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# Towards $\text{Tm}^{2+}$ in Sialon Thin Films Grown by Sputtering Using the Gradient Deposition

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The valence stability parameter ( $E_{\text{FF}}$ ), defined as the difference between the charge transfer energy to the host intrinsic Fermi energy, was used as criterion to analyze the capability of different host materials within the SiAlON class to stabilize divalent thulium. Available data on charge transfer energies and optical bandgap values are reviewed for  $\text{Si}_3\text{N}_4$ ,  $\text{SiO}_2$ ,  $\text{AlN}$ , and  $\text{Al}_2\text{O}_3$ . In addition, new data on thin films, collected by our gradient sputter deposition and characterization method on silicon and aluminum nitrides ( $\text{Si}_{0.75x}\text{Al}_{1-x}\text{N}$ ), are reported. These data are sufficient to show that, at least in the nitride subsection of the SiAlON class, divalent thulium is not expected to be stable due to the presence of high  $E_{\text{FF}}$  values. The use of sub-stoichiometric silicon nitride and oxide is also briefly considered.

## Introduction

Among different candidates of next generation luminescence solar concentrators (LSC) materials for electricity generating glass applications, divalent thulium ( $\text{Tm}^{2+}$ ) doping has been shown to have a high potential. It has been reported to have 65% absorption of the solar spectral range when in di-halide hosts while showing no self-absorption and a relatively neutral color (1). These characteristics are possible because  $\text{Tm}^{2+}$  has a strong absorption from the UV up to 900 nm generated by intense  $4f^{13} \rightarrow 4f^{12}5d^1$  electronic transitions and a single emission line in near infrared (NIR) at 1140 nm credited to the  ${}^2F_{5/2} \rightarrow {}^2F_{7/2}$  transition of the  $4f^{13}$  configuration. However, this specific oxidation state of thulium is rather anomalous: except from  $\text{SrB}_4\text{O}_7$  (2), it was only stabilized so far on di-halide crystals. These crystals are very hygroscopic in nature and therefore not sufficiently attractive for scaling up in industry standard thin film growth techniques. Thus, it is paramount to extend the list of host materials beyond the halides that can host such a promising optically active element with material classes standard to industry. Host materials pertaining to the SiAlON family have the advantage of being wear resistant, and highly stable to temperatures and thermal shock. The possibility of merging a wide absorption range in the UV-VIS with a narrow emission line in NIR with a host material standard to industry shows high scalability potential.

## Rare Earth Valence Stability

When accessing the potential of a given material to host the divalent oxidation state of a given rare earth (RE) ion, one of the accepted and general criteria is given by the  $E_{Ff}$  parameter suggested by Dorenbos in 2005 (3) defined by:

$$E_{Ff} = E^{CT} - E_{F,i} \quad [1]$$

In which  $E^{CT}$  stands for the charge transfer energy of the trivalent ion (4), interpreted as the energy difference ascribed to the effective removal of one electron from the valence band to the ground state of the divalent ion.  $E_{F,i}$  is intrinsic Fermi energy. This relation, essentially borrowed from impurity ionization statistics used in semiconductors, was empirically observed to be a good indicator for RE reduction in solids as well. Without tackling the problem from a first-principles perspective, Dorenbos correlated valence stability of the lanthanides – i.e., the possibility of observing the divalent form of a given lanthanide ion – over more than 300 crystalline compounds to the distance of the lanthanide ground state within the host bandgap and its  $E_{F,i}$  (estimated to be at the middle of the valence and conduction bands). In his work, the intrinsic Fermi energy is determined from the host exciton energy through  $E_{F,i} = 0.54E^{ex}$ .

The main observation in this set of compounds is that a necessary condition for which  $\text{Eu}^{2+}$  can be found is  $E_{Ff} < 0.7$  eV. In this situation, it is reported that  $\text{Eu}^{2+}$  can be obtained even under oxidizing conditions.  $E_{Ff} < 0.7$  eV will be referred to from now on as “Dorenbos’ criterion”. However, there is no apparent sharp line dividing the situations where a lanthanide ion can be found in a divalent or in a trivalent state. Moreover, if  $E_{Ff} > 1.2$  eV, it was reported that “a complete reduction to  $\text{Eu}^{2+}$  is not possible anymore”. This does not necessarily mean that  $\text{Eu}^{2+}$  cannot be found but that it requires extreme reducing treatments – which can also mean that such a condition might not be stable. Similar arguments can be raised for the case of the next ions most easily reduced (in order they are Yb, Sm, and Tm). Furthermore, this analysis does not allow for a complete quantification of the mixture of oxidation states or the expected average oxidation state.

One of the most relevant tools developed for the interpretation of lanthanide luminescence spectroscopy over the last one and a half decade was the emergence of band diagrams which included the lanthanide  $4f$  energy levels (5,6). Those diagrams use optical spectroscopy data such as photoluminescence emission and excitation and sometimes absorption spectroscopy to determine charge transfer energies (4) and optical bandgap information which can then be ascribed, in a simplified picture of the electronic structure of solids, to the RE  $4f$  and  $5d$  (7) levels in both divalent and trivalent oxidation states. This information can also be correlated either relatively to the host valence band edge or to the vacuum energy. Diagrams obtained in such a way are often referred to as host referred binding energy (HRBE) diagrams or vacuum referred binding energy (VRBE) diagrams.

Thus, a first attempt of obtaining criteria for searching possible material candidates hosting  $\text{Tm}^{2+}$  one could use Dorenbos’ criterion on available data regarding the simplest material options within the SiAlON class like  $\text{SiO}_2$ ,  $\text{Si}_3\text{N}_4$ , AlN and  $\text{Al}_2\text{O}_3$ . One of the advantages of the VRBE schemes, or more specifically, the zig-zag shape of the

lanthanide energy levels, is that optical spectroscopy information about one ion of the RE family can offer information about what are the expected charge transfer,  $5d$ , and excitation energies of a different RE ion in the same material. Therefore, information on  $\text{Eu}^{2+}$  and  $\text{Sm}^{3+}$  charge transfer and  $4f \rightarrow 5d$  transition energies are already enough for evaluating if the same compound can stably host  $\text{Tm}^{2+}$ .

Without the aim of being an exhaustive review on the topic of RE reduction in solid compounds, this paper contributes by discussing relevant literature and the progress made so far towards unravelling conditions for obtaining the rather anomalous  $\text{Tm}^{2+}$  in hosts pertaining to the SiAlON class. This is undertaken through the deliberate choice of Dorenbos' criterion to search and analyze for promising compositions and, thus, optical properties.

In the next section, a survey of a selection of available data of the materials from the corners of the SiAlON phase diagram, i.e.,  $\text{SiO}_2$ ,  $\text{Si}_3\text{N}_4$ ,  $\text{Al}_2\text{O}_3$ , and  $\text{AlN}$ , are reviewed for the presence of  $\text{Eu}^{2+}$  and  $\text{Sm}^{2+}$  and correlated with Dorenbos' criterion. This data is then correlated with the ability to host  $\text{Tm}^{2+}$  and with the expected effect of sub stoichiometry on the promotion of divalent formation.

This discussion is followed by experimental results on the feasibility for obtaining  $\text{Tm}^{2+}$  in the nitride composition region of the SiAlON class.  $\text{Tm}^{3+}$  excitation and emission photoluminescence were correlated with Dorenbos' criterion for a wide ratio of Al/Si. The films were prepared by combinatorial RF/DC magnetron sputtering and characterized as a function of position on the substrate through the gradient method (8).

Finally, a small section is dedicated to exploring the implications of recent results available on the VRBE diagrams of sub-stoichiometric  $\text{SiO}_2$  and  $\text{Si}_3\text{N}_4$  in the promotion of  $\text{Tm}^{2+}$ .

### Optical Spectroscopy data on Eu and Sm

In TABLE I there is a collection of available data on charge transfer ( $E^{\text{CT}}$ ), exciton energy ( $E^{\text{ex}}$ ), and the estimated  $E_{\text{FF}}$  parameters for Eu, Sm, and Tm doped in  $\text{Si}_3\text{N}_4$ ,  $\text{AlN}$ ,  $\text{SiO}_2$  and  $\text{Al}_2\text{O}_3$ .  $E_{\text{FF}}$  is more easily accessible for Eu than for the other two dopants. For Sm and Tm,  $E_{\text{FF}}$  can be easily estimated by shifting these energies by 1.2 eV for  $\text{Sm}^{2+}$ , and 1.7 eV for Tm (9). The results contained in TABLE I are schematically gathered in Figure 1.

There is no direct lanthanide spectroscopic data on the CT energies for  $\text{SiO}_2$ . But a reasonable estimate can be done looking at the trends found in different oxide compounds. The  $E^{\text{CT}}$  depends, in a first approximation, on the binding strength and distance experienced by the lanthanide to the anions it is bonded to (10). Additionally, it increases with  $E^{\text{ex}}$ . In the LiRESiO<sub>3</sub> series of elements, oxygen is bonded in  $\text{SiO}_4^{4-}$  tetrahedra which can resemble the neighborhood expected for  $\text{Eu}^{3+}$  in  $\text{SiO}_2$ . The compound with highest  $E^{\text{ex}}$  is  $\text{LiLuSiO}_4$ , thus, a reasonable estimate on  $\text{Eu}^{3+}$  CT energy would lie on similar levels found for  $\text{LiLuSiO}_4:\text{Eu}$ , i.e.,  $\sim 5.5$  eV. There is direct data on the  $\text{Eu}^{3+}$  CT in  $\text{Al}_2\text{O}_3$ . It has been reported to be 3.8 eV (11). Characterization of bandgap energies from XPS results on different  $\text{Al}_2\text{O}_3$  phases gives 8.1 eV for  $\alpha\text{-Al}_2\text{O}_3$ , 7.0 eV for  $\gamma\text{-Al}_2\text{O}_3$ , and 6.5 eV in amorphous  $\text{Al}_2\text{O}_3$  (12,13). Using the assumption that the charge transfer energy

does not strongly change in alumina polymorphs (14) one can estimate the different  $E_{FF}$  values for  $Al_2O_3$ . The lowest  $E_{FF}$  value for Eu, which can be correlated to the higher chances of reducing a RE ion, occurs for  $\alpha$ - $Al_2O_3$ . In fact,  $Eu^{2+}$  can be observed in  $\alpha$ - $Al_2O_3$  (15) but it is not commonly observed in amorphous  $SiO_2$  (16,17). Even in the case of higher  $E_{FF}$  values,  $Eu^{2+}$  can still be observed in SiAlONs (28-20) and in  $\gamma$ - $Al_2O_3$  (11). In all cases, Dorenbos' criterion is satisfied. Despite the high  $E_{FF}$  value retrieved for Sm in  $\gamma$ - $Al_2O_3$ ,  $Sm^{2+}$  can still be observed in  $\gamma$ - $Al_2O_3$  (21).

In the nitrides, the charge transfer energies of  $Eu^{3+}$  are relatively close to each other (10). For the same reasons discussed above,  $E^{CT}$  in  $Si_3N_4$  should be lower than the value known for AlN  $\sim 3.5$  eV (22). This can be estimated to be  $\sim 3.2$  eV, similar to what has been found in GaN (23). So,  $E_{FF}$  are respectively at 0.6 and 0.1 eV, meeting Dorenbos' criterion in both cases. However, Eu tends to be trivalent in AlN (22) but can also be found in divalent form (24). In  $Si_3N_4$  (25) Eu is known to have a mixed valence.  $Eu^{2+}$  characteristic luminescence is known to be stable in a- $Si_3N_4$  even after post deposition thermal treatments on oxidizing conditions. As previously mentioned, Dorenbos' valence stability model is an empirical relationship which can particularly fail in situations where  $E_{FF}$  is around 0.7 eV, e.g., when Eu is doped in the  $Si_3N_4$  (25) or AlN (22). Moreover, the inherent error bars associated with positioning the energy levels within the forbidden bandgap (26) should be taken into account and, very probably, can be an additional limiting factor to the inefficiency in the method to correctly predict the presence of  $Eu^{2+}$ .

**TABLE I.** Charge transfer ( $E^{CT}$ ), exciton energies ( $E^{ex}$ ), and the estimated  $E_{FF}$  parameter for Eu, Sm, and Tm in the simple materials contained in the SiAlON class.

Host	RE	$E^{CT}(3+)$ [eV]	$E^{ex}$ [eV]	$E_{FF}$ [eV]
$Si_3N_4$	Eu	3.2 <sup>b</sup>	4.8 <sup>a</sup>	0.6
	Sm	--		1.8
	Tm	--		2.3
AlN	Eu	3.5 <sup>d</sup>	6.2 <sup>d</sup>	0.1
	Sm	--		1.3
	Tm	--		1.8
$SiO_2$	Eu	5.5 <sup>j</sup>	8.7 <sup>i</sup>	0.8
	Sm	--		2.0
	Tm	--		2.5
$\alpha$ - $Al_2O_3$	Eu	3.8 <sup>e</sup>	8.1 <sup>c</sup>	-0.6
	Sm	5.0 <sup>f</sup>		0.6
	Tm	--		1.1
$\gamma$ - $Al_2O_3$	Eu	3.8 <sup>e</sup>	7.0 <sup>g</sup>	0.0
	Sm	5.0 <sup>f</sup>		1.2
	Tm	--		1.7
Amorp. $Al_2O_3$	Eu	3.8 <sup>e</sup>		0.3
	Sm	5.0 <sup>f</sup>	6.5 <sup>h</sup>	1.5
	Tm	--		2.0

<sup>a</sup> Estimated using the optical bandgap value of 5.2 eV (27) by  $E^{ex} = \frac{E_g}{1.08}$ ; <sup>b</sup> Estimated from GaN (28); <sup>c</sup> Estimated using bandgap value of 8.7 eV (12) - more information on the text; <sup>d</sup> From (28); <sup>e</sup> Estimated using excitation spectra results from (11). An assumption about the CT energy not changing drastically for all alumina polymorphs is made. <sup>f</sup> Estimated using excitation spectra results of (21); <sup>g</sup> Estimated using XPS results from  $E_g = 7.6$  eV (18); <sup>h</sup> Estimated using XPS results from  $E_g = 7.0$  eV (13); <sup>i</sup> From (29); <sup>j</sup> Estimated from LiRESiO<sub>4</sub> series as suggested in (10) - more information is in the text.

As previously mentioned, the higher the bandgap and the lower the charge transfer energy of the trivalent ion, the better the chances of a candidate material to stably host

$\text{Tm}^{2+}$ . As suggested from these results, the simplest compounds, within the SiAlON phase diagram, do not show the necessary conditions for hosting  $\text{Tm}^{2+}$ , i.e., a high optical bandgap and a low charge transfer energy of the trivalent ion. For all materials, the  $E_{\text{FF}}$  values expected for Tm are well above 1.0 eV. Among those options, the best one would be  $\alpha\text{-Al}_2\text{O}_3$  with an  $E_{\text{FF}}$  of 1.1 eV which is still 0.4 eV above the Dorenbos' criterium of 0.7 eV.

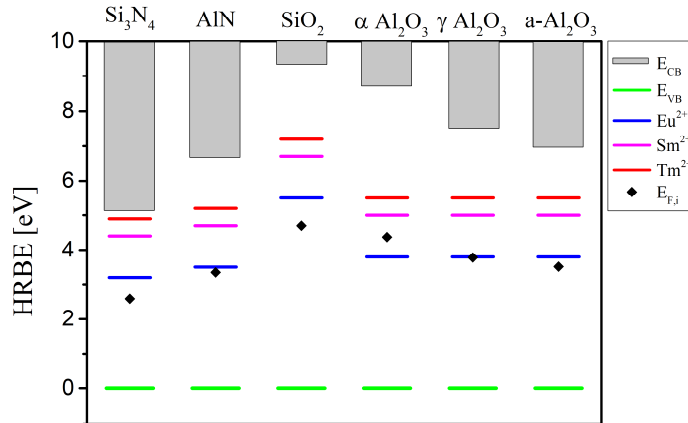


Figure 1. Host referred binding energy (HRBE) diagrams of the simple compounds within the SiAlON class constructed from values in TABLE I. The lines in blue, pink and red represent the ground state energy levels occupied respectively in  $\text{Eu}^{2+}$ ,  $\text{Sm}^{2+}$  and  $\text{Tm}^{2+}$ . The green lines represent the valence band edge location and the grey rectangles the location of the conduction bands.

### Valence stability of Tm in $\text{SiAlN}_x\text{:Tm}^{3+}$

In order to directly access the influence of the Al/Si ratio in silicon and aluminum nitrides ( $\text{Si}_{0.75x}\text{Al}_{1-x}\text{N}$  for short) on the stabilization of  $\text{Tm}^{2+}$ , the gradient characterization method (8) was used to relate local (that is position dependent) film composition – determined by energy dispersive x-ray spectroscopy (EDX) – with local transmission, photoluminescence emission and excitation data. This unique approach made it possible to retrieve a wide composition range in just a handful of different films.

### Experimental

Amorphous thin films of Tm doped in  $\text{Si}_{0.75x}\text{Al}_{1-x}\text{N}$  were prepared by combinatorial magnetron RF/DC sputtering. Four films were deposited on square  $50 \times 50 \text{ mm}^2$  UV grade quartz substrates within an AJA ATC Orion 5 magnetron sputtering system with a base pressure of  $1.5 \times 10^{-7}$  torr. Prior to the deposition, each substrate had been cleaned by rinsing trice with DI water and ethanol. The deposition was carried out with 5.08 cm diameter metal Tm (99.99%, Demaco), Si (99.999%, Lesker), and Al (99.9995%, Lesker) targets which were reactively co-sputtered. The process total gas flow was kept fixed at 32 sccm. 6N purity Ar and 5N purity  $\text{N}_2$  gas flows were chosen to meet the requirements in stoichiometry and deposition rate. TABLE II shows the deposition parameters used for the fabrication of each film. The working pressure was  $3 \times 10^{-3}$  torr and the samples were not actively heated.

Transmission, photoluminescence emission and excitation were carried out with the methods described in detail elsewhere (8). Composition characterization were carried out from SEM-EDS collected from a JEOL JSM-IT100 operated at 15 kV with probing current at 50% for 2 min per measurement. Elemental composition was quantified at 1000 $\times$  magnification. Low vacuum mode (at 35 Pa) was used for to enhance quantitative elemental analysis without the need for a conductive coating. The signal and elemental composition was then compared with simulations carried out with DTSA II software for additional consistency considering the signal from the quartz substrates was also present.

**TABLE II.** Deposition conditions used for each film.

Film	Power Employed			Gas Flow	
	Si [W]	Al [W]	Tm [W]	Ar [sccm]	N <sub>2</sub> [sccm]
Si <sub>3</sub> N <sub>4</sub>	130	0	10	27	5
AlN	0	110	25	22	10
Library 1	110	60	20	22	10
Library 2	90	64	30	22	10

## Results

Figure 2a shows a top view perspective of the metallic precursors used for film fabrication. Figure 2b shows the appearance of the library 1 film right after deposition. Figure 2c shows a surface map of the expected local composition of library 1. Interpolation was carried out using the surface-source evaporation equation (8) to extend the  $[Al]/([Al]+[Si])$  ratio over all substrate – shown as a “heat” map between 0 (blue) and 1 (yellow). Figure 2d shows a surface map representing the distribution of the local thickness obtained from transmission in VIS region. Figure 2e is the typical photoluminescence emission spectrum of Tm<sup>3+</sup> which were used to generate the excitation plots shown in Figure 3.

Figure 3a shows the excitation spectrum obtained from the Si<sub>3</sub>N<sub>4</sub>:Tm thin film. Figure 3b is the excitation spectrum upon one point at library 2. Figure 3c shows the excitation of AlN:Tm. Eu<sup>3+</sup> CT energy in Si<sub>3</sub>N<sub>4</sub> is known to be at 3.2 eV. This results in an approximate position of Tm<sup>3+</sup> CT energy to be at 4.9 eV. The optical bandgap ( $E^{04}$  definition) of Si<sub>3</sub>N<sub>4</sub> is around 5.2 eV (30). Thus, photoluminescence excitation curve fitted to a single Gaussian (Figure 3a) can be interpreted as the convolution of the CT and exciton excitation. The CT energy of AlN:Tm have already been characterized before to the band centered at  $\sim$ 4.5 eV (31). The inset shows the transition assigned to AlN exciton at 6.2 eV. Similar excitation spectra were collected in various locations of libraries 1 and 2 to directly access the influence of  $[Al]/([Al]+[Si])$  on the Tm<sup>3+</sup> CT energy. The graph shown in Figure 3d evidences that the CT band energy in the nitrides hardly changes with the Al/Si ratio remaining close to 4.8 eV.

Using values from TABLE I it is possible to estimate  $E_{Ff}$  to have a very high value for AlN ( $E_{Ff} \equiv E^{CT} - 0.54E^{ex} \approx 1.3$  eV) and for Si<sub>3</sub>N<sub>4</sub> ( $E_{Ff} = 2.1$  eV). In the Si<sub>0.75x</sub>Al<sub>1-x</sub>N alloy,  $E_{Ff}$  is expected to be even higher due to a bandgap contraction with higher Si content and a very similar  $E^{CT}$ . This indicates that obtaining Tm<sup>2+</sup> in this region of the SiAlON phase space might indeed not be possible.

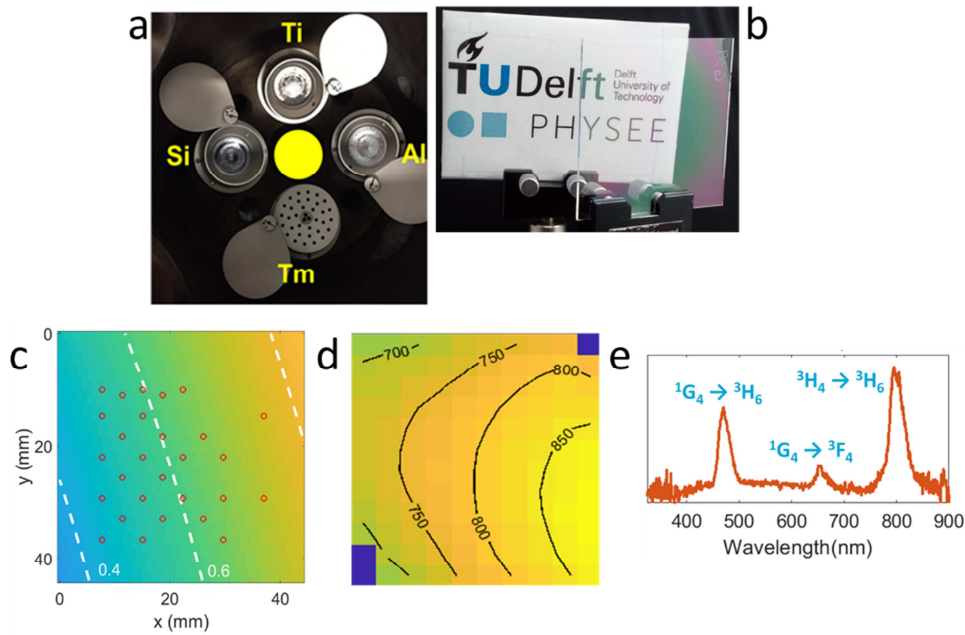


Figure 2. a) Top view of the 4 sputtering guns employed in the combinatorial sputtering system. The yellow circle indicates the position of the substrate relative to the guns. Ti gun was used only for assuring vacuum quality. A metallic “mask” is used over Tm gun to limit its deposition rate; b) Picture of library 1; c) Distribution of the relative concentration of Al over Al and Si -  $[Al]/([Al]+[Si])$  - as obtained by EDS analysis (red points indicate measured positions; color map is the result of evaporation curve fitting method (8,32)); d) Thickness (in nm) distribution of library 1 as obtained by transmission analysis; e) Typical  $Tm^{3+}$  emission spectrum obtained in  $SiAl_{1-x}N_x$  thin gradient films showing the transitions  $^1G_4 \rightarrow ^3H_6$ ,  $^3H_4 \rightarrow ^3H_6$ , and  $^1G_4 \rightarrow ^3F_4$ .

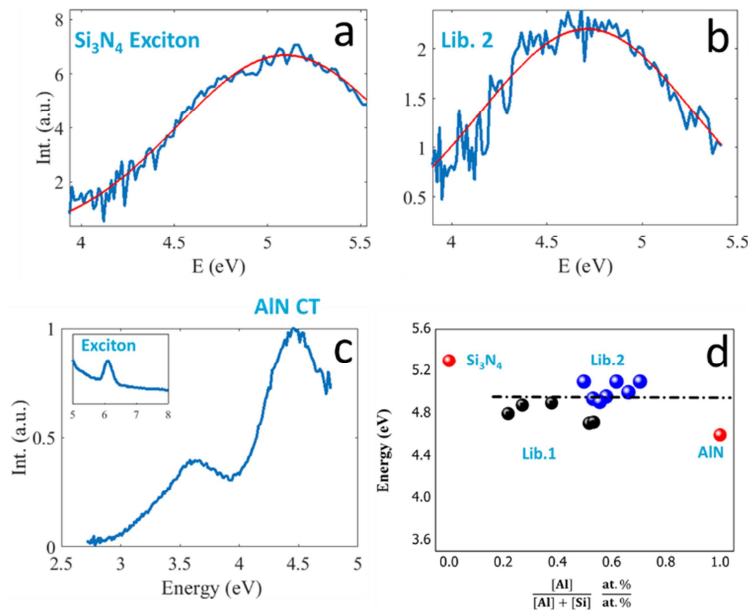


Figure 3. a) Room temperature excitation spectra (blue) and Gaussian fit (red) of the  $Tm^{3+} \ ^3H_4 \rightarrow \ ^3H_6$  line in  $Tm^{3+}$  doped  $Si_3N_4$ ; b) PLE spectra in one point in library 2; c) and in AIN thin films; d) PLE peak energy (correlated with Tm CT energy) vs.  $[Al]/([Al]+[Si])$  obtained from a selected regions in a series of gradient films.

## Sub-stoichiometry as strategy of Tm reduction

Rather than tuning only the cation by replacing Si for Al, a different strategy would be to study the influence of the O or N anion ratio. Since the charge transfer energy of  $\text{RE}^{3+}$  depends on the energy difference between the valence band edge and the  $\text{RE}^{2+}$  energy level, a change in the oxygen (or nitrogen) deficiency in the host material can alter the position of the valence band edge while maintaining the RE energy levels relatively unchanged (4). This lowers the required charge transfer energy offering better chances of a given material for hosting  $\text{Tm}^{2+}$  because this effectively raises the Fermi energy moving it closer to the 2+ ground state in the VRBE diagram.

In this section, this possibility is shortly discussed because tight-binding calculations results on the electronic structure, as a function of oxygen (or nitrogen) content in sub stoichiometric  $\text{Si}_3\text{N}_4$  (27) and  $\text{SiO}_2$  (32), have become recently available.

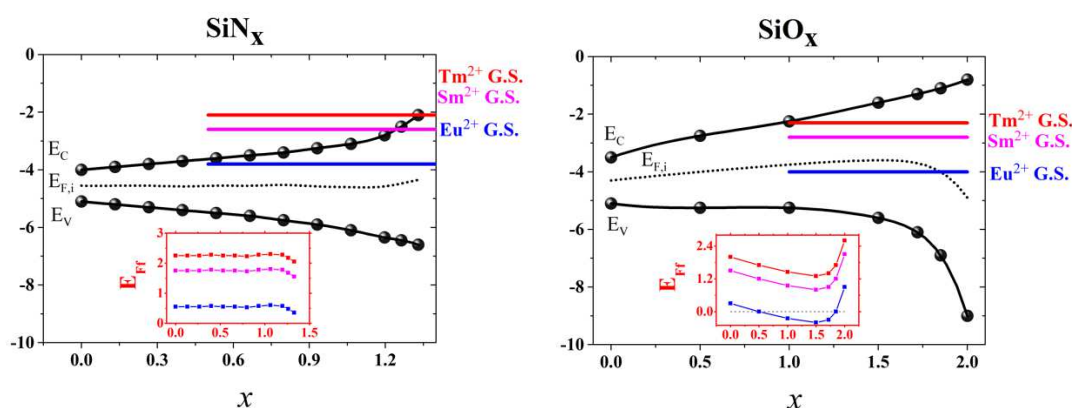


Figure 4. Effect of the nitrogen – or oxygen– ratio over Si on the optical bandgap calculated by a tight binding method in  $\text{SiO}_x$  (32) and  $\text{SiN}_x$  (27). The dashed lines indicate the expected intrinsic Fermi energy located at the midpoint of the bandgap. The insets show explicitly the dependence of the  $E_{Ff}$  parameter upon  $x$ . The colored lines indicate the ground state levels of  $\text{Eu}^{2+}$  (in blue),  $\text{Sm}^{2+}$  (in pink), and  $\text{Tm}^{2+}$  (in red). These positions are placed on the bases of what is known for  $\text{Si}_3\text{N}_4$  and  $\text{SiO}_2$  and are assumed to not change over stoichiometry.

Figure 4 is an adapted plot of the trends in the VRBE of  $E_V$  and  $E_C$  band edges as a function of the stoichiometry of those compounds. The same references used in TABLE I for  $\text{Si}_3\text{N}_4$  and  $\text{SiO}_2$  are used to place the VRBEs of the ground state levels related to  $\text{Eu}^{2+}$ ,  $\text{Sm}^{2+}$ , and  $\text{Tm}^{2+}$ . Because the chemical shift (which determines the VRBE of the RE energy levels) depends, in a first approximation, only on the first shell of anions bonded to the RE (14), the dependence of these levels with the stoichiometry is not strong, so it seems a good approximation to fix these energies across the stoichiometry. The inset of each graphic shows the estimated  $E_{Ff}$  values as a function of stoichiometry. For the sub-nitrides ( $\text{SiN}_x$ ),  $E_{Ff}$  is always positive for all the ions showing a minimum value in the stoichiometric  $\text{Si}_3\text{N}_4$ , so the sub-stoichiometry strategy does not have a positive result in RE reduction. For the sub-oxides ( $\text{SiO}_x$ ), however, there is a minimum in the  $E_{Ff}$  value close to  $x = 1.5$ . In this situation,  $E_{Ff}$  is expected to be negative for Eu,  $\sim 0.8$  eV for Sm, and  $\sim 1.3$  eV for Tm. It is known that  $\text{Sm}^{2+}$  is not stable in  $\text{SiO}_2$  but photoluminescence attributed to it can be observed in  $\text{SiO}_{1.2}$  after annealing on oxidizing atmosphere (33). This is also consistent with recent observations with the raise in  $\text{Sm}^{2+}$  concentration, determined by XPS, in near-stoichiometric Sm doped  $\text{SiO}_2:\text{Al}$  materials (34). The lowest

$E_{\text{FF}}$  value experienced by  $\text{Tm}^{2+}$  seems already too high for any measurable amount of  $\text{Tm}^{2+}$ . Thus, this strategy alone might not be enough for the formation of  $\text{Tm}^{2+}$ .

## Conclusion

The work presented in this paper shows the current status of an ongoing effort on the determination of a suitable host material capable of stably hosting  $\text{Tm}^{2+}$  while still being compatible with current installed industry infrastructure – in particular, this paper presented a  $\text{Tm}^{2+}$  valence stability study focused on materials within the SiAlON class. An attempt of reviewing relevant data for defining new strategies of preparation of these materials was delivered and an analysis was given in terms on what have been referred here as the Dorenbos' criterion.

According to the available data on  $\text{Si}_3\text{N}_4$ ,  $\text{SiO}_2$ ,  $\text{AlN}$ , and  $\text{Al}_2\text{O}_3$ , mixed valence Eu and Sm are expected and observed while the conditions for  $\text{Tm}^{2+}$  formation is not met. Composition and optical characterizations reported on gradient  $\text{Tm}^{3+}$  doped  $\text{Si}_{0.75x}\text{Al}_{1-x}\text{N}$  thin films also suggest that the substitution of Si by Al does not offer the necessary conditions for  $\text{Tm}^{2+}$  stabilization. Other strategies, such as the fabrication of substoichiometric materials seem to improve the probability of stabilizing  $\text{Tm}^{2+}$  but at the expense of lowering severely the material bandgap.

Research on the oxides is still under way and might hold a suitable candidate. As suggested by the results collected for  $\text{SiO}_2$  and  $\text{Al}_2\text{O}_3$ , the replacement of Si by Al can have a more pronounced effect on the CT energy of  $\text{Tm}^{3+}$  favoring  $\text{Tm}^{2+}$  stabilization.

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