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Vertically tapered layers for optical applications fabricated using resist reflow

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Abstract
This paper reports on the IC-compatible fabrication of vertically tapered optical layers for use in linear variable optical filters (LVOF). The taper angle is fully defined by a mask design. Only one masked lithography step is required for defining strips in a photoresist with trenches etched therein of a density varying along the length of the strip. In a subsequent reflow, this patterned photoresist is planarized, resulting in a strip with a local thickness defined by the initial layer thickness and the trench density at that position before reflow. Hence a taper can be flexibly programmed by the mask design to be from 0.001° to 0.1°, which enables the simultaneous fabrication of tapered layers of different taper angles. The 3D pattern of resist structures is subsequently transferred into Si or SiO2 by appropriate etching. Complete LVOF fabrication involves CMOS-compatible deposition of a lower dielectric mirror using a stack of dielectrics on the wafer, tapered layer formation and deposition of the top dielectric mirror. Design principle, processing and simulation results plus experimental validation of the technique on the profile in the resist and after transfer of the taper into Si and SiO2 are presented.

(Some figures in this article are in colour only in the electronic version)

1. Introduction

Single-chip optical micro-spectrometers have huge potential in many applications, such as identification of bio-molecules and in chemical analysis [1]. A small linear variable optical filter (LVOF) integrated with an array of optical detectors is a very suitable candidate for a high-resolution micro-spectrometer [2]. The LVOF is basically a one-dimensional array of many Fabry–Perot (FP) type of optical resonators. Rather than a huge number of discrete devices, the LVOF has a center layer (the resonator cavity) in the shape of a strip and a thickness that changes over its length. Dielectric mirrors are on either side. The spectral resolution of a FP resonator is determined by surface flatness, mirror parallelism and mirror reflectivity. The spectral range of a FP type of LVOF is determined by the thickness variation of the cavity layer over the length of the strip and has to be in the order of a quarter of the wavelength and very well controlled, which makes the fabrication of miniature LVOFs a technological challenge. Conventional approaches for fabrication of tapered optical layers are based on gray tone lithography [3], shadow masking [4], moving mask [5] and polishing.

This paper reports on the IC-compatible fabrication of a vertically tapered layer with a very small taper, using only one lithography step for defining strips in the photoresist and the local trench density therein, as shown schematically in figure 1. The trench density defines the local amount of material removed (figure 1(a)) and is followed by a reflow of this patterned photoresist to locally planarize the remaining strip of material (figure 1(b)). The result is an effective reduction of the resist layer thickness by a value defined by localized trench density. Hence a taper can be flexibly programmed by a mask design to be from 0.001° to 0.1°, which enables the simultaneous fabrication of tapered layers of different taper angles. The structure can also be subsequently transferred into a Si or SiO2 layer by appropriate etching and can be used as a cavity layer of a FP type of LVOF. An LVOF results in case the taper is fabricated in a layer that is deposited on top of a dielectric mirror and another dielectric mirror is deposited...
on top of the tapered layer. The flexibility in the slope fabrication makes it possible to use such a tapered layer also for coupling of optical fibers into optical planar wave guides [3, 6].

A resolution of $\Delta \lambda = 0.5$ nm in an LVOF spectrometer in the entire visible spectral range (400 nm–700 nm) requires about 600 channels. When assuming the minimum feature size of the process enables a pitch between photo detectors about 600 channels. When assuming the minimum feature in the entire visible spectral range (400 nm–700 nm) requires [3, 6].

for coupling of optical fibers into optical planar wave guides fabrication makes it possible to use such a tapered layer also on top of the tapered layer. The flexibility in the slope approach can be implemented using a tapered central layer in which the thickness varies linearly with position. However, an extremely small slope is required: $\psi = \arctan(\Delta z/L) = 0.1/3000 \approx 0.004^\circ$. The actual slope should be designed to match the wavelength range of the intended application and the resolution and the pitch of available detector array.

The photoresist reflows when heated above its glass-transition temperature. The resist reflow is mainly used for surface planarization and as a low-cost IC-compatible technique for the fabrication of micro-lenses [7–9]. The resist reflow technique is also used in the fabrication of sub-micron contact holes [10, 11]. Recently, resist reflow has also been investigated for IC-compatible fabrication of tapered structures with a relatively large taper ( $\Delta z/L = 10 \mu m/1$ mm) for optical coupling [3, 6]. The actual surface profile of the reflowed layer is not critical in the fabrication of sub-micron contact holes [11]. Moreover, the ‘droplet’ shape used in lenses fabricated using reflow is a direct result of the reflow mechanism. However, the shape of a tapered layer for use in an LVOF should be highly linear and requires a more detailed investigation of the mechanisms involved in the reflow process. A more detailed understanding of this mechanism is crucial for assessing the possible angles of the taper and for minimizing the surface roughness. Moreover, an accurate prediction of the resist profile after reflow avoids test runs, hence decreases the processing time and reduces the cost.

2. Resist reflow process

The viscosity of the resist decreases rapidly at temperatures above the glass transition temperature [12]. The material flows because of surface tension, while remaining coherent.

The flow of such a low-viscosity material across a substrate is described by the Navier–Stokes equation:

$$\rho \left( \frac{\partial \vec{V}}{\partial t} + \vec{V} \cdot \nabla \vec{V} \right) = -\nabla p + \mu \nabla^2 \vec{V} + \vec{f}, \quad (1)$$

where $\vec{V}$ denotes the flow velocity, $\nabla \vec{V}$ is tensor derivative of $\vec{V}$, $\rho$ is the fluid density, $p$ is the pressure, $\mu$ is the fluid viscosity and $f$ is external forces on the fluid. The final shape of the resist after reflow does depend on the initial shape of the resist, however, much more significantly on the wetting angle. This is a key thermodynamic variable and defines the angle between the resist sidewall and the substrate after reflow. This wetting angle depends on the interfacial tensions of the surfaces, but not on the initial shape [13]. Since the wetting angle depends on the thermodynamic properties of the two surfaces involved, there is little control apart from the choice of materials and reflow temperature.

Control of the profile of the tapered layers is possible by patterning of the photoresist to define both the outer contour of the rectangular initial layer to be reflowed, but also for the fine definition of the inner area to ensure highly localized reflow only. A detailed analysis of the reflow process in such a structure is required for predicting the final profile, which can subsequently be used for designing the mask that is to be used for the definition of the fine patterning of the layer.

Photoresists AZ9260, AZ4562 and SPR3012 have been used in this work. The glass transition temperature is at 140 °C for AZ9260 and at 120 °C for both AZ4562 and SPR3012. After spin coating AZ9260 and AZ4562 result in layers of about 5 μm thickness and SPR 3012 of about 1 μm thickness. Reflow of the photoresist takes place when its viscosity drops at temperatures beyond the glass transition temperature or when exposed to a solvent.

Cross-linking of the resist is avoided when carrying out the reflow in a chamber saturated with the solvent vapor. This approach for reflow is referred to as chemical reflow.

Figure 1. (a) Trench density variation in the resist after lithography; (b) tapered profile of the resist after reflow of the pattern.
2.1. Thermal reflow

Reflow of the photoresist takes place at temperatures beyond the glass transition temperature, at which the resist becomes soft and starts flowing on the substrate due to surface tension [7]. This approach for reflow is referred to as thermal reflow. The lateral flow of the resist depends on the contact (wetting) angle between the substrate and the resist. A smaller contact angle implies a higher degree of wettability and more flow of the resist on the substrate. Although the glass transition temperature of the resist has a relatively low value, the resist does not flow well over the substrate due to its high viscosity. Therefore, the resist has to be kept at elevated temperatures for a long period of time to ensure sufficient flow over the substrate for a smooth surface. However, the elevated temperature also starts cross linking of the resist, resulting in a hardened material, which does not flow.

Figure 2 shows the shapes of resist strips with different widths before and after thermal reflow. The width of strips varies from 60 μm to 300 μm. The gap between the strips is 20 μm to ensure that the reflowed strips would not join. The contact angle \( \theta \) between the patterned resist and the substrate is 48° before reflow and 17.2° after reflow in all the structures.

This initial experiment demonstrates an edge effect which is limited to about 50 μm from the circumference and is not affected by the width of the strip. The edge effect occurs in strips wider than 72 μm. This property determines the maximum width of the strips that can planarize gaps.

2.2. Chemical (solvent) reflow

Cross-linking of the resist is avoided when carrying out the reflow in a chamber saturated with the solvent vapor. This approach for reflow is referred to as chemical reflow. The wafers are put inside a beaker of the solvent (in the case of the resists used, in PGMEA) with a cover on top. The beaker is heated up to 40–60 °C and the vapor of the solvent is absorbed by the patterned resist. The result is a very soft substance, which flows easily over the substrate. Chemical (solvent) reflow is a very fast reflow technology. However, the main problem in this technique is the uncontrollability of condensation of the solvent vapor on the substrate. When this takes place all structures are flushed away. Nevertheless, this technique is suitable for resist with high viscosity, especially if the resist cross links easily at the glass transition temperature (like AZ4562). The probability of condensation is significantly reduced when the substrate is kept at a higher temperature compared to that of the vapor inside the beaker.

2.3. Chemical–thermal reflow

In the combined approach the wafer is alternatingly exposed for short time to the solvent, so that the resist absorbs some of the solvent vapor and is subsequently exposed to a temperature above glass transition temperature until most of the solvent is removed. Cross-linking of the resist is avoided and a better control of the quality of the surface is obtained. Heating time and solvent exposure time are optimized based on the resist in use and the desired slope angle (pattern) for minimum surface roughness.

Figure 3 shows the same structure as shown in figure 2 after a chemical–thermal reflow process. The sample is exposed to PGMEA vapor at 40 °C for 5 min and subsequently heated on a hotplate at 120 °C for 15 min. Comparison between figures 2 and 3 shows that the wettability is much higher in chemical–thermal reflow. The contact angle between the resist and the substrate is 8° and resist structures smaller than 60 μm
start spreading considerably over the substrate. Considering this fact most structures in the mask are designed with strips not more than 60 μm. This gives a requirement for the minimum trench density.

The edge effect happens in strips wider than 162 μm compared to 72 μm in the case of thermal reflow. This implies that wider resist strips can planarize patterns with lower trench densities and therefore smaller taper angles can be fabricated. Figure 4 also shows the comparison between thermal and chemical thermal reflow for a 50 μm wide resist strip. It can be noted that the resist has smeared beyond the initial contact line with the substrate, which is very important for planarization to fill up the gaps between the strips.

Although the experimental results give good information about wettability and the shape of resist strips after reflow, a more detailed understanding of the reflow process is required to actually solve these problems and to design a practical process for fabrication of the LVOF. The modeling required to achieve that is discussed in the next section.

3. Finite element simulation

The resist profile after reflow is simulated using finite element method software COMSOL.

There are two methods for simulating objects movement in COMSOL. One is based on the moving mesh method [14] and the other is the level set function method. The moving mesh is more suitable for small motions of an object.

Considering the surface tension requires the calculation of curvature of deformation, which is mathematically more difficult. To include surface tension forces in the Navier–Stokes equations, the curvature of the profile has to be calculated at each time step. A border is assumed between the resist (as fluid 1) and air (as fluid 2) in the level set method, by defining a level set function ϕ. The value of the level set function changes from 0 to 1 between the two fluids. The border is thus represented as the 0.5 contour of the level set function. The transport of the border is given by

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \nabla V) + \gamma \times \left[ \left( \nabla \cdot \left( \phi \nabla \phi \right) \right) - \varepsilon \nabla \cdot \nabla \phi \right] = 0,$$

(2)

where ϕ is the level set function and γ and ε are re-initialization and interface thickness control parameters, respectively.

Density and viscosity jump across the border of the two fluids are smoothed using the level set functions such as

$$\rho = \rho_\infty + (\rho_{\text{water}} - \rho_\infty)\phi$$

(3)

$$\mu = \mu_\infty + (\mu_{\text{water}} - \mu_\infty)\phi.$$  

(4)
The curvature of the profile is calculated by taking the second derivative of \( \phi \).

Mesh mapping has to be carefully done. Dense meshing can result in a very accurate simulation result, however, at the expense of an extremely long computation time. Quadratic elements are used in meshing. The boundary condition imposed between the resist and the substrate is that of a wetted wall, which basically implies that a fixed contact angle of \( \theta \) between the resist and the substrate is assumed, and sets the velocity component normal to the substrate to zero:

\[
\mathbf{V} \cdot \mathbf{n}_{\text{wall}} = 0,
\]

where \( \mathbf{n}_{\text{wall}} \) is normal to the surface of the substrate.

The wetted wall boundary condition also adds a frictional boundary force described by

\[
\mathbf{P}_N = -\frac{\mu}{\beta} \nabla \phi
\]

in which, \( \mu \) is the viscosity and \( \beta \) is the slip length.

The slip length is set equal to the mesh element size. For silicon a substrate as in experiments with a defined contact angle of 19° with the resist is selected. The critical surface tension coefficient of the resist is set to 19 mN m\(^{-1}\), according to the datasheet.

Figure 5 shows the simulation result obtained using COMSOL for both long and short structures. An animation type of simulation is done with time steps of 0.1 ms. The border contour obtained from simulation is also compared with the measured profile for a 120 \( \mu \)m wide resist structure and a 540 \( \mu \)m wide one. It can be seen that FEM simulation gives a good prediction of the profile.

The most important thing to find from simulation is the movement of the contact line. However, the resolution to find this value is limited to the mesh size. After measuring the contact angle of the resist and the substrate experimentally, FEM simulation can give an expected shape of the resist after reflow and its contact line movement.

4. Mask design and lithography

A mask is designed composed of an array of holes or trenches with a density based on the amount of the photoresist that has to be locally removed. Using the mask to photolithographically define the pattern into the LVOF layer and a subsequent reflow, transforms the original flat surface into a tapered shape. Figure 6 shows the model used for the calculation of the position of trenches in the resist layer as a function of strip dimensions and the required taper angle, \( \psi \).

The initial layer is of uniform thickness \( t \) and the length of the strip is \( L \). Assume an application, which calls for a taper of thickness \( t - d_1 \) at the long wavelength limit and \( t - d_2 \) at the short wavelength limit. The AB line is the top of the photoresist before any pattern and the line CD is the surface of the photoresist after patterning and reflow. The pattern in principle consists of many thin trench lines of width \( \Delta x \) which are made by lithography in the photoresist layer. \( \Delta x \) is the thinnest line we can be achieved by the available lithography system. The idea used in the design is to start from point A and move toward point B (the length of the tapered structure depends on the application) and make a trench with a width of \( \Delta x \) in the middle of an area in which the volume removed by the trench is the volume that should be removed if we want to
have the desired slope of CD at that area. This can be shown with equation (7):

\[
\frac{(x_1 - x_0)(a + b)}{2} = t \times \Delta x
\]  

(7)

From figure 6 we have

\[
\frac{a}{L' + x_0} = \frac{b}{L' + x_1} = \frac{d_2}{L'} = \frac{a + b}{[(L' + x_0) + (L' + x_1)]}
\]

→ \((a + b) = [(L' + x_0) + (L' + x_1)] \frac{d_2}{L' + L'} \)  

(8)

Replacing \((a + b)\) in (7) with (8) we obtain

\[
\frac{[(L' + x_1) - (L' + x_0)]([(L' + x_0) + (L' + x_1)])}{\Delta x} = \frac{2 \times t \times \Delta x}{[(L' + x_1)^2 - (L' + x_0)^2]}
\]

\[
= \frac{2(L' + L)t \times \Delta x}{d_2}
\]

(9)

By expressing \(x_1\) based on other variables in equation (9) an iterative procedure is derived:

\[
x_1 = \sqrt{2 \Delta x \frac{1}{a_1 - a_2}} \frac{L}{L} \frac{\alpha_1}{\alpha_1 - \alpha_2} \left( x_0 + L \frac{\alpha_2}{\alpha_1 - \alpha_2} \right)^2 - \frac{L}{L} \frac{\alpha_2}{\alpha_1 - \alpha_2}
\]

(10)

in which \(\alpha_1 = \frac{t}{L} \), \(\alpha_2 = \frac{t}{2L} \). In the iterative procedure \(x_0\) is the position at which the necessary trenches are drawn and \(x_1\) is calculated by (10) in a way that the next trench be made in the middle of \(x_0\) and \(x_1\). The initial value for \(x_0\) is zero and the procedure stops when \(x_1 - x_0 < 2 \Delta x\). The slope of the taper can be calculated in degrees with the equation below:

\[
\psi = \text{Arc tan} \left( \frac{d_2 - d_1}{L} \right) = \text{Arc tan} \left( \frac{t}{L} \frac{\alpha_1 - \alpha_2}{\alpha_1 \alpha_2} \right).
\]

(11)

Based on the above equations several patterns with different values of \(\alpha_1\) and \(\alpha_2\) are made on a full wafer 4 inch mask.

The patterns are divided into two parts. For the first part \(L = 2\) mm and there are 10 different values for \(\frac{1}{\alpha_1} = \frac{t}{L}\) and nine values for the slope. The purpose is to have 10 different slopes with nine different heights in 90 patterns. For the second part \(L = 5\) mm and there are 25 patterns with different slopes. The benefit is that several different kinds of tapered structures can be fabricated in one processing time. The layout of the mask is shown in figure 7(a). Figure 7(b) shows one patterned photoresist structure made by the mask.

5. Experimental results

To do the lithography of the patterns, the wafers are exposed to HMDS vapor for 10 min prior to resist coating to enhance adhesion between the resist and the substrate. The resist is subsequently spin coated and soft baked using the standard recipe. After exposure to UV using a contact aligner and subsequent developing, the strips with trench patterns to be used for the reflow experiments are available in the photoresist as shown in figure 7(b). As can be noted in figure 7(b), the density of the trenches increases symmetrically from the center toward both edges. After reflow the result should appear like a symmetrical hill-alike slope.

Figure 8 shows the profile of a pattern on the AZ4562 resist. Reflowed only thermally.

Figure 9 shows the profile measurement on four tapered photoresist patterns fabricated using AZ9260 and the chemical–thermal reflow technique. The surfaces of the resist patterns are very smooth and without ripple. Thickness variations are 0.2 \(\mu\)m, 0.6 \(\mu\)m, 1.5 \(\mu\)m and 2 \(\mu\)m over a 5 mm length.

Figure 10 shows a photograph of four tapered structures fabricated using the combined thermal–chemical reflow of
Figure 9. Profile of four 10 mm long tapered structures with different slopes and constant thickness in AZ9260.

Figure 10. Color change along the resist layer proves the thickness slope. The slope increases from bottom to the top.

Figure 11. Profile of two 4 mm long tapered structures in SPR3012 with different slopes and same thickness in AZ9260 also. The color profile along the structure due to interference demonstrates the sloped layer thickness.

Figure 12. Slope in silicon by transferring resist topography using plasma etching.

Figure 13. Measured reflected laser light during the etching resist on SiO2.

6. Topography transfer by plasma etching

The resist topography can be subsequently transferred into silicon or silicon-dioxide if required by the application. The benefit is that such materials are more stable and more reliable for different devices. Another benefit is better optical properties of these materials than the resist. Moreover, by increasing the resist/oxide or resist/silicon etch-rate ratio the slope angle can be further decreased.

Plasma etching in a CF4, SF6 and O2 mixture has been used to transfer the resist slopes into silicon. Figure 12 shows the results of profile measurement on the etched structure in silicon. Although a number of imperfections can be identified in the replicated structure, which can be significantly reduced in an improved etching process, the result shows a structure that can be used as an LVOF in the infrared (IR) spectral range. The thickness of the silicon structure changes 150 nm in 5 mm length, which gives a slope of 0.0017°.

The process for transferring the resist structure into SiO2 is optimized to get an optically smooth surface. The process uses a mixture of NF3 5 sccm, Ar 50 sccm and O2 30 sccm at 20 mTorr with 100 W. For optimization purposes a test resist is patterned on 500 nm of SiO2. A laser spot is directed to the surface of the resist and reflected light is measured and plotted versus time. The plot is shown in figure 13 for the optimized mixture of gases. Considering the change in optical thickness...
and interference, the etched rates of the resist and SiO₂ can be derived from the graph. The frequency of the plot changes in the interface between the resist and SiO₂. The strong signal of reflected light from SiO₂ in the plot shows that the oxide surface is not getting rough. For this process a resist/SiO₂ etch ratio of 6.2 is achieved. It is possible to increase the etch-rate ratio by increasing the O₂ flow. However, the penalty is a rougher surface on SiO₂. Increasing Ar flow results in a smoother surface by sputtering away the small roughness.

Figure 14 shows a full wafer after pattern transfer into SiO₂ together with a closer look at some of the structures. Color change along the structure illustrates the thickness variation of SiO₂.

Figure 15 shows some profile measurement on tapered SiO₂ structures. The center of the structure in figure 15(b) shows three bumps. The density of the trenches in the patterned resist was very low at the center position. Therefore, the width of the three center resist strips is large and, as discussed in section 4.4, no complete reflow has taken place at that location.

7. Conclusions and future work

Reflow for the fabrication of tapered optical structures has been demonstrated. The intended application is in the IC-compatible fabrication of LVOF’s. Combining thermal reflow with chemical reflow has demonstrated acceptable results. This approach is very suitable for the IC-compatible batch fabrication of LVOF’s of different slopes on one single chip. The resist structures have also been transferred into silicon and silicon dioxide.

For the future work the tapered structure will be fabricated on the top of a dielectric stack. The deposition of the second stack of dielectric mirrors on the tapered layer will complete an LVOF. The surface roughness of the tapered layer has to be minimized in order to get a high-resolution LVOF. More simulation work has to be done to simulate joining of the many strips during the reflow process.

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