 9.5$  nm); M, a cavity fitting the protein size leading to long-term single-protein trapping (here  $d_{\text{coated pore}} = 14.0$  nm) and L, a cavity fitting multiple proteins leading to consecutive trappings of multiple proteins with discrete blockades (here  $d_{\text{coated pore}} = 16.4$  nm). Associated histograms combine many protein trapping events as specified. The protein-accessible cavity diameter is specified:  $d_{\text{coated pore}}$  as defined in the text. **b**, Relative (black) and absolute (blue) current blockade levels of a single ClpP at 100 mV, obtained over a large range of nanopore diameters:  $d_{\text{pore}} = 14 - 46$  nm (where  $d_{\text{pore}}$  is deduced from conductance before coating, Supplementary Note 3). Error bars represent the full-width at half maximum (FWHM) of Gaussian fits to current histograms over 10–30 origami docking events per datapoint. Lines are guides to the eye. The origami sphere in **a** was adapted with permission from ref. <sup>10</sup>, RSC.

trapped ClpP proteins stayed constant with applied voltage. This indicates that the protein stays intact, and protein unfolding by electrostatic or hydrodynamic forces can be excluded for the voltage range probed. In contrast, the current blockade levels of the docked origami sphere grew significantly larger at higher voltages. We attribute this to the DNA-origami sphere that is gradually pressed more tightly onto the nanopore with increasing voltage.

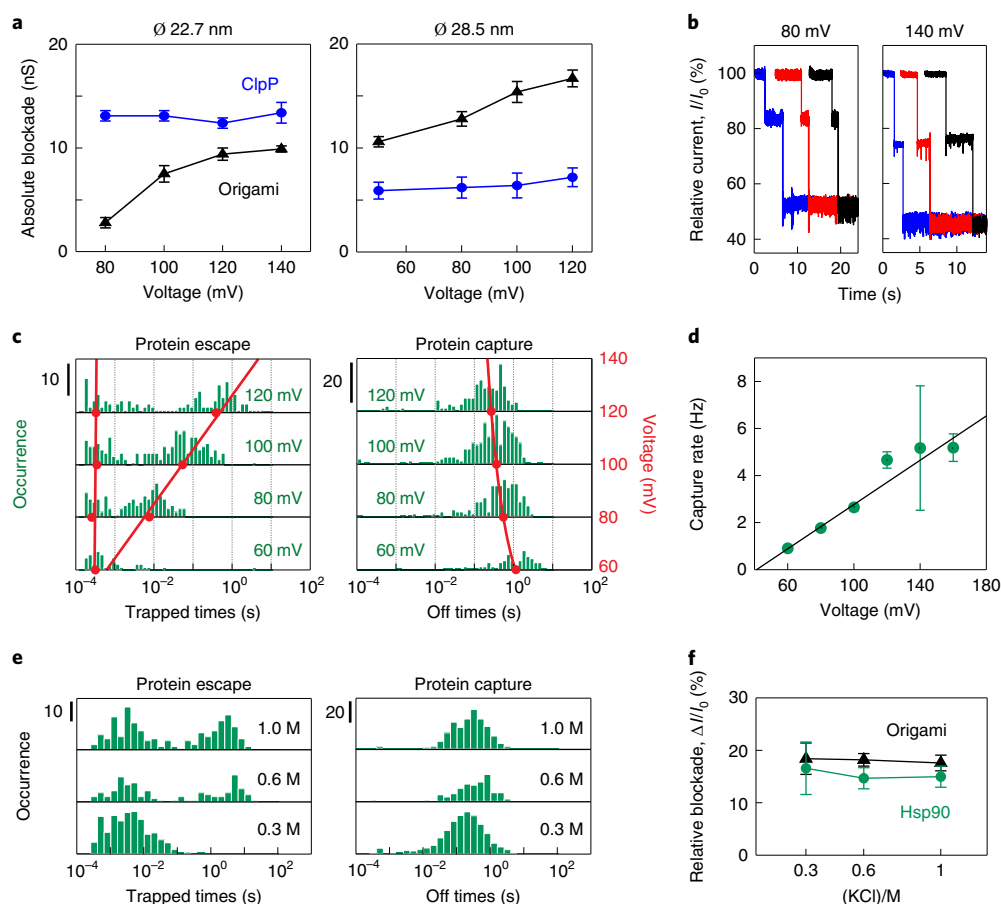
Protein trapping kinetics provide further insight into the trapping mechanism of the NEOtrap. Here, we studied Hsp90 as a model protein, because the very low escape rate (<1 per hour) of the stably trapped ClpP precluded sufficient event statistics. Figure 4c shows histograms of the protein trapped times (protein residence times) and the non-trapped off-times (empty cavity) for Hsp90, which inform on the escape and capture process, respectively.

Two escape pathways are consistently observed, a fast escape process with a trapped time that is voltage independent and a majority escape process that is slower, with a trapped time that grows exponentially with voltage. This exponential dependence is expected for a particle escape process that involves an energy barrier crossing as described by Boltzmann statistics (Supplementary Notes). The voltage-independent fast escape process may involve protein escape through the origami side as a result of temporal fluctuations in the docking of the sphere. Protein capture, on the other hand, exhibits a single-exponential off-time distribution with a linear voltage dependence of the capture rate (Fig. 4c right), as theoretically predicted. The linear fit of the capture rate in Fig. 4d, drops to zero at roughly 40 mV, which is of the order of the minimal voltage required for permanent docking of the origami sphere (here, roughly 60 mV). On the high voltage end, the stability of the lipid coating limits the experimentally accessible voltage range to  $\leq 160$  mV. A voltage of around 100 mV was found to be optimal regarding signal-to-noise ratio and trapping kinetics in the pore diameter range of 20–25 nm.

Scanning three different ionic strengths (Fig. 4e), we find the longest trapped times under an intermediate to high KCl concentration (0.6–1 M), while only the short-trapped population is observed at 0.3 M KCl (see also current traces in Supplementary Fig. 6). The capture rate, on the other hand, remains constant over the three ionic strengths. The relative blockade also remains unchanged within experimental uncertainty (Fig. 4f), which is further evidence for a non-invasive trapping mechanism that does not affect the protein stability, for both ClpP and Hsp90. While the signal-to-noise ratio of nanopore experiments generally decreases with decreasing ionic strength of the buffer, these data show that experiments at lower salt concentrations (currently 300 mM KCl) are possible and physiological salt conditions (150 mM) are well within reach in future studies. Furthermore, many proteins tolerate 600 mM KCl, and higher voltages may be applied to compensate for a shallower trapping potential at lower salt conditions (Fig. 4e).

### Label-free detection of nucleotide-dependent conformational shifts in the chaperone Hsp90

Finally, we demonstrate the NEOtrap's capacity for label-free detection of protein conformational heterogeneity and nucleotide-binding-induced conformational shifts in proteins, such as Hsp90 (Fig. 5). The diverse conformations of Hsp90 play a key role in the functional cycle of this molecular chaperone<sup>21</sup>, which is a central metabolic hub for protein homeostasis in the cell<sup>22</sup> and an enabler of cancer adaption<sup>23</sup>. Figure 5a depicts various conformations of homo-dimeric Hsp90, its two N-terminal nucleotide-binding sites and the many non-covalent bonds that the nucleotide (here ATP) engages in, which critically stabilize multiple domains of the protein<sup>24</sup>. Figure 5b,c compare NEOtrap recordings of Hsp90 under four different nucleotide conditions: with 5 mM AMP-PNP, ATP, ADP or in the absence of any nucleotide termed 'apo'. Clear differences are observed: in the presence of ATP or the non-hydrolysable ATP-analogue AMP-PNP, one dominant peak is observed at a low conductance blockade of  $-5.4 \pm 0.2$  or  $-5.3 \pm 0.2$  nS, respectively, while larger blockades are absent, which we interpret as Hsp90 prevailing in one compact conformation (in line with existing literature discussed below). Larger blockades were occasionally observed in the presence of ATP (see representative traces in Fig. 5b, and the shoulder in the ATP histogram in Fig. 5c). In the presence of ADP, more diverse blockade levels are detected, indicating a more diverse conformational ensemble populated by Hsp90. Last, in the absence of any nucleotide, where Hsp90 is missing essential stabilizing contacts in the N-domain, the blockade levels are distributed the most heterogeneously. Note also that the higher blockade levels exhibit larger fluctuations in the current traces (Fig. 5b), suggesting increased conformational dynamics in the millisecond range.



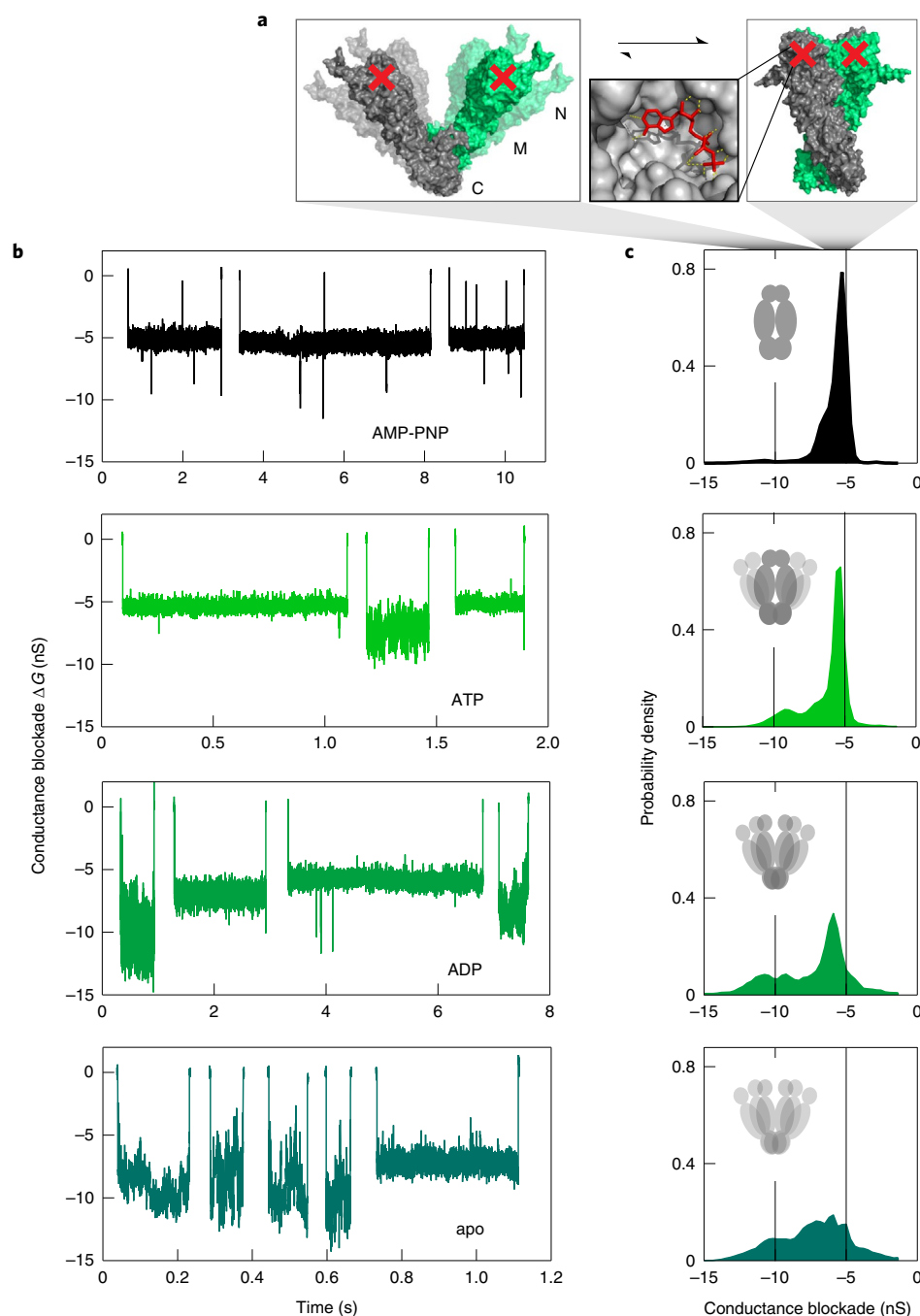
**Fig. 4 | NEOtrap characteristics as a function of voltage and ionic strength.** **a**, Voltage dependence of ClpP (blue) and origami blockades (black) for two pore diameters as specified. Error bars represent the FWHM of Gaussian fits to current histograms of 10–30 docking and trapping events per datapoint. **b**, Example current traces of the  $d_{\text{pore}} = 22.7$ -nm case in **a**, under the maximum and minimum voltage as specified. **c**, Hsp90 protein trapping kinetics (from one origami docking) for a range of applied voltages, as specified. Scale bars specify number of counts. The protein trapped times (left) reproducibly split into two populations—highlighted with red lines, which are linear fits to the population centres (red dots) determined with lognormal fits—indicating two escape pathways (observed in  $\geq 5$  measurements per condition). The off-times (no protein trapped, right panel) form single-exponential distributions described by one voltage-dependent capture rate. The red line (and right axis) is a linear fit to the capture rates as a function of voltage. Corresponding dot plots are shown in Supplementary Figs. 4 and 5. **d**, The protein capture rate grows linearly with applied voltage (here for multiple origami docking events). The black line is a linear fit with slope of  $0.047 \pm 0.003 \text{ Hz mV}^{-1}$  and y-axis intercept of  $-1.9 \pm 0.21 \text{ Hz}$ . The reported rates and error bars were obtained as means and standard deviations from bootstrapping single-exponential fits to off-time histograms (Methods). **e**, Protein trapping kinetics for Hsp90 for different KCl concentrations of 0.3–1 M. Dwell-time histograms include multiple origami docking events. Scale bars specify number of counts. Similar to the data in **c**, the trapped times split into two populations and the off-times (no protein trapped) are single-exponentially distributed. Note that the capture rate barely changes with salt concentration. Long trapped events are underrepresented by the recording time (here 20 s per origami docking, see current traces in Supplementary Fig. 6). Dwell-time distributions of minute-long recordings are shown in Supplementary Fig. 7. **f**, No significant salt dependence was observed for the relative current blockades induced by the docked origami or trapped Hsp90. Error bars represent the FWHM of Gaussian fits to current histograms over 10–30 origami docking events per datapoint. All Hsp90 data in this figure were measured in the presence of 5 mM AMP-PNP.

The striking differences observed in these current recordings of unlabelled Hsp90 reveal different conformational states of the Hsp90 dimer: AMP-PNP binding strongly stabilizes Hsp90's compact conformation<sup>24–26</sup>. Consistently, we observed one clear low-blockade peak. In the cases of ATP, ADP and apo, this structural rigidity and compaction is increasingly lost, and Hsp90 undergoes more and more excursions to multiple thermally accessible open conformations. Our results match with previous findings from single-molecule FRET<sup>25,26</sup>, electron microscopy<sup>27</sup> and biochemical results<sup>28</sup>. However, in contrast to these techniques, the NEOtrap measures proteins in a label-free manner, at room temperature, in solution and at the single-molecule level. And unlike single-molecule FRET, the spatial resolution of the NEOtrap is not limited to a certain inter-dye distance (typically  $\leq 9$  nm), and the observation time is not limited by photobleaching. The ability to

distinguish nucleotide-dependent conformations in a protein complex demonstrates the great potential of the NEOtrap, opening the way for studying conformational dynamics in many protein systems.

## Conclusion

In conclusion, we presented the NEOtrap, which is capable of trapping single unmodified proteins irrespective of their net charge for up to hours while resolving their size and conformations with submillisecond time resolution. This represents a drastic increase in single-protein observation time by a factor of  $>10^6$  compared to previous solid-state nanopore studies<sup>19,29</sup>. A DNA-origami sphere was used to induce local EOFs at will, which creates the trapping potential for the NEOtrap. Voltage inversion can be applied as a clean sweep to trap a new protein molecule in the lipid-passivated nanocavity whose size can be freely chosen between a few and 30 nm,



**Fig. 5 | Label-free NEOTrap detection of nucleotide-dependent conformational shifts of the chaperone protein Hsp90.** **a**, Hsp90 in flexible open (left<sup>26</sup>) and compact closed conformations (right, PDB 2cg9, ref. <sup>24</sup>) of the homo-dimeric chaperone protein Hsp90, consisting of an N-terminal nucleotide-binding domain (N), a middle domain (M) and a C-terminal dimerization domain (C). Hsp90 alternates between these conformations depending on different nucleotide-binding states. The inset shows the many stabilizing contacts (yellow dashed lines) formed by the nucleotide in Hsp90's Bergerat-type ATP-binding site. **b**, Representative conductance blockade events caused by single label-free Hsp90 proteins under the specified nucleotide conditions: non-hydrolysable AMP-PNP, ATP, ADP and apo, and decimated to 20 kHz (Methods). Hsp90's conformational heterogeneity increases strongly from top (AMP-PNP, compact) to bottom (apo, flexible), in line with the reduced number of stabilizing nucleotide contacts and with existing Hsp90 literature. **c**, Histograms of 1,000–4,000 such current blockade events (integral normalized to unity) that show a strong nucleotide dependence, indicative of distinct protein conformations that are visualized using grey icons where dark represents frequent and pale represents rare conformations, respectively. These 'all-points histograms' were composed by summing up each datapoint in each event, without previous averaging. Scatterplots of conductance versus dwell time are provided in Supplementary Fig. 8. Less than 1% of outlier events were excluded as described in Supplementary Fig. 9. The results were reproduced in three or more experiments using different measurement orders of the four nucleotide conditions.

that is, in a size range very relevant for proteins and other molecules, nanoparticles, quantum dots and so on. The NEOTrap features a linear mass dependence of the current blockade for globular proteins

of 54 to 340 kDa; can discriminate between different orientations for disc-shaped ClpX, and distinguishes nucleotide-dependent conformations of the chaperone protein Hsp90. The NEOTrap offers a

widely applicable electrical sensing strategy that addresses a central shortcoming of nanopore detection, namely the impractical high-speed translocation that prohibits single-protein characterization. As a result, the NEOtrap has the potential to become a label-free alternative to classical single-molecule techniques such as fluorescence and force spectroscopies. Due to its technical simplicity and multifaceted applicability, we anticipate that the NEOtrap will find wide-spread application in the study of diverse protein systems, and thus evolve into a new tool in the single-molecule toolbox for studying protein dynamics.

### Online content

Any methods, additional references, Nature Research reporting summaries, source data, extended data, supplementary information, acknowledgements, peer review information; details of author contributions and competing interests; and statements of data and code availability are available at <https://doi.org/10.1038/s41565-021-00958-5>.

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## Methods

Glass chips with free-standing 20-nm-thick silicon-nitride membranes were purchased from Goeppert LLC. Nanopores were drilled by TEM as previously described<sup>8</sup>. The chips were rinsed with MilliQ water, ethanol, acetone, isopropanol and plasma cleaned (SPI Supplies) before the assembly in a custom-made polyether ether ketone (PEEK) flowcell with *cis* and *trans* buffer compartments, placed in a Faraday cage, and connected to an Axopatch 200B amplifier and Digidata 1550B digitizer (both Molecular Devices) using Ag/AgCl electrodes (silver wire chloridized in household bleach). All pores were wetted with MilliQ water, before measuring the conductance in 1 M-KHM buffer (1 M KCl, 50 mM HEPES, 5 mM MgCl<sub>2</sub>, pH 7.5) to obtain a measured nanopore diameter as described in Supplementary Notes.

The nanopore passivation procedure was adapted from ref. <sup>19</sup>: POPC (Avanti Lipids) in chloroform was aliquoted in glass vials, dried for 2–4 h in vacuo and either stored at –20 °C (≤1 month) or resuspended to 1 mg ml<sup>–1</sup> in 600 KHM buffer (600 mM KCl, 50 mM Hepes, 5 mM MgCl<sub>2</sub>, pH 7.5) and vortexed for 10 min. Next, 50 µl of lipid suspension were applied to the *cis* compartment of a chip equilibrated in 600 KHM, while applying a continuous zig-zag potential (peaks at ±50 mV, 5 Hz repetition rate). This step was repeated after 5 min. (If no conductance change was observed, 1 M KHM was added to the *trans* compartment to osmotically drive lipid vesicles to the pore.) After a total of 10–15 min incubation time, the flowcell was disconnected and immersed in 0.5 l of MilliQ water for 10–15 min. Next, the (externally dried) flowcell was reconnected to the amplifier and the nanopore flushed with 200 µl of MilliQ water in *cis* and *trans*, followed by the desired measurement buffer. A reproducible one-step conductance drop indicated a stable and consistent coating. About 10% of the pores could not be stably coated in this way and were not used further. Other chips were recoated and reused several times using this procedure.

Nanopore experiments were controlled using ClampEX (Molecular Devices), and data was analysed using self-written code in Igor Pro v.6.37 (Wavemetrics). The open-source NeuroMatic binary file importer was used<sup>30</sup>. Median-aware decimation was used to exploit oversampling: for example, data recorded at 500 kHz sampling were decimated to 5 kHz by replacing each 100 samples by their median value. Threshold criteria were used for event detection. Bootstrapping was used to estimate uncertainties of (1) kinetic rate constants from single-exponential fits and (2) histogram peak positions as specified: for a dataset of size *n*, 10*n* subsets of size *n* were randomly chosen with replacement, individually evaluated and the mean and standard deviations across all subsets are reported. Protein structures were visualized with the PyMOL Molecular Graphics System, v.2.0.6, Schrödinger, LLC. Data and code are available at <https://doi.org/10.5281/zenodo.5059802>.

The DNA-origami sphere was described previously<sup>10</sup>. Details were provided by Tilibit Nanosystems GmbH. Folding reaction mixtures contained 7,560-nt long scaffold DNA at 50 nM final concentration and 227 staple strands at 175 nM final concentration, for each strand. The folding buffer contained 5 mM TRIS, 1 mM EDTA, 5 mM NaCl (pH 8) and 20 mM MgCl<sub>2</sub>. Folding reaction mixtures were subjected to thermal annealing ramps using TETRAD (MJ Research, now Biorad) thermal cycling devices: 15 min at 65 °C, followed by 1-h intervals for each temperature, starting at 64 °C down to 48 °C, decreasing by 1 °C every step. The folding reaction mixtures were then incubated at 20 °C before purification steps. Excess staple strands were removed by ultrafiltration (Amicon Ultra 0.5-ml Ultracel filters, 50 K) with buffer containing 5 mM TRIS, 1 mM EDTA, 5 mM NaCl and 5 mM MgCl<sub>2</sub> (ref. <sup>31</sup>). Then, 0.5 ml of buffer was added to the filter and centrifuged at 10,000g for 3 min at 25 °C. The filtrate was discarded, and 0.1–0.2 ml of the folded object sample and 0.3–0.4 ml of buffer were then added to the filter and centrifuged at 10,000g for 5 min at 25 °C. The filtrate was discarded, and 0.5 ml of buffer were added to the filter and centrifuged at 10,000g for 5 min at 25 °C. The filtrate was discarded again. The third step was repeated three more times. The filter inset was then removed from the tube, placed upside-down into a new tube and centrifuged at 10,000g for 3 min at 25 °C as a final sample retrieving step.

ClpP and ClpX were expressed and purified in-house as previously described<sup>32</sup>. Hsp90 was a gift from B. Hermann and T. Hugel<sup>25</sup>. Avidin (mono-valent SAe1D3, ref. <sup>33</sup>) was a gift from M. Howarth.

Unless stated differently, all measurements were performed under 100 mV bias, 500 kHz sampling, 100 kHz low-pass filter (four-pole internal Bessel filter), in 600 KHM, at room temperature of 21 ± 1 °C. Unless stated otherwise, chemicals and nucleotides were purchased from Sigma Aldrich.

## Data availability

Data are available at <https://doi.org/10.5281/zenodo.5059802>.

## Code availability

Code for data analysis of nanopore recordings as described herein are available at <https://doi.org/10.5281/zenodo.5059802>.

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## Author contributions

S.S. and C.D. conceived the project. S.S. performed all nanopore experiments, analysed the data and purified proteins. H.D. and P.S. advised on DNA origami, and P.S. folded and characterized it. S.S. wrote the paper with C.D. All authors discussed the results and commented on the paper.

## Competing interests

The authors declare no competing interests.

## Additional information

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