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Thin-film amorphous silicon germanium solar cells with p- and n-type hydrogenated silicon oxide layers

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A R T I C L E   I N F O

Keywords: Hydrogenated amorphous silicon germanium Hydrogenated silicon oxide Light management Parasitic losses

A B S T R A C T

Mixed-phase hydrogenated silicon oxide (SiOₓ:H) is applied to thin-film hydrogenated amorphous silicon germanium (a-SiGe:H) solar cells serving as both p-doped and n-doped layers. The bandgap of p-SiOₓ:H is adjusted to achieve a highly-transparent window layer while also providing a strong electric field. Bandgap grading of n-SiOₓ:H is designed to obtain a smooth transition of the energy band edge from the intrinsic to n-doped layer, without the need of an amorphous buffer layer. With the optimized optical and electrical structure, a high conversion efficiency of 9.41% has been achieved. Having eliminated other doped materials without sacrificing performance, the sole use of SiOₓ:H in the doped layers of a-SiGe:H cells opens up great flexibility in the design of high-efficiency multi-junction thin-film silicon-based solar cells.

1. Introduction

Hydrogenated amorphous silicon germanium (a-SiGe:H) is a meaningful building block in multi-junction thin-film silicon-based solar cells. Its electronic bandgap decreases as the Ge content increases in the matrix [1–3]. As an absorber material, a-SiGe:H offers an adjustable, intermediate bandgap between those given by the widely used hydrogenated amorphous silicon (a-Si:H, 1.7 eV) and hydrogenated nanocrystalline silicon (nc-Si:H, 1.1 eV). A combination of these absorbers can lead to better spectral utilization in triple- and quadruple-junction solar cells [4–7].

To form a p-i-n junction, p-type hydrogenated amorphous silicon carbide (a-SiC:H), a-Si:H, nc-Si:H, hydrogenated silicon oxide (SiOₓ:H), and n-type a-Si:H, nc-Si:H, SiOₓ:H have been used in the p- and n-layer, respectively, of a-Si:H solar cells. Among these materials, SiOₓ:H is worth extra attention because of its multifunctionality and adaptability demonstrated in the application of thin-film silicon-based solar cells. SiOₓ:H is a mixed-phase material comprising nanocrystalline silicon filaments embedded in an amorphous silicon oxide matrix [8–14]. The anisotropic growth of the material results in good transverse and poor lateral conductivity, thus improving the electrical performance of the solar cell. Optically, its relatively wide bandgap and low absorption coefficient suppress own parasitic absorption, while providing a high built-in voltage as a p-layer. The tunable refractive index offers a means to reduce the reflection at the window layer, or to enhance the intermediate reflection when used after an absorber layer [5,15–17]. Given that well-performing tunnel recombination junctions (TRJs) can be formed between p-type and n-type SiOₓ:H [17], SiOₓ:H opens a great versatility for the light management in multi-junction solar cells. Lastly, n-type SiOₓ:H/Ag was also used to replace the typical ZnO:Al/Ag back reflector, as the low refractive index of SiOₓ:H can limit the plasmonic absorption in Ag [18–20], just like in ZnO:Al/Ag stack. Fig. 1 shows the optical properties of the p- and n-type SiOₓ:H, in comparison to other typical materials used in a-SiGe:H cells.

So far, the use of SiOₓ:H in a-SiGe:H solar cells is not as popular as it is in a-Si:H or nc-Si:H cells. N-type SiOₓ:H has been used as the n-layer of a-SiGe:H cells behind either intrinsic a-Si:H [6] or a-Si:H/nc-Si:H [21] buffer layer(s). On the p-side, Schüttauf et al. tested several p-layer configurations and reported a significant drop in Voc when the p-layer was solely constituted of p-type SiOₓ:H, compared to those which also include a a-SiC:H or a-Si:H layer between the oxide layer and the absorber [6]. Nevertheless, taking into account the adaptability of SiOₓ:H materials and its successful applications in a-Si:H and nc-Si:H solar cells, it is natural to think of making a a-SiGe:H cell using only SiOₓ:H for the doped layers, without compromising the performance. From the optical perspective, it can reduce the parasitic absorption caused by the other more absorptive supporting materials, and at the same time provide the most flexibility in light management for both single- and multi-junction devices. Electrically, the doped SiOₓ:H materials can be tuned such that they give favorable band offset at the interfaces, or form efficient TRJs with minimized losses. In this paper, we explore the possibility of using both p- and n-type SiOₓ:H in thin-film a-SiGe:H solar cells without extra buffer layers adjacent to the

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absorber layer. Additionally, we evaluate the performance of n-SiO2:H/Ag back reflectors and optimize the light management of such device.

2. Experimental

The thin-film silicon alloy materials, including intrinsic a-SiGe:H, p- and n-type SiO2:H and n-type a-Si:H, were deposited in a cluster tool using plasma-enhanced chemical vapor deposition (PECVD) at radio frequency (RF) of 13.56 MHz. The maximum substrate area and tool using plasma-enhanced chemical vapor deposition (PECVD) at multiple incident angles using an M – 2000D® Spectroscopic Ellipsometer (J.A. Woollam Co.), which is equipped with a dual-lamp light source covering the wavelength range of 193–1690 nm. To determine the optical bandgap of a-SiGe:H, the Tauc-Lorentz model [28,29] was fitted to the SE data measured from samples of single thin films deposited on glass substrates, and the result gave the optical bandgap $E_{\text{opt}}$. In addition, optical bandgap $E_{\text{opt}}$ was found as the energy with which a photon has an absorption coefficient of $1 \times 10^{4}$ cm$^{-1}$ in the material.

The performance of the solar cells was examined by the illuminated current-voltage ($I-V$) measurement and external quantum efficiency (EQE) measurement. The $I-V$ measurement was conducted at a controlled cell temperature of 25°C, with a dual-lamp continuous solar simulator (WACOM WXS-908-L2, class AAA). The two filtered lamps in the solar simulator were adjusted with two monocrystalline silicon reference cells manufactured by and traceable to the Fraunhofer Institute for Solar Energy Systems (ISE), to provide an incident irradiance of 1000 W m$^{-2}$ with optimum spectral matching with AM1.5G solar spectrum. The EQE measurement was performed using an in-house system, in which the electrical signal is detected by a lock-in amplifier, and the chopped monochromatic light is provided by a xenon light source, a 3-grating monochromator and a chopper. A silicon photodiode, which was regularly calibrated by Fraunhofer ISE, was used to calibrate the light source in the EQE measurement. As for the external parameters of the solar cells, the open-circuit voltage ($V_{\text{OC}}$) and fill-factor (FF) are determined by the $I-V$ measurement, while the short-circuit current density ($J_{\text{SC}}$) is calculated by weighting the measured EQE with the AM1.5G solar spectrum. The cell performance reported in this work is the average taken from the 12 better performing cells (40%) out of the 30 cells on each sample. While the EQE measurement was conducted on the best performing cell in each sample, the reported deviation in $J_{\text{SC}}$ was obtained from the $I-V$ measurement.

3. Results and discussion

3.1. Use of SiO2:H in p-layer

In this section, p-SiO2:H is applied in a-SiGe:H cells as the main p-layer, and the cell performance is examined. To compare the effect of different p-SiO2:H, the device structure except the p-layer was kept the

Table 1  Deposition conditions of the PECVD materials used in this work. The constitution of gaseous precursors is shown by their flow rates in standard cubic centimeter per minute (sccm).

<table>
<thead>
<tr>
<th>Material</th>
<th>T (°C)</th>
<th>$p$ (mbar)</th>
<th>$P$ (mW cm$^{-2}$)</th>
<th>SiH$_4$</th>
<th>GeH$_4$</th>
<th>H$_2$</th>
<th>B$_2$H$_6$</th>
<th>PH$_3$</th>
<th>CO$_2$</th>
</tr>
</thead>
<tbody>
<tr>
<td>p-nc-Si:H</td>
<td>180</td>
<td>2.2</td>
<td>243.1</td>
<td>0.8</td>
<td>–</td>
<td>190</td>
<td>0.004</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>p-SiO$_2$:H</td>
<td>2.2</td>
<td>90.3</td>
<td>30</td>
<td>0.8</td>
<td>–</td>
<td>190</td>
<td>0.004</td>
<td>–</td>
<td>0.8–2.2</td>
</tr>
<tr>
<td>n-nc-Si:H</td>
<td>2.6</td>
<td>20.8</td>
<td>30.0</td>
<td>3.65</td>
<td>200</td>
<td>–</td>
<td>–</td>
<td>10.8</td>
<td>–</td>
</tr>
<tr>
<td>n-a-Si:H</td>
<td>0.6</td>
<td>27.8</td>
<td>40.0</td>
<td>–</td>
<td>101</td>
<td>–</td>
<td>0.22</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>n-SiO$_2$:H</td>
<td>1.5</td>
<td>69.4</td>
<td>1.0</td>
<td>–</td>
<td>101</td>
<td>–</td>
<td>0.024</td>
<td>0.5–2.6</td>
<td>–</td>
</tr>
</tbody>
</table>
same as described in Section 2. No buffer layers were used between the p-i-n interface. A series of solar cells were fabricated with different CO2 flow rate during the deposition of p-layer. The different oxygen content in p-SiOx:H results in different optical bandgap and electrical conductivity of the material [17].

The optical effect of the different p-SiOx:H layers is first investigated. The p-layers were deposited with a flow rate ratio \( R(\text{CO2}/\text{SiH4}) \) from 1.00 to 2.75. External quantum efficiency was measured at a bias voltage of −0.5 V to reduce the influence of inefficient carrier collection and to estimate the optical absorption in the absorber layer. The measured result is shown in Fig. 2. It is clear that in the short wavelengths between 300 nm and 600 nm, the greater the CO2 flow during the p-layer deposition, the greater the spectral response of the solar cell exhibits. The difference in blue response comes mainly from the parasitic absorption in the p-layer. With more CO2, the deposited p-SiOx:H has a larger optical bandgap and lower absorption coefficient, so it allows more high-energy photons to be utilized in the intrinsic layer and contribute to the photocurrent. Beyond 600 nm, the EQE is very similar between the cells (not shown) because the p-layers barely absorb any photons in that region.

Electrically, the cells in Fig. 2 show similar performance. With a lowest \( E_{\text{Tauc}}(E_\text{gap}) \) of 1.41 eV (1.62 eV) in the intrinsic layer, all the cells have a FF in the range of 54–56%, as shown in Fig. 3. The \( V_{OC} \) shows a slight still monotonic increase with increased CO2 flow from 654 mV to 660 mV. For a a-SiGe:H absorber with a wider bandgap, however, the p-SiOx:H has a greater impact on the electrical performance of the solar cell. When the absorber has an \( E_{\text{Tauc}}(E_\text{gap}) \) of 1.45 eV (1.67 eV) at its lowest point, both \( V_{OC} \) and FF show a monotonic increase with ratio \( R(\text{CO2}/\text{SiH4}) \).

### 3.2. Use of SiOx:H in n-layer

In a-Si:H solar cells, both n-a-Si:H and n-SiOx:H acting as the n-layer can lead to decent electrical performance of the cells. The use of n-SiOx:H is more favorable when the optical losses and the behavior in forming TRJs [5,30] are considered. Based on such experience, several configurations of n-layer have been tested in this work to see how well they work for a-SiGe:H cells. The tested structures are summarized in Table 2. Configurations a and c comprise only a single layer of n-a-Si:H and n-SiOx:H, respectively. The ratio \( R(\text{CO2}/\text{SiH4}) \) for n-SiOx:H was chosen to be 2.0 after some preliminary optimization. Configuration b, c, and d have the same n-SiOx:H layer as e, but with an n-a-Si:H layer between the n-SiOx:H and the intrinsic layers. Lastly, the deposition of configuration f started with a bandgap grading of n-SiOx:H, in which the ratio \( R(\text{CO2}/\text{SiH4}) \) was gradually changed from 1.0 to 2.0. The second half of the n-SiOx:H deposition was then continued at \( R(\text{CO2}/\text{SiH4}) \) of 2.0. The influence of different n-layer configurations was examined by the performance of the solar cells. A series of a-SiGe:H cells were fabricated using identical device structure except for the n-layer. p-SiOx:H with \( R(\text{CO2}/\text{SiH4}) \) of 2.0 was used as the p-layer and a-SiGe:H with a lowest \( E_{\text{Tauc}} \) of 1.45 eV was used as the i-layer.

The external parameters of the solar cells measured under AM1.5G spectrum are shown in Fig. 4. It is the most noticeable that configuration a provides significantly less photocurrent than all the others do. In this structure, the cell performance suffers from severe plasmonic absorption at the Ag back reflector because the adjacent material, which is n-a-Si:H, has a high refractive index [19]. From configuration a, with an additional n-SiOx:H layer acting as an optical spacer, configuration b exhibits improved \( J_{SC} \) and slightly dropped FF. When comparing configurations b, c, and d, the reduction of n-a-Si:H thickness mitigates the parasitic absorption, leading to the increase in \( J_{SC} \). Furthermore, shrinking the defective n-a-Si:H layer also slightly improves \( V_{OC} \) and FF. n-SiOx:H can act as n-type material to provide the necessary electric field, so the role of n-a-Si:H in these configurations is more about interface passivation and diminishing band offset. Without such buffer in between, configuration e gives worse \( V_{OC} \) and FF probably because of the band discontinuity between the absorber and the wide-gap n-SiOx:H. To tackle this problem without using n-a-Si:H, configuration f aims at achieving a smooth transition from the intrinsic layer to narrow-gap n-SiOx:H to wide-gap n-SiOx:H by means of CO2 grading during the deposition. As a result, all external parameters are improved from configuration e to f. Remarkably, this n-layer is solely constructed of n-SiOx:H while its best sample tops the conversion efficiency among all compared configurations.

### 3.3. Optical spacer at back reflectors

It is demonstrated in Section 3.2 that because of its low refractive index, n-SiOx:H can function as optical spacer at the back reflector to suppress the plasmonic absorption in Ag. Conventionally, ZnO:Al is

### Table 2

<table>
<thead>
<tr>
<th>Configuration</th>
<th>n-layer tested for a-SiGe:H cells.</th>
</tr>
</thead>
<tbody>
<tr>
<td>a</td>
<td>20 nm n-a-Si:H</td>
</tr>
<tr>
<td>b</td>
<td>20 nm n-a-Si:H/60 nm n-SiOx:H</td>
</tr>
<tr>
<td>c</td>
<td>10 nm n-a-Si:H/60 nm n-SiOx:H</td>
</tr>
<tr>
<td>d</td>
<td>5 nm n-a-Si:H/60 nm n-SiOx:H</td>
</tr>
<tr>
<td>e</td>
<td>60 nm n-SiOx:H</td>
</tr>
<tr>
<td>f</td>
<td>30 nm graded n-SiOx:H/30 nm n-SiOx:H</td>
</tr>
</tbody>
</table>
used together with Ag back reflector for the same purpose [19]. On the other hand, the free carrier absorption in ZnO:Al [31] can potentially deteriorate the spectral response of the solar cell at long wavelengths.

In this section, we investigate the feasibility of fully replacing ZnO:Al with n-SiO$_2$:H in a-SiGe:H solar cells. A series of cells were fabricated with different configurations of optical spacer before the Ag back reflector. The basic structure includes p- and i-layers identical to the ones used in Section 3.2, as well as a 30-nm-thick n-SiO$_2$:H layer with bandgap grading. On top of it, wide-gap n-SiO$_2$:H and/or ZnO:Al with various thicknesses was deposited before completing the back reflector with Ag. The thickness of wide-gap n-SiO$_2$:H and ZnO:Al are 0/30 nm and 0/40/80/120 nm, respectively.

Fig. 5 shows the measured EQE of the a-SiGe:H cells with different spacer configurations. Without an additional layer of wide-gap n-SiO$_2$:H, it can be clearly seen in Fig. 5(a) that the spectral response at wavelengths longer than 700 nm is enhanced when the thickness of ZnO:Al is increased from 0 nm to 40 nm and further to 80 nm, showing the effect of ZnO:Al on reducing plasmonic absorption. When the thickness of ZnO:Al is further increased to 120 nm, the spectral response is not as good as the 80-nm-thick counterpart. We speculate that 80-nm-thick ZnO:Al is capable of quenching most of the plasmonic absorption so the extra thickness of 40 nm only increases the parasitic absorption in ZnO:Al.

A direct comparison in Fig. 5(a) shows that the cell with 60/0 nm of n-SiO$_2$:H/ZnO:Al (i.e. no ZnO:Al) exhibits comparable spectral response to that of the cell with thinner n-SiO$_2$:H and 40-nm-thick ZnO:Al spacer. The addition of 30-nm-thick wide-gap n-SiO$_2$:H layer provides a similar effect as 40 nm of ZnO:Al. It confirms the capability of n-SiO$_2$:H material being an optical spacer to avoid parasitic losses. Fig. 5(b) shows the EQE spectra of cells with different thickness of ZnO:Al on top of the extended n-SiO$_2$:H layer. A similar improvement in the long wavelengths is observed when the thickness of ZnO:Al is increased, even with the thicker n-SiO$_2$:H layer. Two remarks should be noted for clarification: Firstly, all solar cells reported in Fig. 5 show similar $V_{oc}$ and FF, so better spectral response can be translated to better cell performance. Secondly, the strong EQE variation between the samples in the wavelengths of 600–700 nm resembles the reflectance spectrum of the respective cells (not shown). It is probably due to the interference effect, which is sensitive to the thickness of the films.

By extrapolation, one can expect comparable optical performance from replacing the thick ZnO:Al layer with a thick wide-gap n-SiO$_2$:H layer. Nonetheless, a thick wide-gap n-SiO$_2$:H layer can be too resistive to support efficient carrier transportation in the device. Therefore, a combination of n-SiO$_2$:H and ZnO:Al both with medium thickness should give the best overall performance.

### 3.4. Improvement by light scattering substrates

Light scattering is particularly important for a-SiGe:H solar cells. The intrinsic a-SiGe:H materials are rather defective compared to a-Si:H, so the absorber layer in a-SiGe:H cells has to be very thin to ensure adequate carrier collection efficiency of the device. Light scattering for enhancing the light absorption in the long wavelengths is therefore necessary for achieving high photocurrent within the thin absorber. In this section, the importance of light scattering for a-SiGe:H solar cells is highlighted by comparing cells fabricated on different light-scattering substrates.

The native roughness on APCVD SnO$_2$:F is too small to provide effective light scattering near the band edge of a-SiGe:H. On this respect, the surface of ZnO:B deposited by low-pressure CVD (LPCVD) typically features a greater roughness and larger (average) local surface inclination angle thus better light scattering capability among the longer wavelengths [32,33]. For comparison, solar cells were deposited on these two types of substrate, in specific, SnO$_2$:F (Asahi VU-type) and LPCVD ZnO:B with a thickness of 2.3 μm and smoothened by Ar plasma treatment for 4 min. The optimized device structure being used contains p-SiO$_2$:H with $R(CO_2/SiH_4)=2.0$, 60 nm of n-SiO$_2$:H with bandgap grading in its first half and $R(CO_2/SiH_4)=2.0$ in the second half, and 40 nm of ZnO:Al. In addition, the hydrogen dilution during a-SiGe:H deposition was slightly increased to improve the material
quality, while the ratio between SiH₄ and GeH₄ was maintained the same.

The performance of the optimized solar cells is shown in Fig. 6. On Asahi-VU substrates, the optimized device structure using doped SiOₓ:H for both p- and n-layers delivers a conversion efficiency of 8.8%. For the cells on ZnO:B substrates, the Voc and FF are fairly similar to that on Asahi-VU substrates. An obvious difference appears in the spectral response shown in Fig. 6(b). Except for wavelengths below 390 nm, at which ZnO:B is very absorptive, the cell made on ZnO:B shows a higher EQE over the whole spectrum of interest. It clearly demonstrates the difference in light scattering from different surface textures. The absorption enhancement leads to an improvement in JSC from 18.4 mA/cm² on Asahi-VU to 19.5 mA/cm² on ZnO:B.

Finally, the optical loss from the reflection at the air/glass interface is tackled. An anti-reflection (AR) foil (produced by temicon GmbH) was applied onto the incident surface of the cell made on ZnO:B substrate. The influence of the AR treatment on the cell reflectance and EQE is shown in Fig. 6(b). At the wavelength of 550 nm, the reflectance is reduced from 5.8% to 3.9%, which corresponds to a reduction of 33%. The EQE in the visible spectrum is enhanced, consequently. In spite of the parasitic absorption induced by the AR foil in the short wavelengths, an increased JSC of 19.8 mA/cm² is reached. In the end, as a result of the improved light management and optimized device structure, an efficiency of 9.41% is achieved.

4. Conclusions

Doped SiOₓ:H was implemented in thin-film a-SiGe:H solar cells. In the window layer, the bandgap of p-SiOₓ:H was adjusted to obtain a high Voc and decent FF, while the transparency of the wide-gap p-layer also resulted in high spectral response in the short-wavelength region. The use of n-SiOₓ:H instead of n-a-Si:H reduces the parasitic absorption in the n-layer and suppresses the plasmascopic absorption in the neighboring metal layer. A combination of wide-gap n-SiOₓ:H and TCO layers with appropriate thicknesses in front of a Ag back reflector provides the most favorable optical properties to enhancing the long-wavelength response of the solar cells. Bandgap grading of n-SiOₓ:H was introduced and its application near the i-n interface mitigates the band offset between the wide-gap n-SiOₓ:H and the absorber material. As a consequence, thin-film a-SiGe:H cells using only SiOₓ:H doped layers have been realized without giving up the opto-electrical performance. Having optimized the cell structure, a conversion efficiency of 9.41% in a thin-film a-SiGe:H solar cell has been achieved on a ZnO:B-coated substrate with anti-reflection treatment. This study paves the way for designing high-efficiency thin-film silicon-based multi-junction solar cells with the use of a-SiGe:H absorber and multifunctional SiOₓ:H layers.

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