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DOI

[10.1093/ptep/ptae011](https://doi.org/10.1093/ptep/ptae011)

Publication date

2024

Document Version

Final published version

Published in

Progress of Theoretical and Experimental Physics

Citation (APA)

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Pointing Calibration of GroundBIRD Telescope Using Moon Observation Data

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Received August 22, 2023; Revised January 17, 2024; Accepted January 19, 2024; Published January 22, 2024

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 Understanding telescope pointing (i.e. line of sight) is important for observing the cosmic microwave background (CMB) and astronomical objects. The Moon is a candidate astronomical source for pointing calibration. Although the visible size of the Moon (30') is larger than that of the planets, we can frequently observe the Moon once a month with a high signal-to-noise ratio. We developed a method for performing pointing calibration using observational data from the Moon. We considered the tilts of the telescope axes as well as the encoder and collimation offsets for pointing calibration. In addition, we evaluated the effects of the nonuniformity of the brightness temperature of the Moon, which is a dominant systematic error. As a result, we successfully achieved a pointing accuracy of 3.3'. This is one order of magnitude smaller than an angular resolution of 36'. This level of accuracy competes with past achievements in other ground-based CMB experiments using observational data from the planets.

Subject Index F10, F11, F14, F15

1. Introduction

Precise measurements of the cosmic microwave background (CMB) provide us with knowledge about the beginning of the universe, such as cosmic inflation prior to the Big Bang [1,2]. An important predictor of the cosmic inflation is the existence of primordial gravitational waves [3]. Primordial gravitational waves imprint weak odd-parity patterns (*B*-modes) in the CMB polarization [4–10]. Recent CMB experiments have focused on detecting *B*-modes from primordial gravitational waves, which are expected to be very faint. Therefore, it is important to establish calibration methods and obtain sufficient statistics by implementing a large number of detectors. In particular, the calibration of the line-of-sight of each detector (hereafter referred to as pointing) is essential. This is because time-ordered data (raw data) cannot be converted into map data (patterns in the sky coordinates) without pointing information. In previous studies, the primary calibration sources used were planets [11–13]. Given their small size compared with the large beam widths of CMB telescopes, they can be safely considered point sources. This is because we do not need to consider possible nonuniformities of them. For ground-based CMB experiments, the previous study achieved a ratio of pointing accuracy to their beam width of 0.13–0.14 [12,14]. This is because the pointing error is related to the beam width and leads to a smearing of the beam. The beam is defined as the telescope response in relation to the angle from the pointing direction. The angular resolution is represented by the beam width, which is the full width at half maximum (FWHM).

Jupiter and Saturn are popular pointing calibration sources for observations in the millimeter wavelength range. Their positions in the sky were calculated precisely using `astropy` [15,16]. Their visible sizes (0.2'~0.3') are sufficiently small compared with the beam width, which is typically 1' to 1° for CMB telescopes. The maximum elevation of the planets varies periodically on the order of 10 years, as shown in Fig. 1. Therefore, they cannot be used if the lowest elevation of the telescope is above them. This difficulty was encountered by our telescope, GroundBIRD [17], from 2019 to 2023 because the lowest elevation of the telescope is 60°. Therefore, a calibration source other than Jupiter or Saturn was required. Although Mars and Venus are other candidates for pointing calibration, it is difficult to use them because they are located close to the Sun. We excluded the observation region at a distance of $\pm 11^\circ$ from the Sun from the perspective of GroundBIRD, which we describe in Sect. 2. In addition, these planets (i.e. Jupiter, Saturn, Mars, and Venus) are not sufficiently bright for detection at high signal-to-noise ratios because of the beam width of GroundBIRD (i.e. 0.6°).

The Moon is a potential astronomical source for pointing calibration. GroundBIRD can frequently observe the Moon at least once a month. Its visible size is 30' (i.e. 0.5°), and its brightness temperature is high (~ 200 K). Although ideally point sources are used for pointing calibrations, the Moon is sufficiently bright for measurements with a high signal-to-noise ratio. Even though we used Jupiter which is the brightest of the point sources, telescopes whose beam width is sub-degrees need to integrate data to identify a position of it. Therefore, the measurement with the high signal-to-noise ratio introduces an advantage to reduce the effect of atmospheric fluctuation. Recently, in the CLASS experiment, the Moon was used as a pointing calibration source because the beam width was sufficiently large (1.5°) compared with the visible size of the Moon [18]. The beam width of GroundBIRD (36') is slightly larger than the Moon's visible size, and the dynamic range of the detector response is sufficiently large to

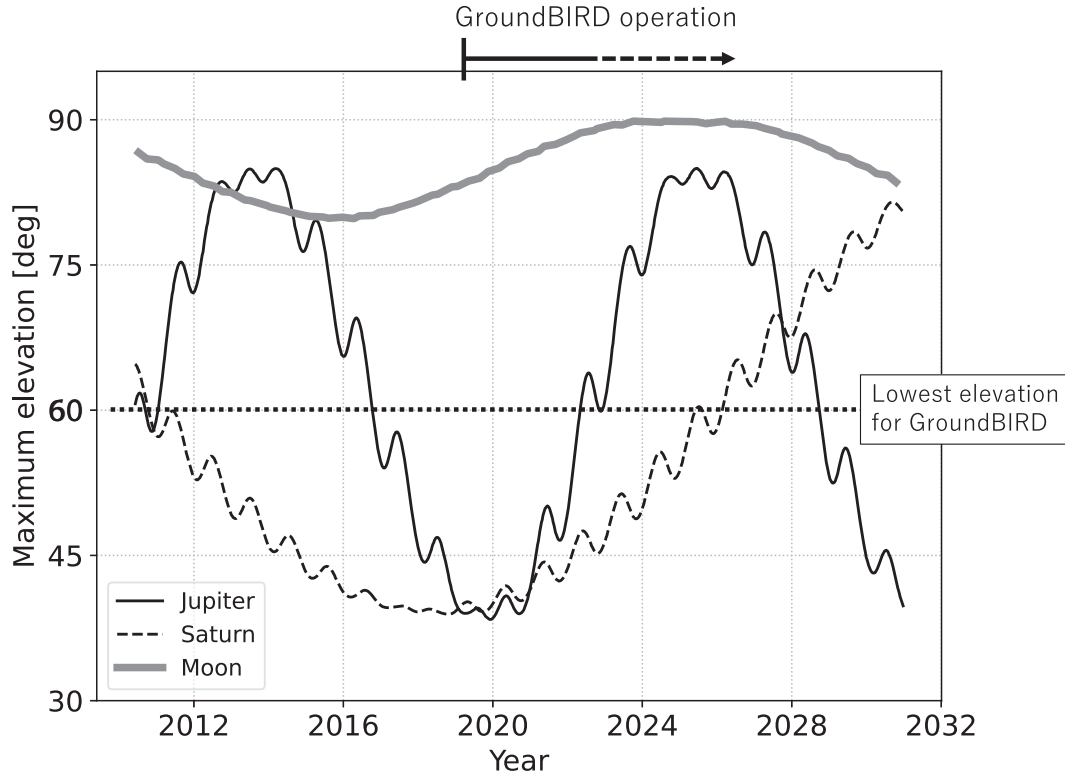


Fig. 1. Maximum elevations of Jupiter, Saturn, and the Moon between 2011 and 2031 at the Teide Observatory in the Canary Islands, Spain. The orbital periods of Jupiter and Saturn are 12 years and over 20 years, respectively. In the years around 2020, their maximum elevations were both lower than 60° . However, the Moon has always had a high elevation throughout these years.

measure the Moon signal. Therefore, Moon observational data can be used for pointing calibration, as demonstrated in the present study.

The methodology of pointing calibration using the Moon has not been well established yet. For instance, there has not been any systematic error study in 145 GHz. In this study, we demonstrate pointing calibration using the Moon. We also performed a systematic error study. Because the Moon is brighter compared with the point-like planets, methods established in this study should be useful for other experiments whose beam width is sub-degrees as with GroundBIRD. In Sect. 2, the GroundBIRD telescope used in this study is described. In Sect. 3, we describe the Moon's angular response. In Sect. 4, the methodology used to extract the central position of the Moon using each detector is described. In Sect. 5, we define the pointing calibration model and present the calibration results. In Sect. 6, we discuss the systematic uncertainties. Finally, in Sect. 7, our conclusions are presented.

2. GroundBIRD telescope

GroundBIRD is a ground-based CMB polarization experiment (Fig. 2). Our telescope is located at the Teide Observatory in the Canary Islands, Spain. Its longitude, latitude, and altitude are $28^\circ 18'$ N, $16^\circ 30'$ W, and 2,400 m, respectively. The telescope rotates continuously in the azimuthal direction at a fixed elevation. The maximum speed of the azimuth rotation is 20 revolutions per minute (RPM). This rapid scan modulation mitigates the effects of atmospheric fluctuations [17]. The region above an elevation of 60° can be observed, which allows for the observation of up to 45% of the entire sky (declination range of -2° to 58°).

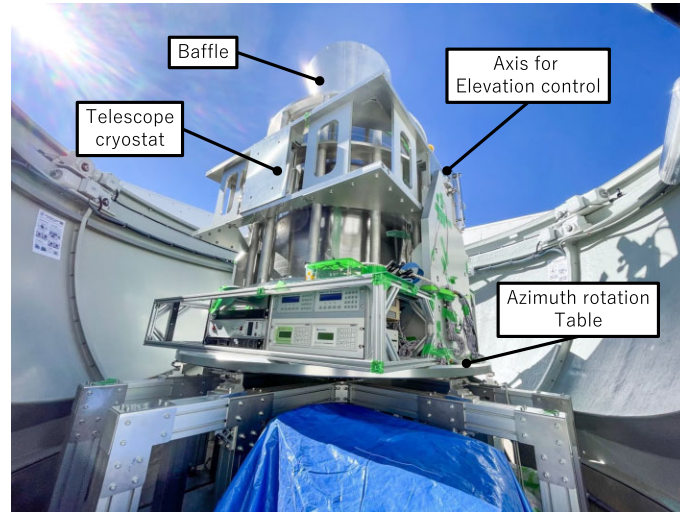


Fig. 2. GroundBIRD telescope. The telescope cryostat is placed on the azimuth rotation table, and it is rotated in the azimuth with a fixed elevation.

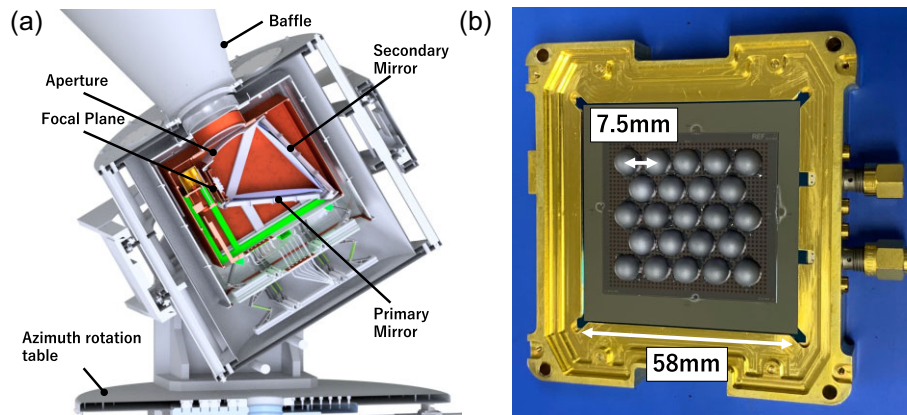


Fig. 3. (a) Cross-sectional illustration of the GroundBIRD telescope. (b) Photograph of the detector array module used in this research. There are 23-antenna-coupled KIDs behind each silicon lens.

Figure 3(a) shows a cross-sectional illustration of the GroundBIRD telescope. The sky signal enters the cryostat through a polyethylene window [19] and is focused onto focal-plane detectors using a Mizuguchi–Dragone dual reflector [20–22]. The signals were detected by antenna-coupled Kinetic Inductance Detectors (KIDs) [23] using silicon lenses in front of each antenna. The responses of each KID were measured as variations of the resonant phase and frequency [24,25]. The data used in this study were obtained with a 1 kHz sampling rate using a 23-KID array module for the 145 GHz band, as shown in Fig. 3(b) [26,27].¹ They are maintained at approximately 280 mK.

According to the simulation study by the CST Microwave Studio, the beam width of the telescope is $36'$ (i.e. 0.6°) [28]. Its ellipticity is at most 1% according to this simulation.² As shown

¹The module for the 145 GHz band was only installed for the commissioning observations in 2021–2022. We fully upgraded it in May 2023. They were fabricated at the Delft University of Technology and the Netherlands Institute for Space Research in the Netherlands.

²In Ref. [28], we confirmed this simulation by laboratory measurements with a small aperture setup ($\sim 1/4$ size) in a different frequency (90 GHz). We did not have sufficient space in the laboratory for a full-sized aperture. We did not have a setup for 145 GHz with sufficient signal-to-noise ratio. These are

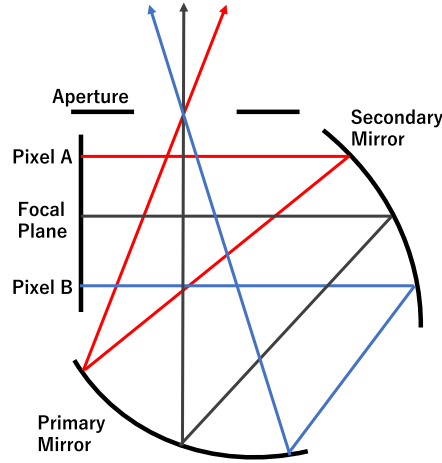


Fig. 4. The pointing of each detector pixel is dependent on its location in the focal plane. The aperture diameter is 220 mm. We use a Mizuguchi–Dragnone dual reflector which is a combination of a primary paraboloid and a secondary hyperboloid mirror. For the two mirrors, the diameters at the optical area are the same. The diameters for the minor axis and the major axis are 360 mm and 490 mm, respectively.

Table 1. Optical specifications of the GroundBIRD telescope.

Field of view	$\pm 11^\circ$
Typical collimation interval between each detector pixel	51'
Beam width	36'
Beam ellipticity	$< 1\%$
Precision of elevation encoder	4.0''
Precision of azimuth encoder	3.4''
Pointing accuracy (requirement)	4.7'

in Fig. 4, each detector in the focal plane points to a different position in the sky. The angular interval between the collimations of each detector in the sky is typically 51'. Therefore, the field of view of the telescope is $\pm 11^\circ$. The elevation and azimuth of the telescope are monitored using a laser rotary encoder (R-1SL, Canon) and a magnetic rotary encoder (ERM220, HEIDENHAIN), respectively. Their precisions are 4.0'' and 3.4'' for the elevation and azimuth, respectively [29]. Table 1 lists the optical specifications of the GroundBIRD telescope.

In other CMB experiments, the ratio of pointing accuracy to their beam width is 0.13–0.14 [12,14]. In this case, the impact on the B -modes power measurement was estimated to be $\sim 10^{-4} \mu\text{K}^2$ at a multipole of 100, which approximately corresponds to a power at the tensor-to-scalar ratio (r) of 0.001 [30]. This is two orders of magnitude lower than our sensitivity with three years of observations ($r \approx 0.1$). Based on this knowledge, we set the required pointing accuracy to 4.7' (i.e. $36' \times 0.13$) in this study.

3. Angular response to the Moon

In the millimeter wavelength range, thermal radiation from the Moon's surface is dominant compared with reflected sunlight. The brightness of the Moon signal in Kelvin (T_{moon}) is

the reasons for the small aperture and the different frequency. Validation using other planets such as Jupiter is the subject of a future study.

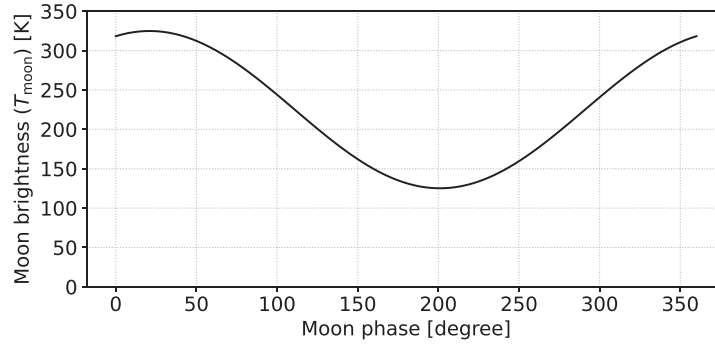


Fig. 5. Brightness of the Moon signal as a function of the Moon phase for 145 GHz. The Moon phase at $0'$ corresponds to the full Moon.

modeled using the Moon phase (ψ), as follows [31]:

$$T_{\text{moon}} = 225 \left\{ 1 + \frac{0.77}{\sqrt{1 + 2\delta + 2\delta^2}} \cos \left(\psi - \arctan \frac{\delta}{1 + \delta} \right) \right\}, \quad (1)$$

$$\delta \equiv 0.3\lambda,$$

where λ is the wavelength in millimeters. Figure 5 shows the brightness temperature at the Moon's surface at 145 GHz ($\lambda \sim 2$ mm) as a function of the Moon's phase. The timing of the maximum brightness is delayed compared with the full Moon owing to a delay in the temperature variation at the Moon's surface. The maximum and minimum Moon brightness temperatures were 325 and 125 K, respectively.

From the Earth, the visible brightness (T_{eff}) at the edge of the Moon is different from T_{moon} because of the different refractive indices of the Moon's surface. Considering the Moon as a black body sphere covered by a dielectric material layer [32], we obtained the angular distribution of the visible brightness, as shown in Fig. 6 (left). This is almost a top-hat distribution with a diameter of $30'$. The mean of T_{eff} within the visible range of the Moon is equal to the brightness temperature (T_{moon}). Outside the Moon, T_{eff} corresponded to the CMB temperature (2.725 K). For real observations, the measured signals were convolved with the beam. We used an angular response model convolved with T_{eff} and a Gaussian beam, as shown in Fig. 6 (right). The uncertainties owing to the nonuniformity of T_{moon} and the beam shape are discussed in Sect. 6.

4. Reconstruction of the Moon's position

We performed 19 observations of the Moon using an azimuth rotation scan at 10 RPM with a fixed elevation at $70'$ between 8 February 2022 and 18 April 2022, as listed in Table 2. The rotation of the Earth then leads the Moon to form a map through telescope scans. With regards to the Moon observations, we used time-ordered data (TOD) for each detector as well as telescope encoder data. The duration of each observation was 38–60 min. One of the 23 detectors was not used because its resonant frequency was outside our readout bandwidth. Prior to the analysis, we made corrections for the microwave phase delay owing to the cable length and detector responses [33].

We selected good-quality data from each detector for each observation based on the following criteria:

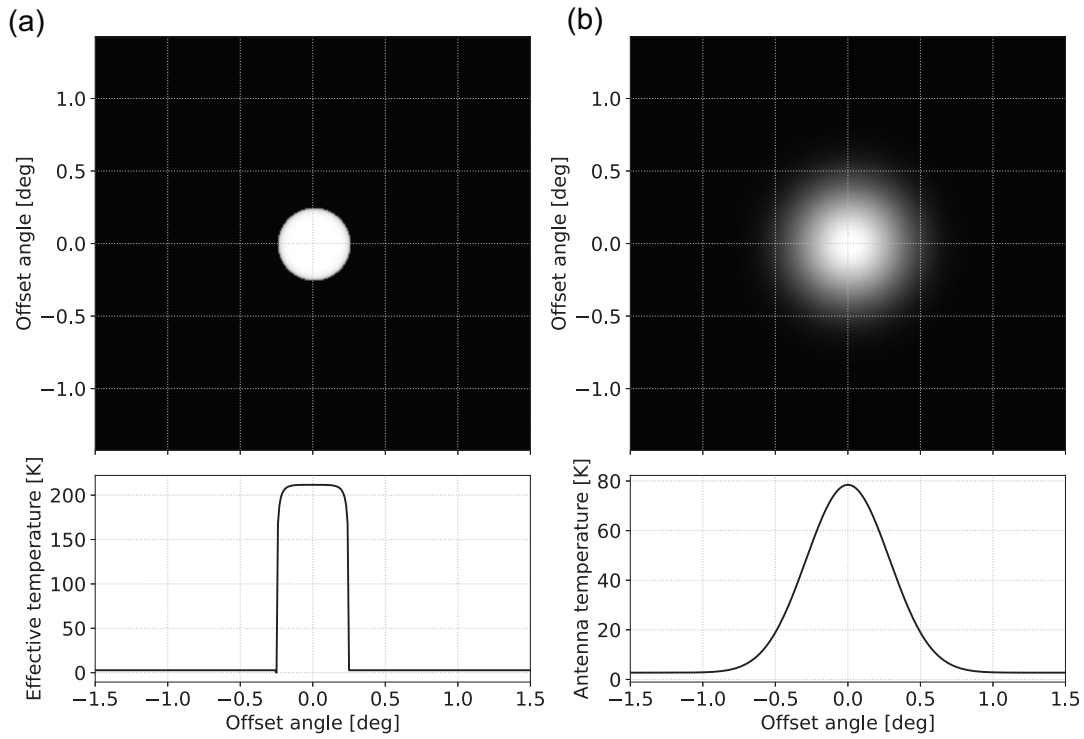


Fig. 6. (a) Distribution of visible brightness temperature around the Moon from the Earth. (b) Convolved distribution with the beam modeled using the simple Gaussian (beam width of 36'). The plots in the bottom row present a cross-sectional view.

Table 2. Information on the Moon observations used in this study.

Observation date (UTC)	Duration [min]	Ascent or descent	Azimuth [deg]	Elevation [deg]
8/2/2022 17:11–18:11	60	ascent	106–126	61–73
8/2/2022 20:14–21:01	47	descent	238–253	63–72
14/2/2022 22:12–22:50	38	ascent	95–103	64–72
15/2/2022 22:56–23:56	60	ascent	101–120	62–74
16/2/2022 23:48–24:48	60	ascent	110–133	61–72
17/2/2022 02:36–03:36	60	descent	224–247	61–72
7/3/2022 15:29–16:29	60	ascent	115–146	65–75
7/3/2022 18:01–19:01	60	descent	230–252	61–72
8/3/2022 15:47–16:47	60	ascent	100–117	62–74
9/3/2022 16:30–17:30	60	ascent	93–104	62–75
9/3/2022 19:50–20:43	52	descent	258–268	62–73
10/3/2022 20:42–21:36	54	descent	266–273	61–73
5/4/2022 17:41–18:28	47	descent	255–265	63–73
6/4/2022 15:26–16:26	60	ascent	91–103	65–78
6/4/2022 18:34–19:26	53	descent	264–272	62–73
7/4/2022 19:25–20:23	58	descent	269–275	61–73
8/4/2022 16:50–17:50	60	ascent	86–92	62–75
8/4/2022 20:16–21:16	60	descent	269–275	60–73
12/4/2022 20:28–21:28	60	ascent	114–140	62–73

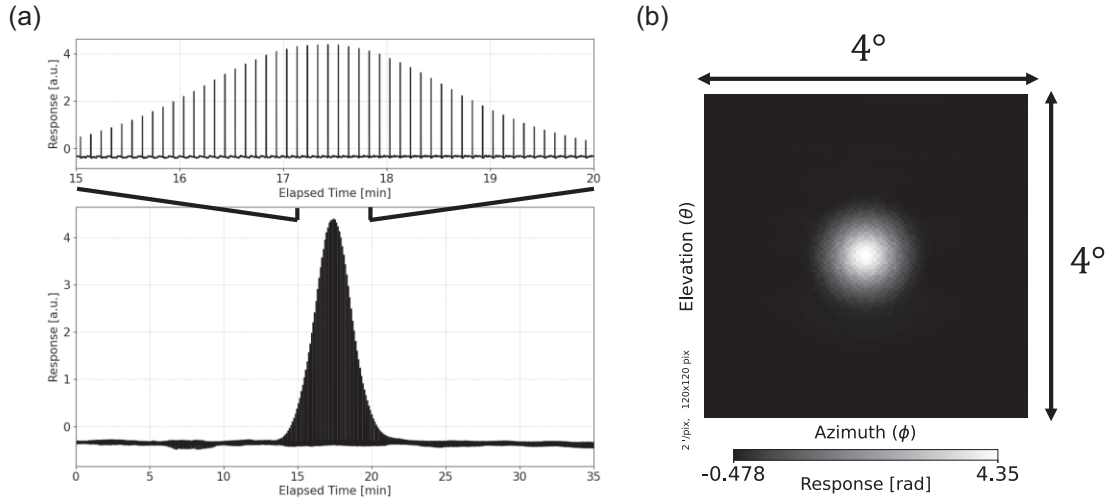


Fig. 7. (a) Detector response as a function of the elapsed time. (b) Reconstructed Moon image using the Moon-centered coordinates for azimuth (horizontal axis) and elevation (vertical axis). This image is created using `healpy` with the parameter $N_{\text{side}} = 1024$. The baseline fluctuations due to atmospheric radiation are subtracted as described in the text.

- The intervals of the resonant frequencies for each detector must be greater than 0.5 MHz for the same observation. This condition eliminates potential crosstalk among the detectors. We did not use 2 of the 22 detectors due to this criterion.
- Data within an elevation of 1.5° from the Moon center must be obtained continuously for each detector, i.e. each detector must observe the outside region of the Moon as well as the Moon itself.
- The Moon signal must be within the dynamic range of each detector. The detector response is proportional to $\tan(\frac{\psi}{2})$ where ψ is a phase of fed microwaves for the detector readout [34]. Thus, we observe a rapid jump in the detector response when the signal intensity is higher than the dynamic range. We identified it as rapid variations of detector response (more than 30 times) within the visible size of the Moon in the azimuthal scan. We did not use 1 of the 20 detectors based on this criterion.

On the basis of these criteria, 345 Moon observation data were selected.

Figure 7(a) presents the TOD for a single observation. The overall shape of the bottom figure traces the elevation pattern of the Moon, and the spikes in the top figure trace the azimuth pattern of the Moon. The time information between the detector data and telescope encoder data was synchronized by distributing a common clock signal. Thus, a map of the elevation and azimuth coordinates for each detector was constructed. For each TOD point, we can calculate the true position of the Moon using `astropy`. Using the true position of the Moon, we further converted the map to Moon-centered coordinates, the axes of which comprise the elevation and azimuth from the Moon's center. One of the reconstructed Moon images in the Moon-centered coordinates is shown in Fig. 7(b). In this study, we subtracted the baseline offset for each azimuthal scan to eliminate the effects of atmospheric fluctuations. The baseline offset should be calculated from the detector response at away from the Moon. Therefore, each baseline offset was calculated as the mean value of the signals between 1.5° and 2.0° from the observed center of the Moon. Figure 8 presents the reconstructed Moon images obtained from each detector for one of the observations. The locations of the observed Moon images are

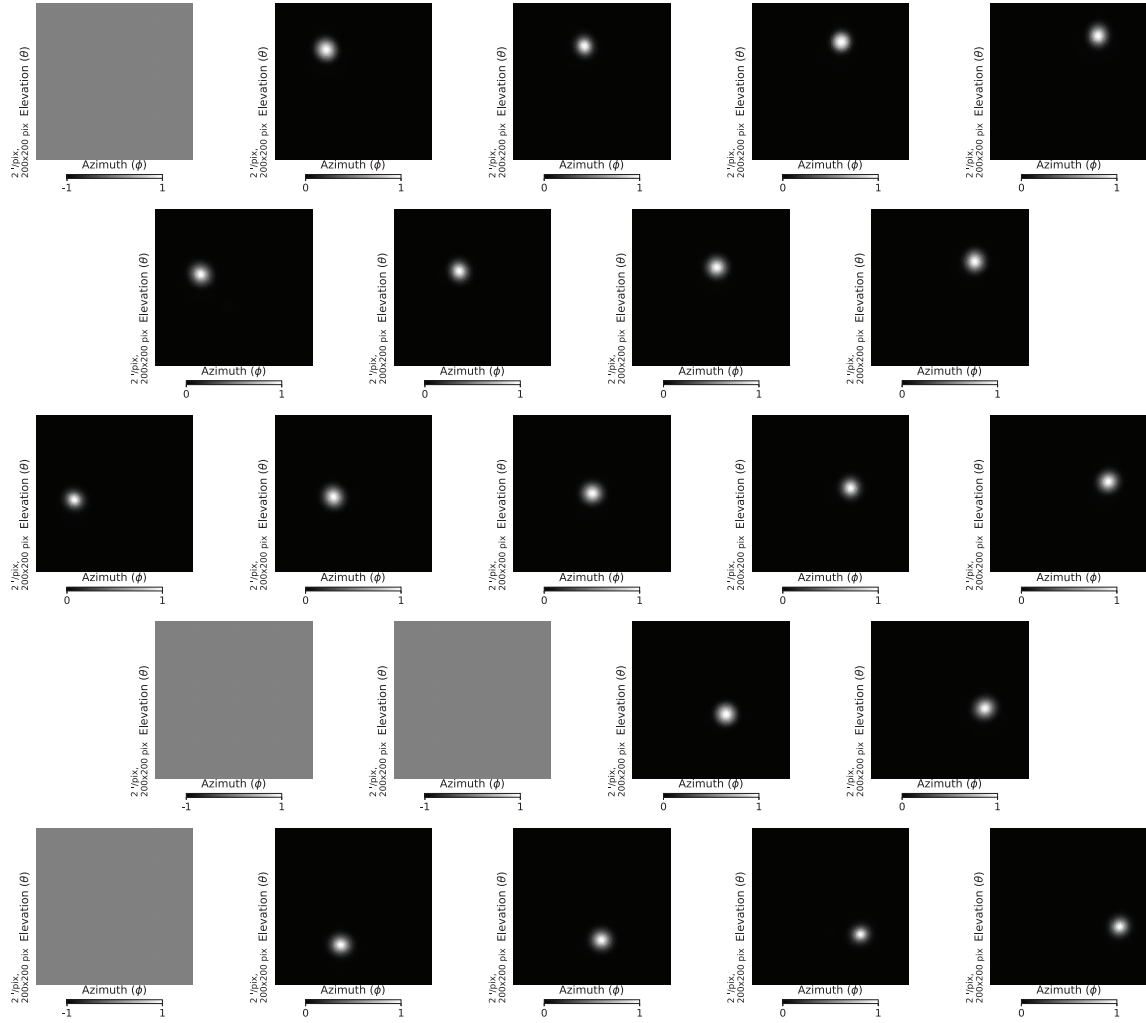


Fig. 8. Reconstructed Moon images for each detector in the Moon-centered coordinates. The highest signal is normalized to 1. The positions of the plots indicate the locations of each detector pixel on the focal plane. The data of four detectors are not used; they are indicated by the gray squares. The locations of the Moon images in each detector image are different because each detector has a different collimation offset. The angular ranges in all plots are $400'$ times $400'$. These images are created using `healpy` with the parameter $N_{\text{side}} = 1024$.

different for each detector because each detector has a different collimation point in the sky, as described in Sect. 2.

Based on the Moon's response defined in Sect. 3, we extracted the Moon's center positions for each detector and for each observation, by performing an unbinned likelihood fit for the data within a radius of 1.5° from the Moon's center. The fitting parameters were the elevation and the azimuth of the Moon center, detector gain, beam width, and constant offset as the background residual. Figure 9 presents an example of a result obtained by the fitting procedure. The difference in the beam between the data and the model has a negligible effect on the pointing calibration, which is discussed in Sect. 6. The extracted Moon positions for each detector and each observation comprised 345 samples and were used to calibrate the pointing model described in the next section. The average collimation interval (the distance between neighboring detectors at the sky) was $52' \pm 1'$, which is consistent with the design ($51'$).

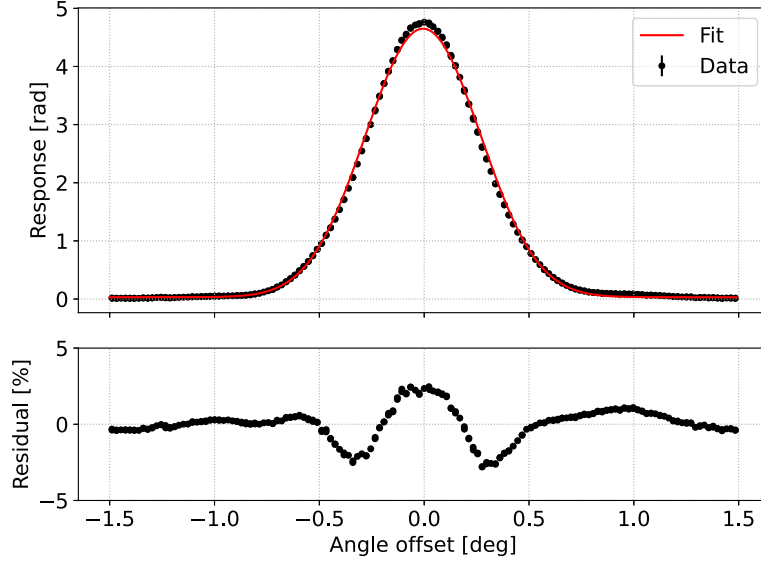


Fig. 9. Data from a Moon scan fitted with a Gaussian beam in an azimuthal scan. The non-Gaussianity of the beam shape leads to approximately 3% residuals because we use the simple Gaussian model in the fit. This effect is considered to be a systematic error, as described in Sect. 6.

5. Pointing calibration

The tilts of the elevation and azimuth axes produce pointing shifts in terms of the elevation (θ) and azimuth (ϕ), as shown in Fig. 10. Therefore, we modeled the pointing shifts as follows [35]:

$$\left\{ \begin{array}{l} \delta_{\theta}(\phi) = \delta_{\text{NS}} \cos \phi + \delta_{\text{EW}} \sin \phi, \\ \cos \theta \delta_{\phi}(\theta, \phi) = \delta_{\text{NS}} \sin \theta \sin \phi - \delta_{\text{EW}} \sin \theta \cos \phi, \end{array} \right. \quad \begin{array}{l} (2a) \\ (2b) \end{array}$$

where δ_{NS} is the tilt in the north–south direction and δ_{EW} is the tilt in the east–west direction. Figure 11 shows pointing shifts as a function of the true elevation and the true azimuth for $\delta_{\text{NS}} = -1.8'$ and $\delta_{\text{EW}} = -2.8'$ calculated from Eqs. (2a) and (2b).

In addition to the tilts of the axes, the encoder and collimation offsets of each detector cause constant pointing shifts, as shown in Fig. 8. We define the parameters for the encoder offsets in terms of elevation (θ_e) and azimuth (ϕ_e). We also define the collimation offsets of each detector for the elevation (θ_i) and azimuth (ϕ_i). Here, i is the index for each detector. Using these parameters, the pointing model for the data is defined as follows:

$$\left\{ \begin{array}{l} \theta^{\text{model}} = \theta^{\text{data}} - \theta_e - \theta_i - \delta_{\theta}(\phi^{\text{true}}), \\ \phi^{\text{model}} = \phi^{\text{data}} - \phi_e - \phi_i - \delta_{\phi}(\theta^{\text{true}}, \phi^{\text{true}}), \end{array} \right. \quad \begin{array}{l} (3a) \\ (3b) \end{array}$$

where θ^{data} and ϕ^{data} are the reconstructed Moon positions for each detector for each observation (i.e. the reconstructed data samples in Sect. 4), and θ^{true} and ϕ^{true} are the true positions of the Moon, which are calculated using `astropy`.

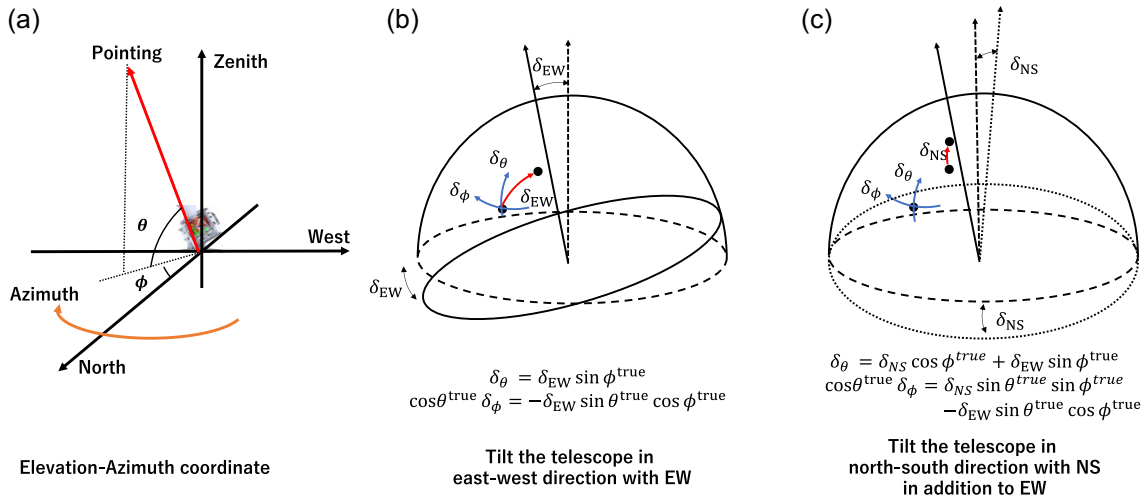


Fig. 10. (a) Definition of the elevation (θ) and azimuth (ϕ). (b) Schematics of the axis tilt to the east–west direction. (c) Schematics of the axis tilt to the east–west direction and north–south direction.

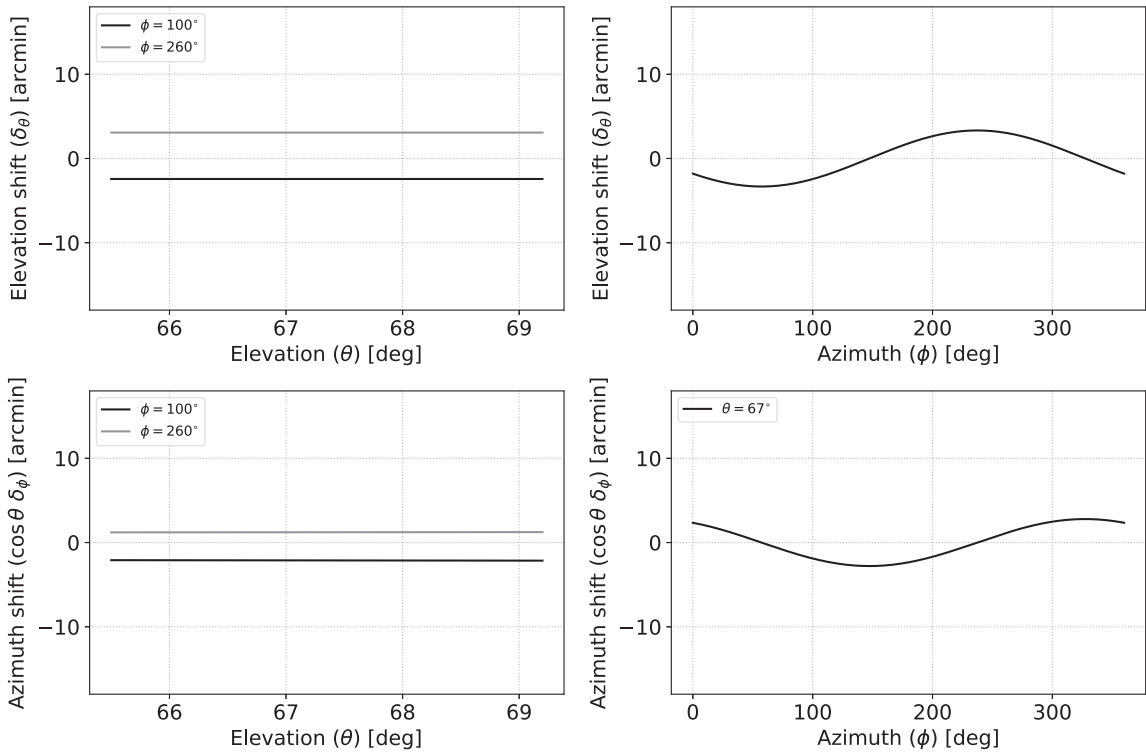


Fig. 11. Pointing shifts as a function of the elevation or azimuth with δ_{NS} of $-1.8'$ and δ_{EW} of $-2.8'$. The shift in the elevation depends on the azimuth pointing. In contrast, the shift in the azimuth depends on both the elevation and the azimuth pointing.

These parameters were optimized to minimize the angular distance (dl) between the true and model positions; i.e. $dl^2 = (\theta^{model} - \theta^{true})^2 + [\cos \theta^{true}(\phi^{model} - \phi^{true})]^2$. Although dl is calculated under the approximation of the small angular distance, this calculation method makes a negligible effect ($\leq 0.01''$). The optimized tilt angles were $\delta_{NS} = -1.8'$

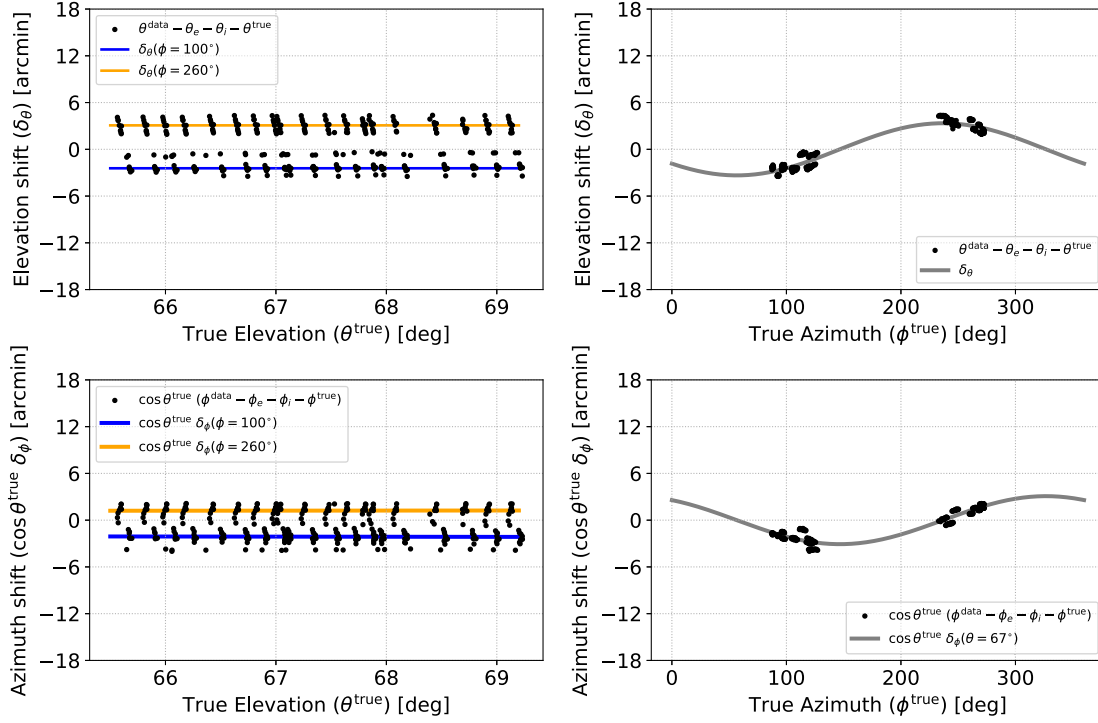


Fig. 12. Pointing shifts as a function of the elevation or azimuth of the true Moon position. The points comprise reconstructed data calculated using Eqs. (4a) and (4b). The lines were calculated using Eqs. (2a) and (2b) with the same values for δ_{NS} and δ_{EW} .

and $\delta_{EW} = -2.8'$. Using the extracted encoder and collimation offsets, we calculated the pointing shifts as follows:

$$\begin{cases} \delta_\theta(\phi^{\text{true}}) = (\theta^{\text{data}} - \theta_e - \theta_i) - \theta^{\text{true}}, & (4a) \\ \cos \theta^{\text{true}} \delta_\phi(\theta^{\text{true}}, \phi^{\text{true}}) = \cos \theta^{\text{true}} [(\phi^{\text{data}} - \phi_e - \phi_i) - \phi^{\text{true}}]. & (4b) \end{cases}$$

Figure 12 presents the pointing shifts as functions of the true Moon position at that elevation or azimuth. For comparison, we overlay the lines calculated using Eqs. (2a) and (2b). These values were consistent within the required precision (4.7').

In the CMB analysis, we calculate the pointing based on the optimized model. Although θ^{true} and ϕ^{true} are unknown in Eqs. (3a) and (3b), accurate pointing can be calculated iteratively with sufficient precision [36] instead of solving Eqs. (3a) and (3b) exactly. The following formulae are used:

$$\begin{cases} \theta_i^{(0)} = \theta^{\text{data}} - \theta_e - \theta_i, & (5a) \\ \phi_i^{(0)} = \phi^{\text{data}} - \phi_e - \phi_i, & (5b) \end{cases}$$

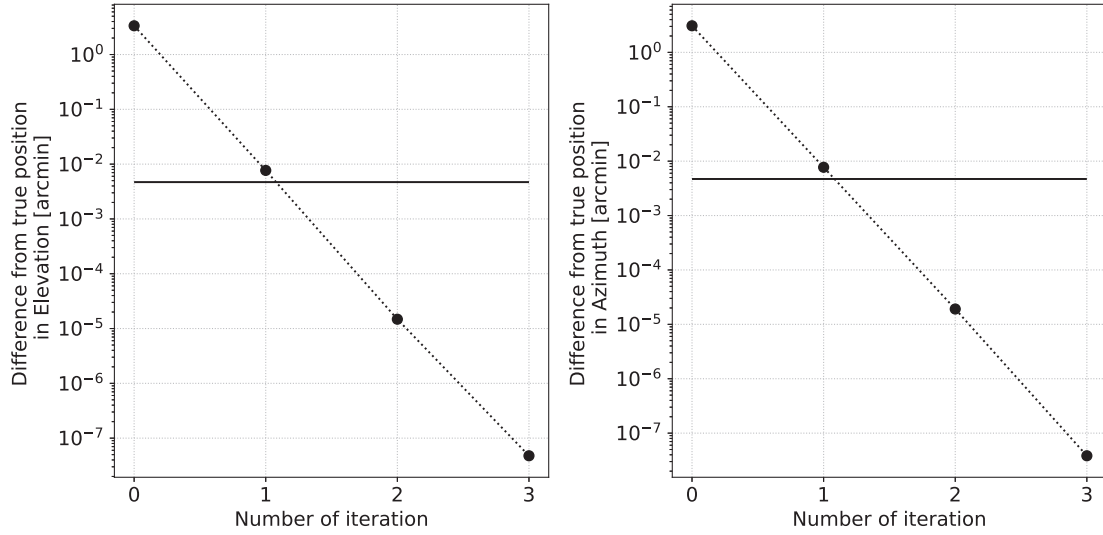


Fig. 13. Pointing differences between a calibrated position with Eqs. (5a–6b) and the true position as a function of the number of iterations for the elevation and azimuth. The elevation angle is assumed to be 67° , which is a typical elevation angle. The solid lines are $1/1,000$ of the requirement. We realize a sufficient precision at $n \geq 2$.

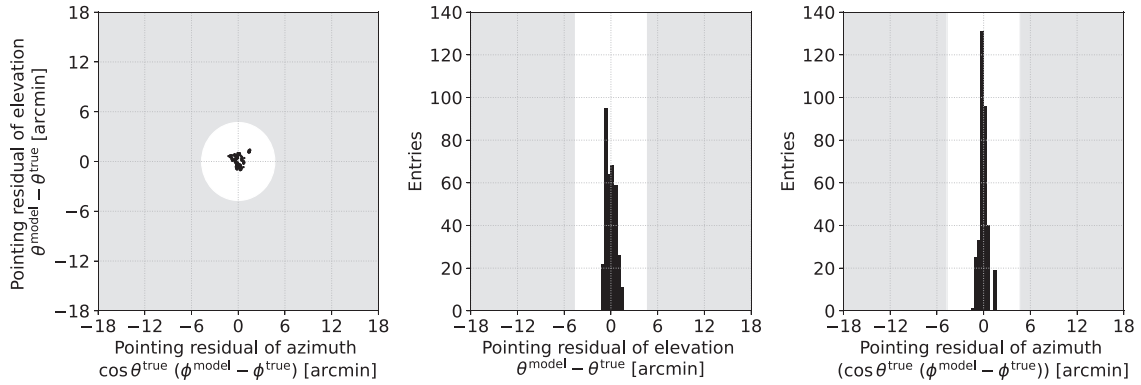


Fig. 14. Residuals of the calibrated pointing for all 345 samples from each true position. The dynamic range of the elevation and azimuth axes corresponds to the beam width of GroundBIRD ($36'$). They are within the requirement, which is indicated by the unshaded regions.

$$\left\{ \begin{array}{l} \theta_i^{(n)} = \theta^{\text{data}} - \theta_e - \theta_i - \delta_\theta \left(\phi_i^{(n-1)} \right), \\ \phi_i^{(n)} = \phi^{\text{data}} - \phi_e - \phi_i - \delta_\phi \left(\theta_i^{(n-1)}, \phi_i^{(n-1)} \right), \\ n = 1, 2, 3, \dots, \end{array} \right. \quad \begin{array}{l} (6a) \\ (6b) \end{array}$$

where n indicates the n -th iteration. Using Eqs. (6a) and (6b) and the extracted parameters, we numerically calculated the differences from the true position for each n . We could achieve sufficient precision at $n \geq 2$, as shown in Fig. 13. We conclude that two iterations should be performed in the calculation of the pointing to reduce this effect on the pointing calibration to a negligible level.

Figure 14 shows the residuals of the Moon position as reconstructed from the true positions for all 345 samples. These values were within the range of our requirements. Their root mean

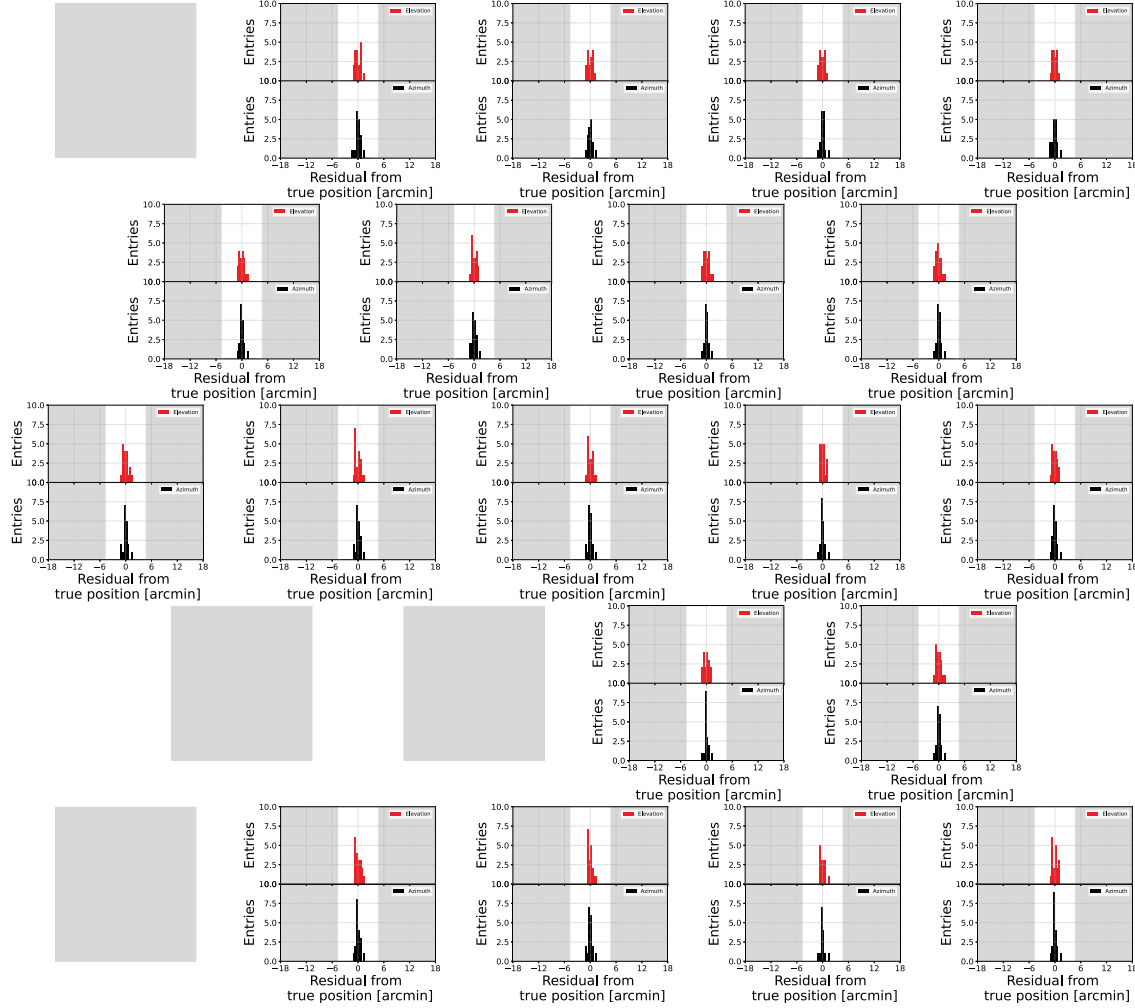


Fig. 15. Pointing residuals of the elevation, $\theta^{\text{model}} - \theta^{\text{true}}$, and azimuth, $\cos\theta^{\text{true}}(\phi^{\text{model}} - \phi^{\text{true}})$, for each detector. The mean values for each detector are consistent with zero.

squares were $0.6'$ for the elevation and $0.5'$ for the azimuth times $\cos\theta$. Similar plots for each detector are shown in Fig. 15. Their mean values were consistent with zero; i.e. no bias was found.

6. Systematic uncertainty

The uncertainty associated with the beam shape is estimated by changing the response model of the fit. We compared the results of the baseline analysis with a model that uses a fixed beam width as the design, one that uses an elliptical Gaussian beam with extracted ellipticity of $2.6 \pm 1.0\%$ from the data, and one that uses a polynomial (i.e. nonlinear) responsivity. The pointing differences among the beam models are less than $0.26'$. Note that the assumed ellipticity of the systematic error study is larger than that of the simulation study ($<1\%$). The uncertainties in the Moon's position owing to the time constant of the detector and the astronomical calculations are $0.14'$ and $0.31'$ [37], respectively. The uncertainties in the elevation and azimuth encoders are $0.066'$ and $0.057'$, respectively, as described in Sect. 2. A possible thermal effect due to ambient temperature change is estimated by comparing two calibration results at high temperature (average of 10°C) and low temperature (average of 3°C).

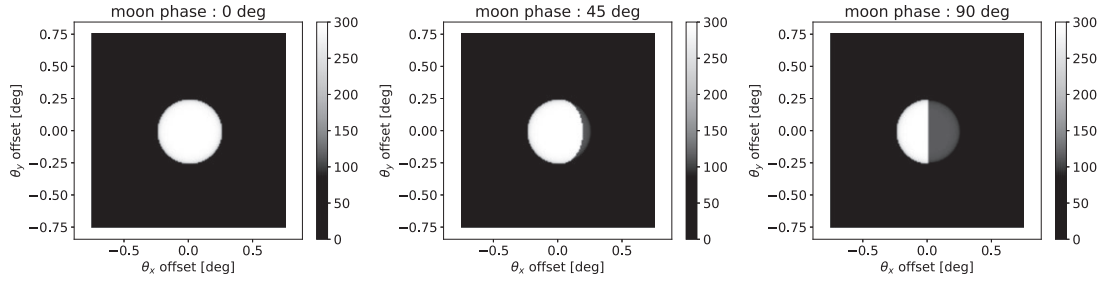


Fig. 16. Distribution of the Moon brightness temperature, which varies with the Moon phase, where θ_x is the direction of the phase shift, and θ_y is perpendicular to it. These models are used only for the systematic error study. The white and gray regions correspond to the brightness temperatures of 325 K and 125 K, respectively.

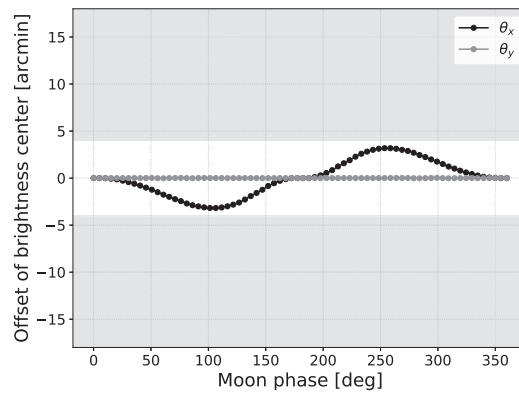


Fig. 17. Offset of the brightness center as a function of the Moon phase, where θ_x and θ_y are the directions of the offsets, as shown in Fig. 16. The maximum offset found here is smaller than the requirement, which is indicated by the unshaded area.

We found a difference of $0.4'$, while the statistical fluctuation is dominant in this comparison. A possible mechanical variation associated with the scan is estimated from the difference between the residual from the pointing model at azimuth $<180'$ and that at azimuth $>180'$. It was $0.057'$.

The systematic uncertainty is driven by the nonuniformity of the Moon's brightness temperature. We estimated this systematic uncertainty using a simulation based on simplified temperature distributions for each Moon phase, as shown in Fig. 16. Considering Eq. (1), we set the maximum and minimum temperatures in the bright and shaded regions, respectively. This assumption yielded the most conservative systematic error. Using the same method as that used for the real data analysis, we extracted the central position of the Moon for the simulation data. Figure 17 presents the difference between the extracted center position and the input position in the simulation. The maximum difference ($3.2'$) was assigned to the systematic error. Although there are variations of brightness on the Moon that are not due to the Moon phase (e.g. Ref. [38]), their effects are much smaller than those that are due to the Moon phase. It would be worthwhile to consider them in a future study.

Table 3 presents a summary of the systematic uncertainties. We calculate the square root of the quadrature sum of all errors. In total, we assigned a systematic error of $3.2'$ to the pointing calibration using the Moon.

Table 3. Systematic uncertainties in pointing.

Source	[arcmin]
Beam shape	2.6×10^{-1}
Moon position	3.4×10^{-1}
Elevation encoder	6.6×10^{-2}
Azimuth encoder	5.7×10^{-2}
Ambient temperature	4.0×10^{-1}
Mechanical variation	5.7×10^{-2}
Nonuniformity temperature	3.2
Total	3.2

7. Conclusions

We calibrated the pointing of the GroundBIRD telescope based on Moon observation data. The angular response of the Moon was modeled with respect to the Moon phase as well as the beam of the telescope. We modeled the pointing shift due to the tilts of the elevation and azimuth axes as well as the collimation offsets and the encoder offsets. We obtained the optimized pointing model using them. The residuals of the calibrated pointing from the true pointing were less than $1.6'$, whereas our requirement was $4.7'$. The root mean squares were $0.6'$ and $0.5'$ for the elevation and azimuth, respectively. Consequently, we successfully achieved an uncertainty of $3.3'$ including all systematic uncertainties, which is lower than our requirement. For the CMB telescope, this was the first attempt at using pointing calibration with the Moon for a beam width of $36'$ against the Moon's angular size of $30'$. Thus, we realized sufficient pointing precision.

We can frequently observe the Moon at high elevations. The orbital period of the Moon (monthly) is much shorter than that of planets (~ 10 years). In addition, the observation of the Moon with the high signal-to-noise ratio allows us to perform the unbinned likelihood fit to extract the center position of the Moon. It introduces the advantage of reducing the effect of atmospheric fluctuation. In addition, we avoid degradation of the angular resolution of detector response due to the binning of the data for the astronomical point sources. We discussed the systematic uncertainties related to the Moon. The nonuniformity of the Moon's brightness was assumed as the most conservative case. Nevertheless, it was one order of magnitude lower than the beam width. In conclusion, the established method in this study is applicable to other CMB telescopes whose beam width is sub-degrees.

Acknowledgements

This work is supported by Japan Society for the Promotion of Science (JSPS) KAKENHI under grant numbers JP21H04485, JP20KK0065, and JP21H04485, JSPS Core-to-Core Program number JPJSCCA20200003, and JSPS Bilateral Program numbers JPJSBP120219943 and JPJSBP120239919. Y.S. also acknowledges grant no. JP21J20290. K.L. and E.W. also acknowledge grant no. 2022R1A2B5B02001535 by the National Research Foundation of Korea (NRF).

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