Ship motion compensation platform for high payloads dynamic analysis and control MSc Project at GustoMSC – Wouter de Zeeuw Prof.dr. D.J. Rixen, Ir. M. Wondergem



Introduction: Pooltable on cruise ship







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Two ship motion compensation platforms:











Offshore windturbine installation with Jack-up units **GustoMSC**

Present method





Goal: Complete windturbine installation from a floating unit Motion stabilizing platform to extend operating limits

1000[t] Fast feeder barges ////

Jack up unit

stays at site





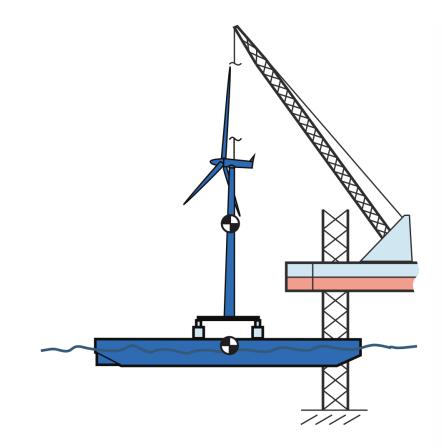
Small overview

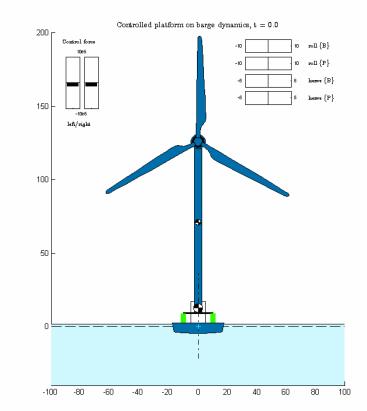
Preliminary 2D model

- 1. Analysis of Ampelmann scalemodel tests
- 2. 3D modeling of new mechanism on ship
- 3. Controlling the system



Goal: Complete windturbine installation from a floating unit Preliminary 2D model showed feasibility ...

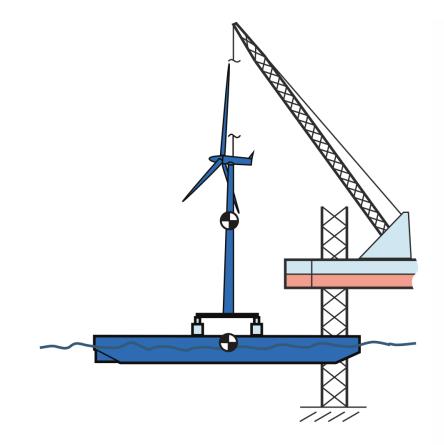


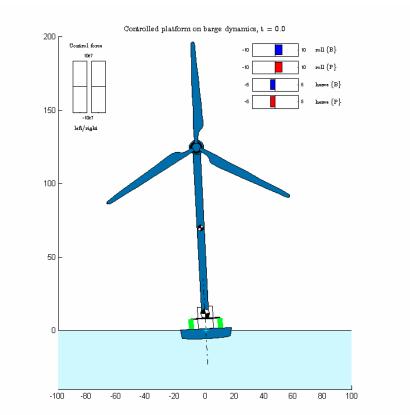




Goal: Complete windturbine installation from a floating unit

Preliminary 2D model showed feasibility but dynamic instability







1. (In-)stability due to the quasistatic control?

 k_P

 \bullet

k_S≨∏

•

 $\frac{k_P}{c_D}$

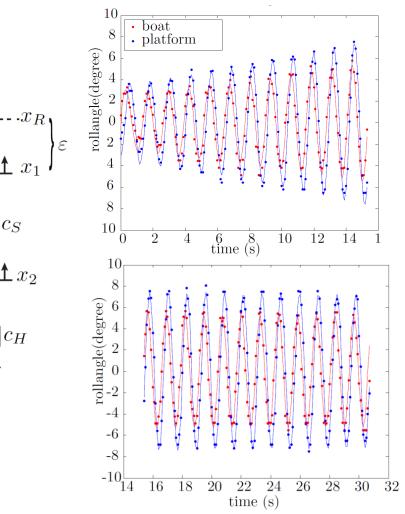
 ε

 c_S

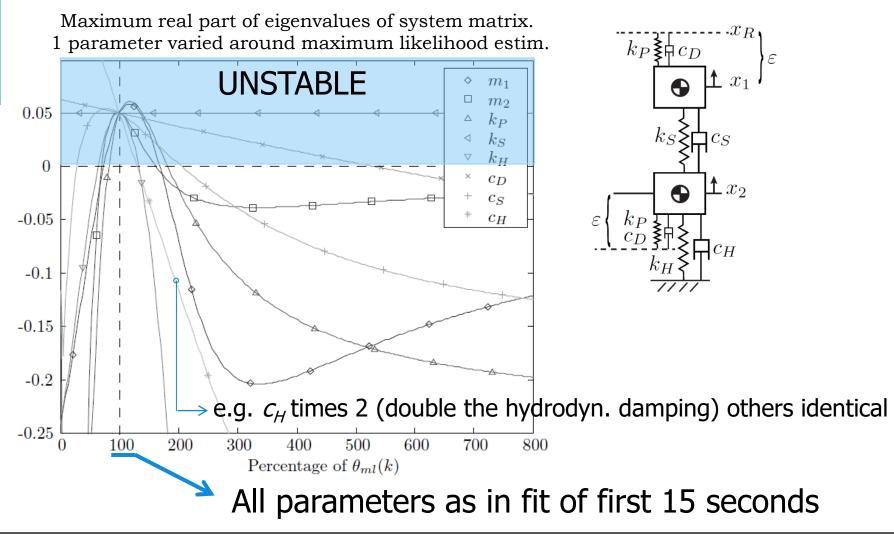
Similar mass system: Ampelmann scale model tests



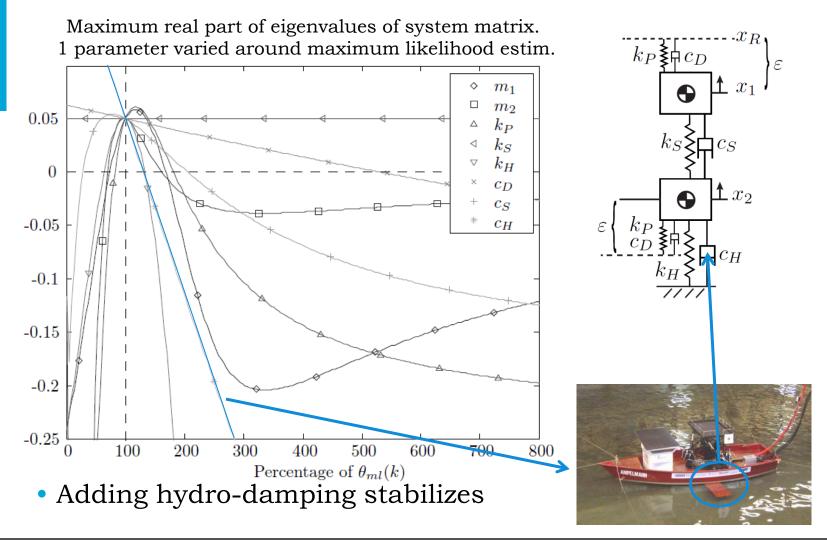




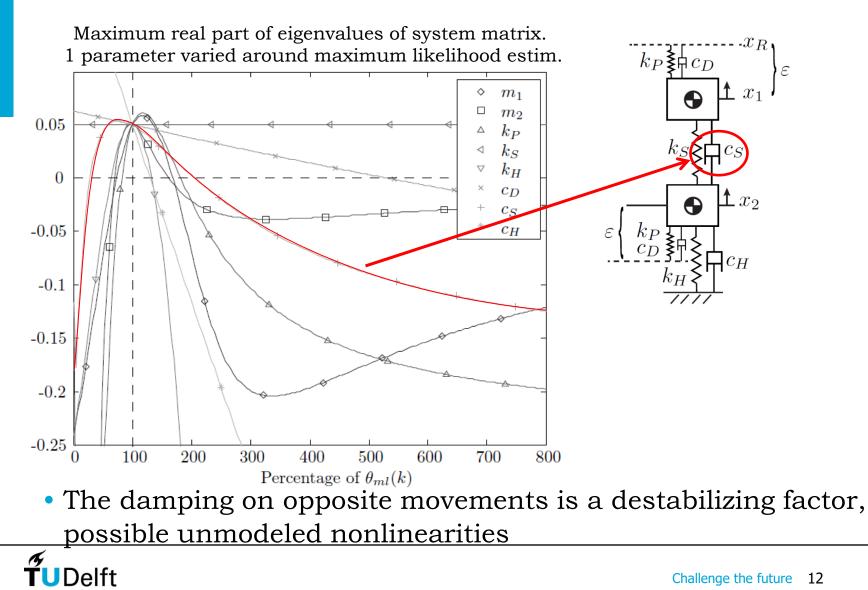




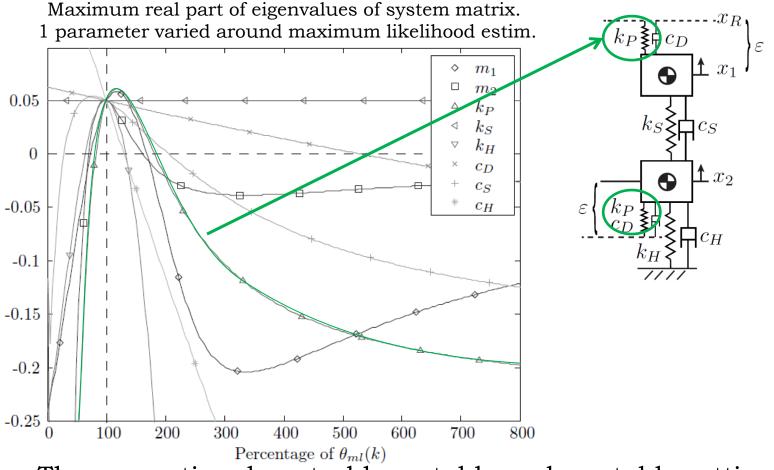




TUDelft



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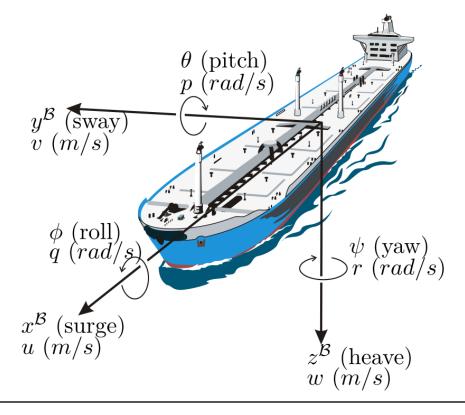


• The proportional control has stable and unstable settings

TUDelft

2. 3D modeling - ship movements

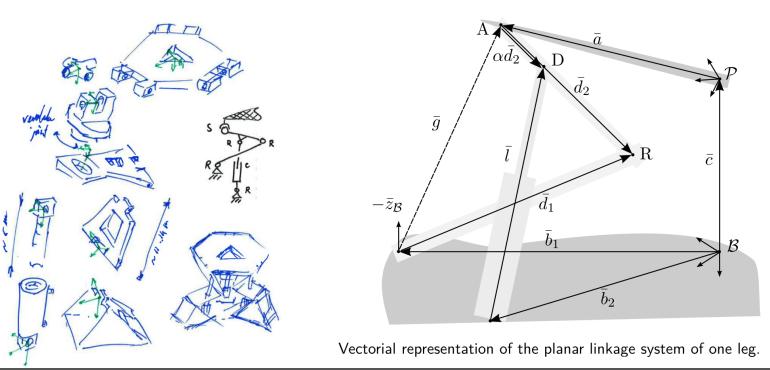
- Accelerations due to planar movements surge, sway and yaw are smaller than due to off planar movements
- Platform should compensate heave, roll and pitch



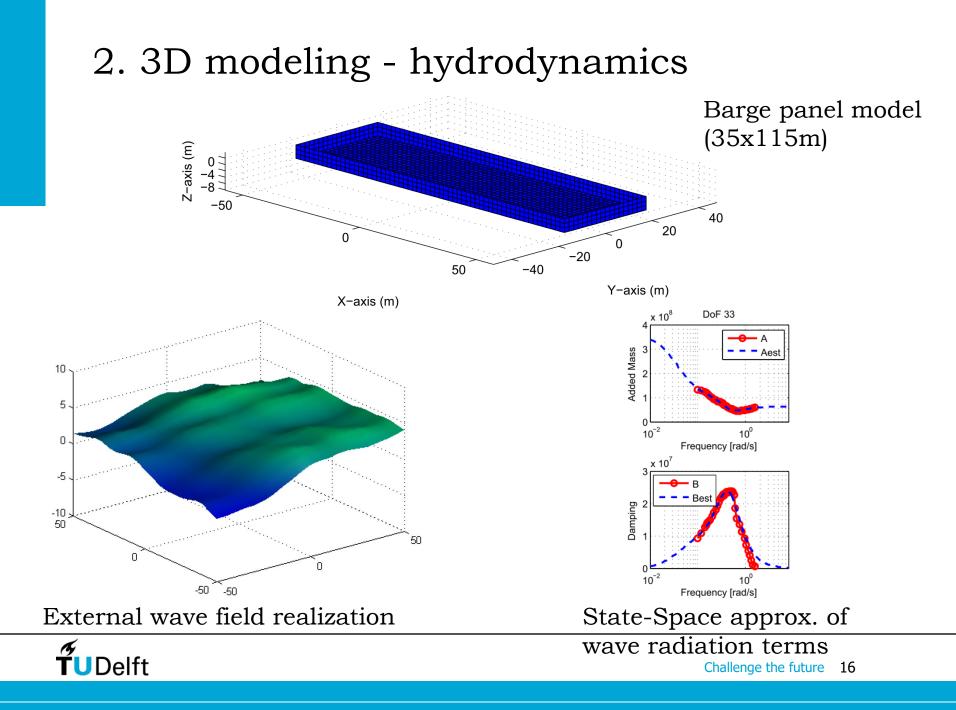


2. 3D modeling - platform mechanism

- New mechanism for a 3 degree of freedom platform
- Planar movements are constrained by 3 Sarrus type linkages
- Force vs. Reach variable via α

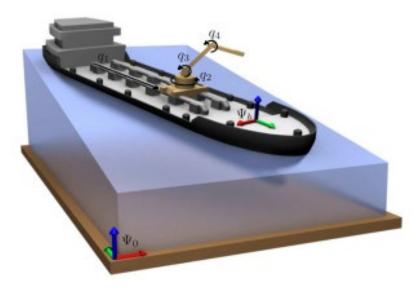






2. 3D modeling - vessel+platform

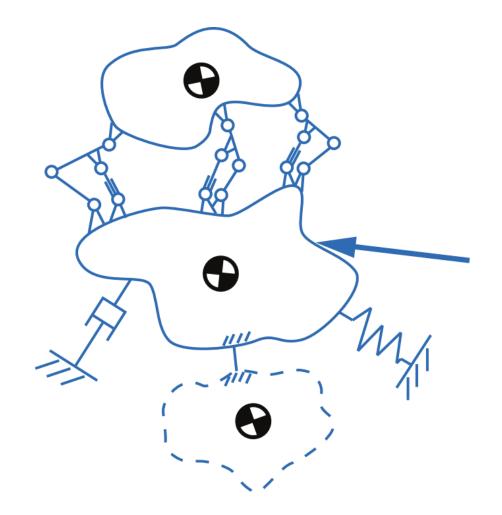
- Lagrangian dynamics (body fixed)
- Extension of serial robot on ship to parallel robots





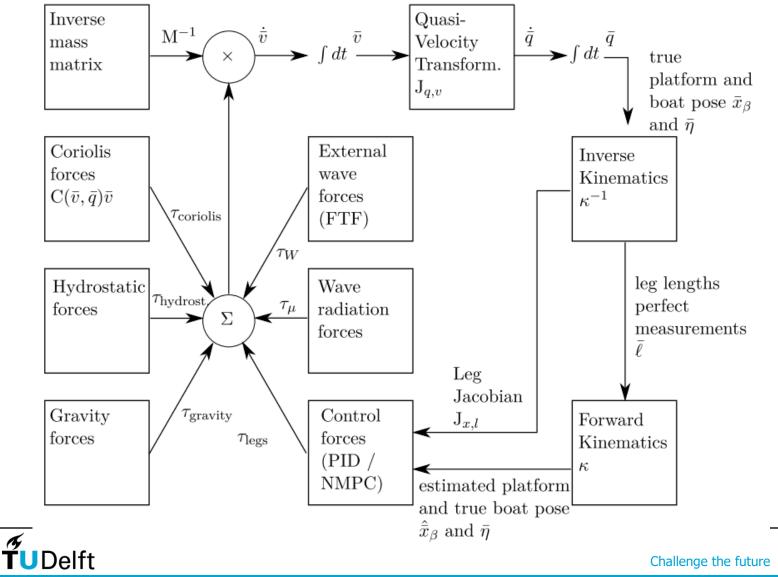
2. 3D modeling - total dynamics

- Nonlinear kinematics
- Coriolis terms
- Pose dep. mass matrix
- External waveloads
- Hydrostatics
- Hydrodynamics
 - Wave radiation
 - Added mass/damping





2. 3D modeling - total dynamics



3. Controllers - Naïve Quasistatic vs. Model Based

Quasistatic:

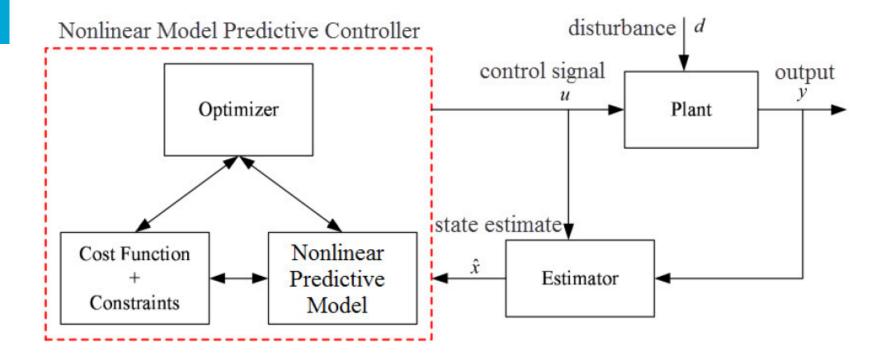
- Calculate leg length error assuming fixed boat position
- 2 Proportional-Integral-Derivative (PID) controllers
 - On mean error
 - On asymmetric errors

Model Based:

• Nonlinear Model Predictive Control

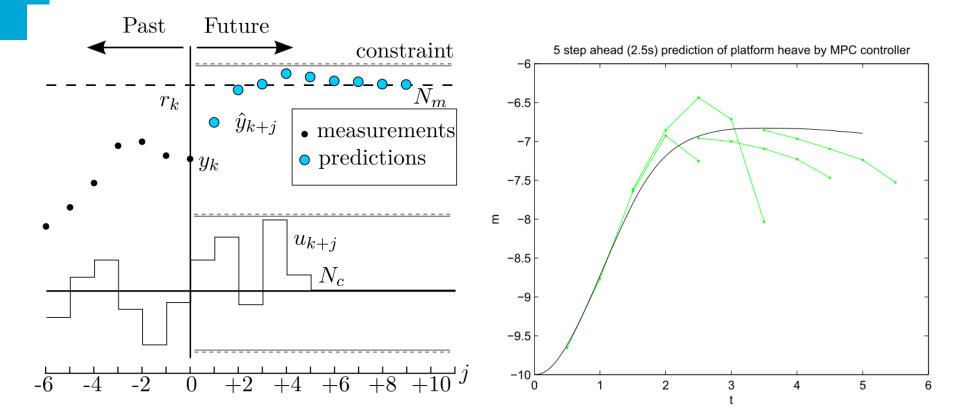


3. Control - Nonlinear Model Predictive Control



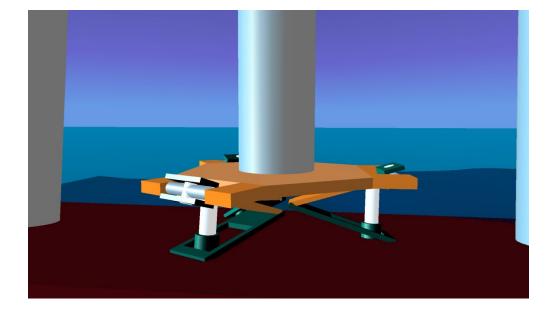


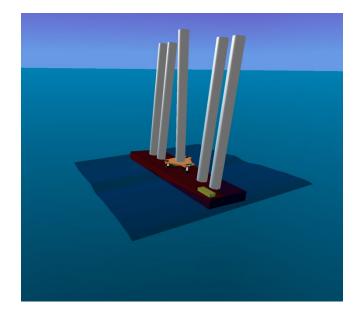
3. Control - Nonlinear Model Predictive Control





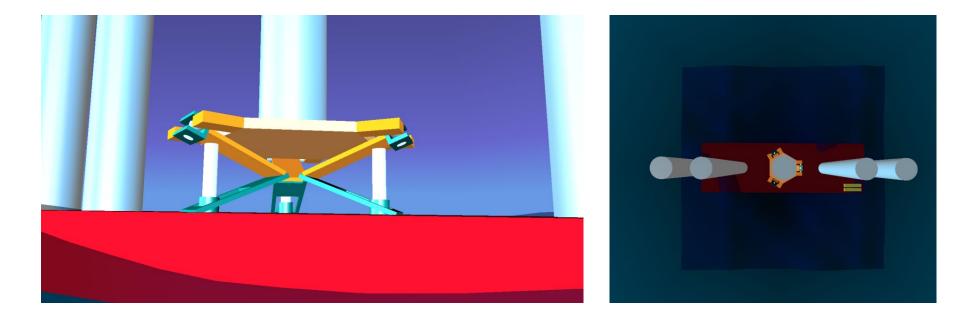
Visualizations





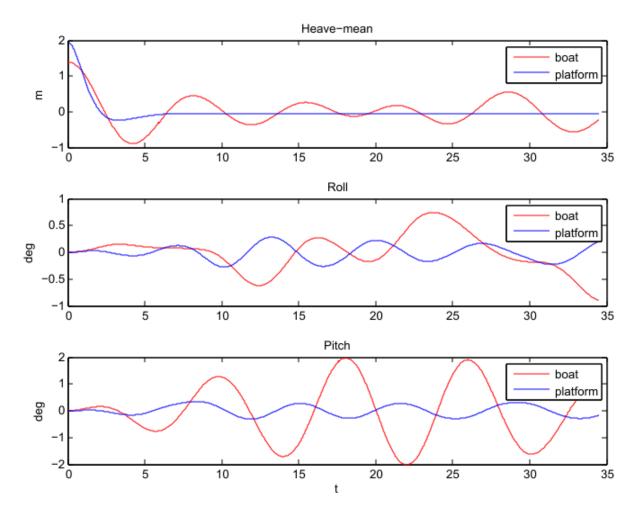


Visualizations





Heave-Roll-Pitch in storm conditions Head sea, seastate Hs=4m T1=6.5s.





Energy usage in disturbance rejection Milder sea

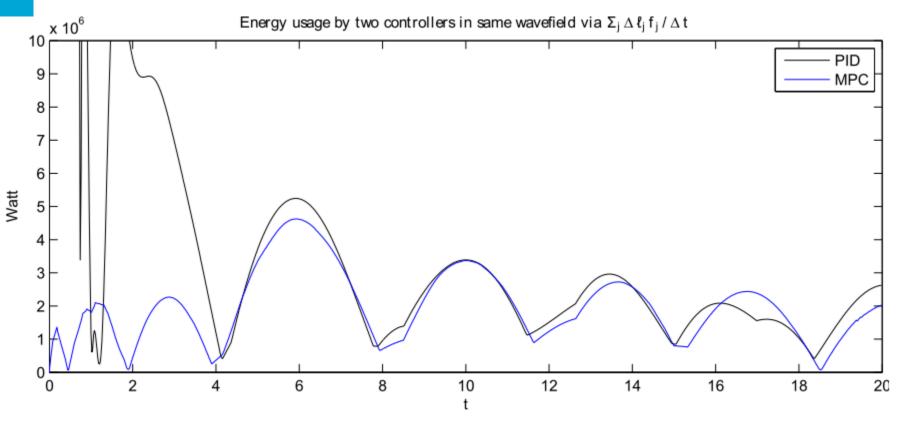


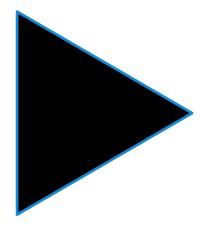
Figure 4-18: Comparison of the power required by the platform.



Conclusions

- The scalemodel roll instability can be reproduced by a linear model with quasistatic control and influential parameters can be recognized.
- The coupled ship parallel platform dynamics are derived and he new platform can compensate the ship movements.
- MPC is shown to be a successful candidate for control, requires less power than PID in disturbance rejection and is less hard to tune and to stabilize.





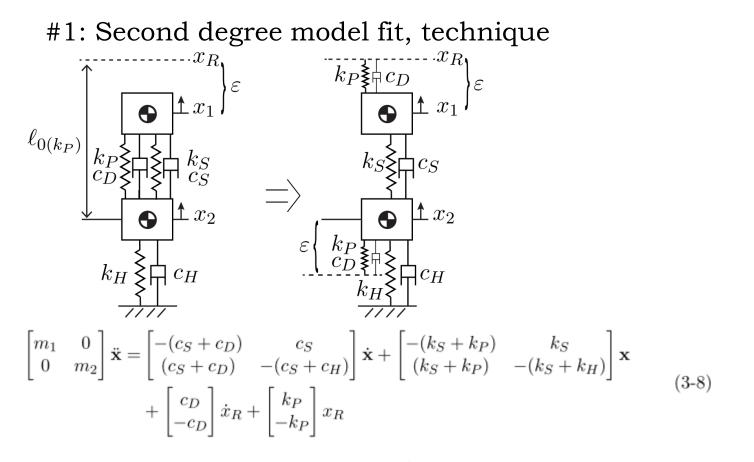


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Thank you. Questions?



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Or for with a non changing reference of zero $(x_R = 0)^1$

$$\ddot{\mathbf{x}} = \begin{bmatrix} -\frac{(c_D + c_S)}{m_1} & \frac{c_S}{m_1} \\ \frac{(c_D + c_S)}{m_2} & -\frac{(c_H + c_S)}{m_2} \end{bmatrix} \dot{\mathbf{x}} + \begin{bmatrix} -\frac{(k_P + k_S)}{m_1} & \frac{k_S}{m_1} \\ \frac{(k_P + k_S)}{m_2} & -\frac{(k_H + k_S)}{m_2} \end{bmatrix} \mathbf{x} \qquad SSR = \sum_{1}^{N} (\mathbf{y} - \hat{\mathbf{y}})^T (\mathbf{y} - \hat{\mathbf{y}})$$
Fitting technique:
$$A_s = \begin{bmatrix} C & K \\ I & \underline{0} \end{bmatrix} \quad \hat{\mathbf{y}}(t) = e^{A_s t} \mathbf{x}_0 - \mathbf{x}_0 - \mathbf{x}_0, C, K$$
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#1b: Second degree model fit, results

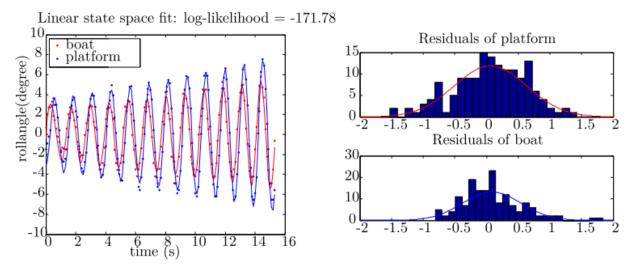


Figure 3-8: Fit for the first 15 seconds using the parametric physical model. $m_1 = 20kg$ and the first half of the dataset is used so t = 0..T/2. The dots are the measured data points and the full line is the estimated model.

$\dot{\phi}_{p0}$	23.25	(4.77)	$\dot{\phi}_{b0}$	-1.64	(-0.76)
ϕ_{p0}	-0.37	(-0.71)	ϕ_{b0}	-3.09	(-10.01)
m_1	20.00	(fixed)	m_2	72.97	(275.16)
k_P	755.63	(54.57)	k_S	0.00	(0.01)
k_H	2068.72	(164.39)	c_D	3.37	(1.40)
c_S	63.00	(43.21)	c_H	35.39	(18.46)
σ	0.74	(16.85)			

Table 3-3: Physical parameterized fit of first 15 seconds with $\log \mathcal{L} = -171.78$.



#2: Kinematics – leg joint velocity Jacobian construction

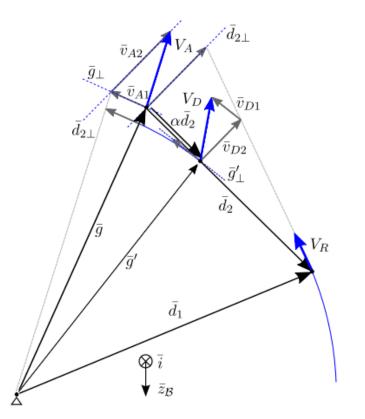


Figure 4-6: Velocity of joints in planar linkage system. The linkage is connected to the platform in A and to the hydraulics in D. The vectors \overline{d}_1 and \overline{d}_2 form the parallel mechanism of the linkage, and the auxiliary vectors \overline{g} and \overline{g}' are used in calculation and represent no parts. The main strategy of the velocity transformation is to decompose velocity vector V_A in two independent directions that represent the two rotations around the revolute joints at the floor and in R.



#3: Hydrodynamics – rad forces

Cummins eqn. in hydrodyn. ref. frame:

$$(\mathbf{M}_{\mathrm{RB}}^{\mathcal{H}} + \mathbf{A}_{\infty})\ddot{\bar{\xi}}^{\mathcal{H}} + \int_{-\infty}^{t} \mathbf{K}(t - t')\dot{\bar{\xi}}^{\mathcal{H}}(t)dt' + \mathbf{G}\bar{\xi}^{\mathcal{H}} = \bar{\tau}^{\mathcal{H}}$$

With the retardation function $\mathbf{K}(t)=\frac{2}{\pi}\int_0^\infty \mathbf{B}(\omega)\cos(\omega t)d\omega.$

Retardation forces (vector):

$$\bar{\mu}^{\mathcal{B}} = \int_{0}^{t} \mathbf{K}(t - t')\bar{\nu}(t)dt'$$
$$= \int_{0}^{t} \left[\frac{2}{\pi} \int_{0}^{\infty} \left(\mathbf{J}^{*T}\mathbf{B}(\omega)\mathbf{J}^{*}\right)\cos\left(\omega(t - t')\right)d\omega\right]\bar{\nu}(t)dt'$$

State space approx. per radiation component (scalar):

$$\dot{\bar{\chi}} = A_{rad}\bar{\chi} + B_{rad}\bar{\nu}$$

 $\mu^{\mathcal{B}} \approx \hat{\mu}^{\mathcal{B}} = C_{rad}\bar{\chi}$

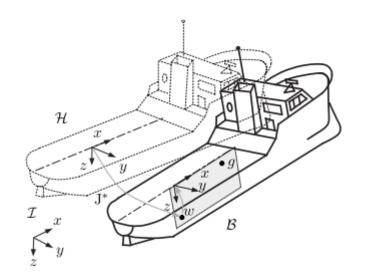
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Known values via hydrodynamic code (WAMIT):

$$\bar{\tau}_{rad}^{\mathcal{B}} = \mathbf{A}_{c}^{\mathcal{B}}(j\omega)\bar{\nu}^{\mathcal{B}} = \left[\frac{\mathbf{J}^{*T}\mathbf{B}(j\omega)\mathbf{J}^{*}}{j\omega} + \mathbf{J}^{*T}\mathbf{A}(j\omega)\mathbf{J}^{*}\right]\bar{\nu}^{\mathcal{H}}$$

Approx. model: (Gauss-Newton iter)

$$\hat{\tau}_{rad_{kl}}^{\mathcal{B}}(s) = \left[\mathbf{A}_{\infty_{kl}}s + \frac{\mathbf{P}_{kl}(s)}{\mathbf{Q}_{kl}(s)}\right]\nu_l(s)$$



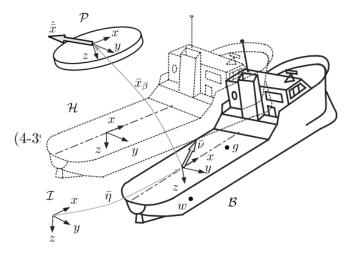
#4a: Dynamics – Pose dep. mass matrix

The kinetic energy of the platform is now found using the platform rigid body inertia matrix $M_{RBp} \in \mathbb{R}^{6\times 6}$ with the reduction matrices H from Eq. (4-34) and the velocity of the platform with respect to the inertial frame expressed in the platform frame $\mathcal{P}(\bar{V}_{0p}^{\mathcal{P}})$ as:

$$\mathcal{K}_{\text{platform}} = \frac{1}{2} (\bar{V}_{0p}^{\mathcal{P}})^T M_{RBp} \bar{V}_{0p}^{\mathcal{P}}$$
(Relative velocity) = $\frac{1}{2} (\bar{V}_{0b}^{\mathcal{P}} + \bar{V}_{bp}^{\mathcal{P}})^T M_{RBp} (\bar{V}_{0b}^{\mathcal{P}} + \bar{V}_{bp}^{\mathcal{P}})$
(ref. frame transf.) = $\frac{1}{2} (\bar{V}_{0b}^{\mathcal{B}} + \bar{V}_{bp}^{\mathcal{B}})^T \text{Ad}_{\mathcal{BP}}^T M_{RBp} \text{Ad}_{\mathcal{BP}} (\bar{V}_{0b}^{\mathcal{B}} + \bar{V}_{bp}^{\mathcal{B}})$
(notation) = $\frac{1}{2} (\bar{\nu} + \text{H}\dot{\bar{x}})^T \text{Ad}_{\mathcal{BP}}^T M_{RBp} \text{Ad}_{\mathcal{BP}} (\bar{\nu} + \text{H}\dot{\bar{x}})$

$$= \frac{1}{2} (\bar{\nu}^T + \dot{\bar{x}}^T \text{H}^T) \text{Ad}_{\mathcal{BP}}^T M_{RBp} \text{Ad}_{\mathcal{BP}} (\bar{\nu} + \text{H}\dot{\bar{x}})$$
(expand) = $\frac{1}{2} \bar{\nu}^T \underbrace{\text{Ad}_{\mathcal{BP}}^T M_{RBp} \text{Ad}_{\mathcal{BP}}}_{M_{vv}} \bar{\nu} + \bar{\nu}^T \underbrace{\text{Ad}_{\mathcal{BP}}^T M_{RBp} \text{Ad}_{\mathcal{BP}} \text{H}}_{M_{vc}} \dot{\bar{x}}$

$$+ \frac{1}{2} \dot{\bar{x}}^T \underbrace{\text{H}^T \text{Ad}_{\mathcal{BP}}^T M_{RBp} \text{Ad}_{\mathcal{BP}} \text{H}}_{M_{cc}} \dot{\bar{x}}$$



Where we've recognised the velocity of the ship expressed in the ship frame (only heave, roll and pitch) as $\bar{V}_{0b}^{\mathcal{B}} = \mathbf{H}^T \bar{\nu}$ and the velocity of the platform relative to the ship expressed in \mathcal{B} as $\bar{V}_{bp}^{\mathcal{B}} = \dot{\bar{x}}$. Now the total mass matrix for the platform dynamics becomes:

$$\mathbf{M}_{p} = \begin{bmatrix} \mathbf{M}_{cc} & \mathbf{M}_{vc} \\ \mathbf{M}_{vc}^{T} & \mathbf{M}_{vv} \end{bmatrix} = \begin{bmatrix} \mathbf{H}^{T} \\ \mathbf{I}_{6\times 6} \end{bmatrix} \mathbf{M}_{vv} \begin{bmatrix} \mathbf{H} & \mathbf{I}_{6\times 6} \end{bmatrix} \in \mathbb{R}^{9\times 9}$$
(4-40)

And the kinetic energy and massmatrix for the ship are found by the rigid body mass plus the infinite period added mass in the body frame (Eq. (4-32)):

$$\begin{aligned} \mathcal{K}_{\rm ship} &= \frac{1}{2} \bar{V}_{0b}^{\mathcal{B}^T} \mathbf{M}_s \bar{V}_{0b}^{\mathcal{B}} \\ &= \frac{1}{2} \bar{\nu}^T \mathbf{M}_s \bar{\nu} \end{aligned}$$

The total mass matrix is now (4-41) (platform) pose dependent

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#4b: Dynamics - Lagrange

Now using the velocity transform coupling the position variables to the quasi-velocities defined as:

(body fixed general velocities) $\bar{v} = S(\bar{q})\dot{\bar{q}}$ (euler angle rates) (4-44) Gives the Lagrangian \tilde{L} expressed in the derivative of the position variables (\bar{q}):

$$\tilde{L}(\bar{q},\bar{v}) = \frac{1}{2}\dot{\bar{q}}^{T}S^{T}(\bar{q})M(\bar{q})S(\bar{q})\dot{\bar{q}} - \mathcal{U}(\bar{q})$$

$$\frac{\partial\tilde{L}}{\partial\bar{v}} = S^{T}(\bar{q})\frac{\partial L}{\partial\bar{v}}$$

$$\frac{d}{dt}\left(\frac{\partial\tilde{L}}{\partial\bar{v}}\right) = \dot{S}(\bar{q})\frac{\partial L}{\partial\bar{v}} + S^{T}(\bar{q})\frac{d}{dt}\left(\frac{\partial L}{\partial\bar{v}}\right)$$

$$\frac{\partial\tilde{L}}{\partial\bar{q}} = \frac{\partial L}{\partial\bar{q}} + \frac{\partial^{T}(S(\bar{q})\dot{\bar{q}})}{\partial\bar{q}}\frac{\partial L}{\partial\bar{v}}$$
(4-45)

Now Lagrange's equations are found as:

$$\frac{d}{dt} \left(\frac{\partial \tilde{L}}{\partial \bar{v}} \right) - \frac{\partial \tilde{L}}{\partial \bar{q}} = S^{T} \tau$$
(4-46)
$$\mathbf{M}(\bar{q})\dot{\bar{v}} + \underbrace{\dot{\mathbf{M}}(\bar{q})\bar{v} - S^{-T}(\bar{q})\frac{1}{2}\frac{\partial^{T}\mathbf{M}(\bar{q})\bar{v}}{\partial \bar{q}}}_{\text{Multibody Coriolis terms}} + \underbrace{S^{-T}(\bar{q})\left(\dot{\mathbf{S}}^{T}(\bar{q}) - \frac{\partial^{T}(\mathbf{S}(\bar{q})\dot{\bar{q}})}{\partial \bar{q}} \right) \mathbf{M}(\bar{q})\bar{v}} + S^{-T}(\bar{q})\frac{\partial \mathcal{U}(\bar{q})}{\partial \bar{q}} = \tau$$
(4-47)

Coriolis terms identical to single body case

