Piezoelectric deformable mirror with adaptive multiplexing control

A. N. Simonov, S. Hong, and G. Vdovin
Delft University of Technology,
Electronic Instrumentation Laboratory,
Mekelweg 4, 2628 CD Delft, The Netherlands

Abstract. We describe a simple and efficient implementation of adaptive multiplexing control for high-order piezoelectric deformable mirrors. The relatively high capacitance of piezoelectric actuators allows the electrical charge to be stored in a disconnected actuator, retaining its displacement while the other actuators can be addressed. Adaptive multiplexing, consisting of selective addressing of only those actuators that need to change their elongation in the current cycle, improves the mirror performance and simplifies the driver electronics. In experiment, a 12-channel prototype of a deformable mirror with multiplexing control has been characterized. At appropriate update rates with a fixed set of control signals, the shape of the deformable mirror remains nearly constant. A surface displacement error does not exceed \( \sim \lambda/100 \) rms at a multiplexing frequency of 700 Hz with a full interactuator stroke of \( \sim 2 \mu \)m. © 2006 Society of Photo-Optical Instrumentation Engineers.

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Deformable mirrors (DM) are the key components of adaptive optics (AO) systems.1,2 A number of scientific,3,4 medical,5,6 and industrial2,5,7 applications require inexpensive, reliable, and high-quality multichannel DMs.8 The number of channels in a DM can reach 102—103, especially in astronomical AO systems.1,2,4,9 Typically, each DM actuator is driven by an individual control unit, resulting in complex electronics and cabling for a high-order DM. This makes the AO system rather bulky, vulnerable to handling, and very expensive. The total cost of the driver electronics for a high-order AO system may account for two-thirds of the total system cost.9

An attempt to simplify the DM electronics by sequential addressing of actuators was made by Kibblewhite et al.10 They built a 59-channel piezoelectric faceplate mirror that was driven by several high-voltage amplifiers (HVA) by way of 16 high-voltage switches assembled from discrete components. Due to the implementation and control complexity, this approach has not won general acceptance.2

In this letter, we report on the possibility of a high-order piezoelectric DM with internal multiplexing and simple driver electronics that implements adaptive addressing. This approach implies that the addressing of piezoelectric actuators is accomplished not continuously, but as needed through a pregenerated lookup table to follow the required DM figure. Only those actuators that must adjust the DM shape are addressed in each cycle. Other actuators are updated at the slowest possible rate, to keep their size. The simplicity and compactness of the multiplexing electronics allows it to be potentially integrated with the mirror, reducing the complexity of cable interconnections.

In our experiment, we used a 12-channel experimental piezoelectric DM with a 25.4-mm aperture from OKO Technologies.11 The mirror uses standard 3.2-mm-diameter and 30-mm-long tubular PZT actuators (produced by PI Ceramic), positioned in a rectangular grid with 7-mm pitch. With a control voltage ranging from 0 to 300 V, the DM full stroke reaches \( \sim 7 \mu \)m (hysteresis \(<10\%\)), and its interactuator stroke, i.e., the maximum displacement between the adjacent actuators, amounts to \( \sim 2 \mu \)m. The actuators have comparatively low capacitance of \( C_m=12 \) nF, with \( 1/e \) leakage time exceeding 30 s while mounted in the mirror.

The driver electronics for the DM (see Fig. 1) includes a midpower unipolar MOSFET-based HVA, loaded by a 12-channel multiplexer. The HVA is capable of driving a single actuator with 500-V amplitude at 18 kHz. Actuators were multiplexed using 12 miniature high-voltage optotriacs (Q1—Q12) with optically isolated low-voltage controls. A computer-integrated 8-bit DA converter board was used to generate all the control signals for the mirror operation. Figure 2 shows the simplified timing diagram for 1, 2, and 12 channels of the DM. As seen, in a multiplexing period, the voltages \( V_1 \) to \( V_{12} \) produced by the HVA are sequentially applied to the DM channels via the optotriac switches Q1—Q12 activated by channel selection signals.

The operating bandwidth and surface stability are the critical parameters of the multiplexed DM. To characterize the multiplexed DM in open-loop mode, we used the setup shown in Fig. 1. A light beam from a He-Ne laser at \( \lambda \approx 633 \) nm is collimated and split by a 50/50 beam splitter BS1. Part of the collimated beam illuminates a low-finesse asymmetric Fabry-Pérot interferometer formed by the DM mirror and a second 50/50 beam splitter BS2. The light reflected from the interferometer is picked up by BS2 and goes through the imaging optics (IO) to a CCD for fringe registration, or to a photodiode for measurements of the surface stability.

Limited bandwidth of the driver electronics and mechanical DM resonances reduce the overall DM performance at high frequencies. In each situation, there is a threshold frequency (\( F_{th} \)) at which the multiplexing leads to interchannel crosstalk and temporal instability of the mirror profile. Figure 3 presents the interferograms of the DM influence function at (a) \( F< F_{th} \) and (b) \( F> F_{th} \). In this experiment, a voltage of 150 V was applied to the DM channel 8 (see Fig. 1), while all other channels were kept at 0 V. The threshold frequency of multiplexing was found to be \( F_{th} \approx 1.5 \) kHz without adaptive optimization of actuator addressing.

The DM operation with a large interactuator stroke can be considered as the extreme case of multiplexing. Actually, in this case a high voltage difference (\( U_A \)) is required between adjacent DM actuators, and the HVA needs to perform full-amplitude switching at its maximum frequency. This frequency is given (see Fig. 2) by:

\[
F_{th} = \left( \frac{1}{t_{HVA} + t_{on} + t_{off}} \right)^{-1},
\]

where \( t_{HVA} \) is the HVA settling time, and \( t_{on} \), \( t_{off} \) are the switching times of optotriacs. At \( U_A=300 \) V, we have \( t_{HVA}=60 \mu s, t_{on}=8 \mu s, t_{off}=30 \mu s, \) and \( F_{th} \) reaches...
The experimental dependence of $F_{th}$ on $U_{\Delta}$ (and the corresponding interacuator stroke) is shown in Fig. 4. In the experiment, $U_{\Delta}$ was applied to the even channels of the mirror, whereas the odd channels were maintained at 0 V. As seen, a maximum interacuator stroke of $\sim 2 \mu m$ that corresponds to $U_{\Delta}=300$ V was obtained at $F < F_{th}=700$ Hz. This value is in agreement with the $F_{th}$ estimate above. A stroke of such magnitude, however, is rarely needed in common AO applications. At the opposite extreme, the multiplexed DM provides a $0.1$-$\mu m$ interacuator stroke $U_{\Delta}=12$ V at $F_{th}=3.6$ kHz. The equivalent Nyquist frequencies (i.e., single-channel bandwidths of the multiplexed DM) are 350 Hz and 1.8 kHz, respectively. These values meet or exceed the requirements for typical atmospheric correction.2,10

Figure 4 shows that the maximum update rate $F_{th}$ can be adjusted, depending on the required DM deformation. For small displacements, the addressing can be carried out at higher frequencies, and vice versa. This allows optimization of the overall mirror performance by choosing an appropriate update rate for each DM channel.

To estimate the figure stability of the multiplexed DM, the surface displacement error ($\Delta d$) was evaluated as a function of $F$ for the DM with a single activated channel. We estimated the measured intensity variation ($\Delta I$) of the lowest-order fringe of the DM interference pattern (see Fig. 3), to determine very small displacements $\Delta d$ of the mirror surface, caused by the interchannel crosstalk. The reflected intensity ($I+\Delta I$) of an asymmetric Fabry-Pérot interferometer,12 can be written as:

$$I + \Delta I = I_0 R(d_0 + \Delta d),$$

where $R(d) = \left| \frac{-\sqrt{R_1 + \sqrt{R_2}} \exp(2ikd)}{1 - \sqrt{R_1 R_2} \exp(2ikd)} \right|^2$.

$I_0$ is the intensity of incident light; $d_0$ is the sum of the DM stroke caused by the addressed actuator and the distance between BS$_2$ and DM; $k=2\pi/\lambda$ is the wave number; and $R_1$, $R_2$ are the reflectivities of BS$_1$ and DM, respectively. Assuming that the surface displacement error is small, i.e., $|\Delta d| \ll \lambda$, Eq. (1) yields: $\Delta I \approx \gamma \Delta d$, where $\gamma = I_0 R(d_0)/\partial d$ is the interferometer sensitivity at $d_0$. The inset in Fig. 5 illustrates intensity variation with $\Delta d$. The sensitivity parameter $\gamma$ can be calibrated by applying a low-frequency AC voltage with the amplitude $\Delta U$ to the selected PZT actuator at $F \approx F_{th}$. Taking into account that $d_0$ linearly

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**Fig. 1** Experimental setup. P, polarizer; SF, spatial filter; BS$_1$, BS$_2$, beam splitters; IO, imaging optics; DM, deformable mirror; HVA, high-voltage amplifier; Q$_1$—Q$_{12}$, optotriacs.

**Fig. 2** Simplified timing diagram for 1, 2, and 12 channels of the DM.

**Fig. 3** Averaged over 1-s period interferograms of the piezo DM at (a) low-refresh frequency $F=5$ Hz and (b) high-refresh frequency $F=1700$ Hz, above the threshold value $F_{th} \sim 1500$ Hz.

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changes with applied voltage, we obtain:
\[\gamma = 2U_{\lambda 2} \Delta U / (\lambda \Delta U),\]
where \(U_{\lambda 2}\) is the voltage required for a
\(\lambda/2\) shift (see inset in Fig. 5). In our case, \(U_{\lambda 2}\) is measured to be \(\sim 38\) V. Using these relations, the required surface
displacement error \(\Delta d\) can be evaluated.

The experimental dependence of \(\Delta d\) on \(F\) is shown in
Fig. 5. In this case, a voltage of 20 V is applied to the DM
channel 8 while multiplexing. As seen, the displacement
error does not exceed \(\sim \lambda/100\) rms at multiplexing
frequencies of less than 3.5 kHz. Higher multiplexing
frequencies result in the loss of the DM figure stability.

An important advantage of multiplexing control is its
scalability to \(10^{\text{th}}\)—\(10^{\text{th}}\) DM channels. This is of particular
interest for piezoelectric DMs that employ low-capacitance
actuators \((\sim 10\,\text{nF})\) operating at frequencies \(<1\,\text{kHz}\). In
this case, low-power switching electronics generating low
heat can be integrated with the mirrors. Actually,
for our system the average power \((W_d)\) dissipated by a
single switch can be estimated as:
\[W_d = t_{\text{on}}^2 P_{\text{on}} + t_{\text{off}}^2 P_{\text{off}},\]
where \(t_{\text{on}}, t_{\text{off}}\) are the charging and discharging
currents, respectively; \(t_{\text{on}}, t_{\text{off}}\) are the switching
times (see Fig. 2); and \(P_{\text{on}}/P_{\text{off}}\) is the optotriac output resist-
ance. For an actuator operated with a maximum voltage of
300 V at \(F = 700\,\text{Hz}, t_{\text{on}} = 8\,\mu\text{s}, t_{\text{off}} = 30\,\mu\text{s},\) we obtain
\(W_d = 9\,\text{mW}\). With these parameters, the power \(W_{\text{PTZ}}\) gen-
erated in the actuator is \(\pi \delta^2 \rho^2\), where \(\delta = 12\,\text{mW},\) for a
loss factor \(\gamma \sim 0.01). Thus, the heat produced by the
switch is lower than the actuator dissipation. The example
above gives the maximum estimate for the heat dissipation;
practical numbers will be at least two orders of magnitude
lower. As discussed by Aldrich,2 the multiplexing of high-
capacitance actuators has no significant advantages over di-
rect control.

In conclusion, we present the adaptive multiplexing
control for piezoelectric deformable mirrors. The adaptive
addressing of actuators through the program-based lookup
table is implemented to obtain a stable DM figure at a high

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